



This is a digital copy of a book that was preserved for generations on library shelves before it was carefully scanned by Google as part of a project to make the world's books discoverable online.

It has survived long enough for the copyright to expire and the book to enter the public domain. A public domain book is one that was never subject to copyright or whose legal copyright term has expired. Whether a book is in the public domain may vary country to country. Public domain books are our gateways to the past, representing a wealth of history, culture and knowledge that's often difficult to discover.

Marks, notations and other marginalia present in the original volume will appear in this file - a reminder of this book's long journey from the publisher to a library and finally to you.

### **Usage guidelines**

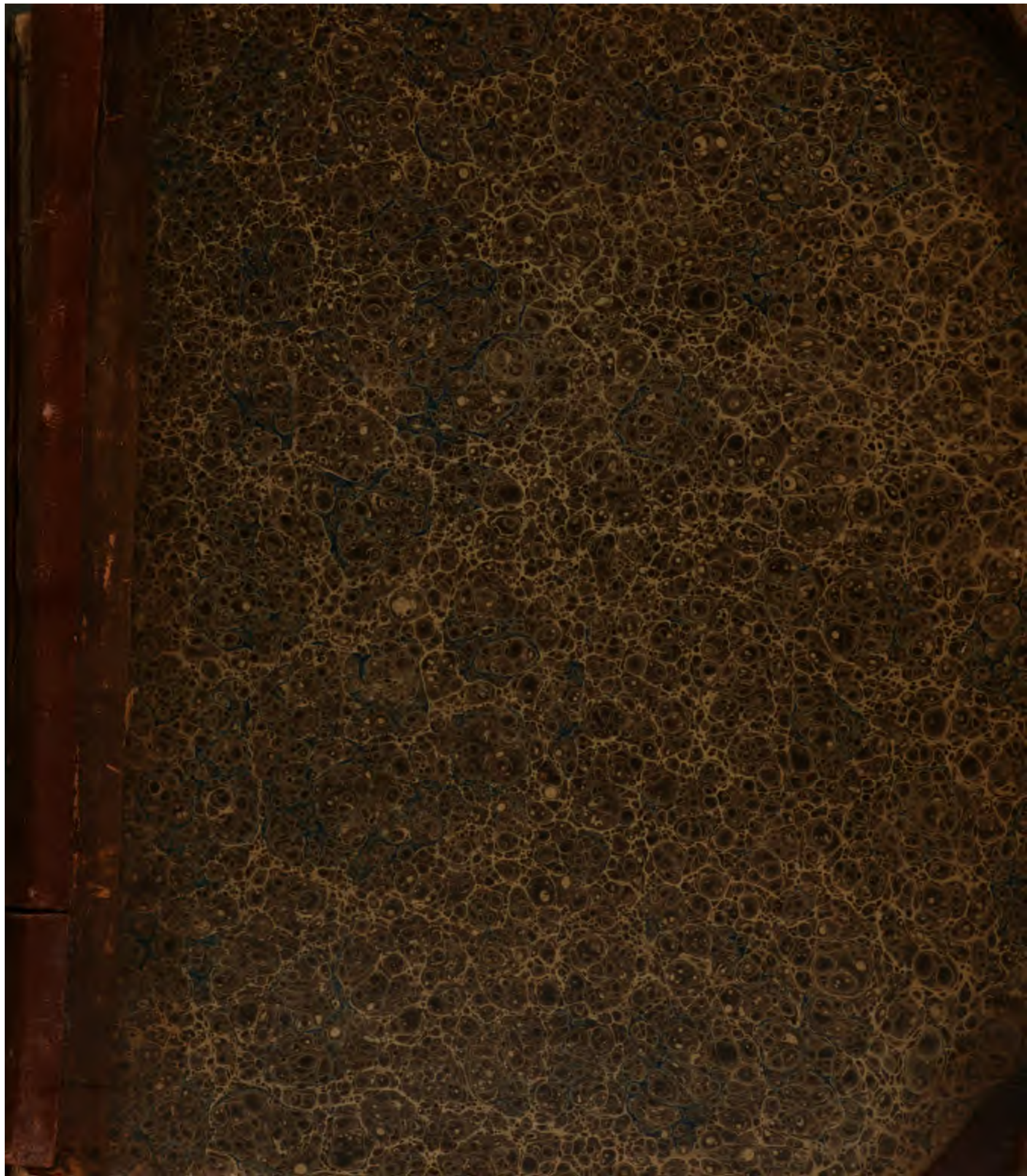
Google is proud to partner with libraries to digitize public domain materials and make them widely accessible. Public domain books belong to the public and we are merely their custodians. Nevertheless, this work is expensive, so in order to keep providing this resource, we have taken steps to prevent abuse by commercial parties, including placing technical restrictions on automated querying.

We also ask that you:

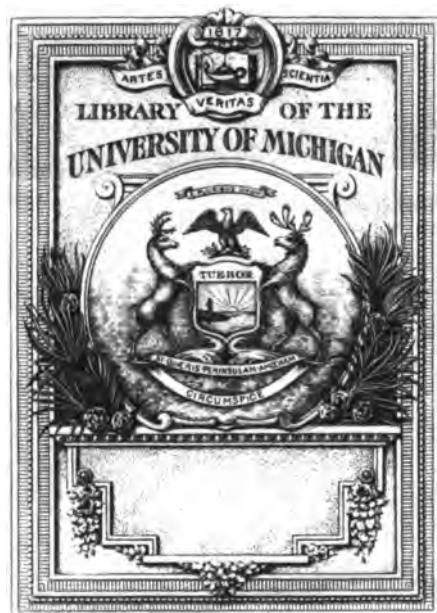
- + *Make non-commercial use of the files* We designed Google Book Search for use by individuals, and we request that you use these files for personal, non-commercial purposes.
- + *Refrain from automated querying* Do not send automated queries of any sort to Google's system: If you are conducting research on machine translation, optical character recognition or other areas where access to a large amount of text is helpful, please contact us. We encourage the use of public domain materials for these purposes and may be able to help.
- + *Maintain attribution* The Google "watermark" you see on each file is essential for informing people about this project and helping them find additional materials through Google Book Search. Please do not remove it.
- + *Keep it legal* Whatever your use, remember that you are responsible for ensuring that what you are doing is legal. Do not assume that just because we believe a book is in the public domain for users in the United States, that the work is also in the public domain for users in other countries. Whether a book is still in copyright varies from country to country, and we can't offer guidance on whether any specific use of any specific book is allowed. Please do not assume that a book's appearance in Google Book Search means it can be used in any manner anywhere in the world. Copyright infringement liability can be quite severe.

### **About Google Book Search**

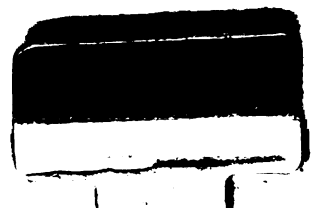
Google's mission is to organize the world's information and to make it universally accessible and useful. Google Book Search helps readers discover the world's books while helping authors and publishers reach new audiences. You can search through the full text of this book on the web at <http://books.google.com/>



712a



math  
(u)



QA  
3  
L25

100

# MATHEMATICAL MEMOIRS

RESPECTING

A VARIETY OF SUBJECTS;

WITH AN

A P P E N D I X

CONTAINING

T A B L E S of T H E O R E M S

FOR THE

C A L C U L A T I O N of F L U E N T S.

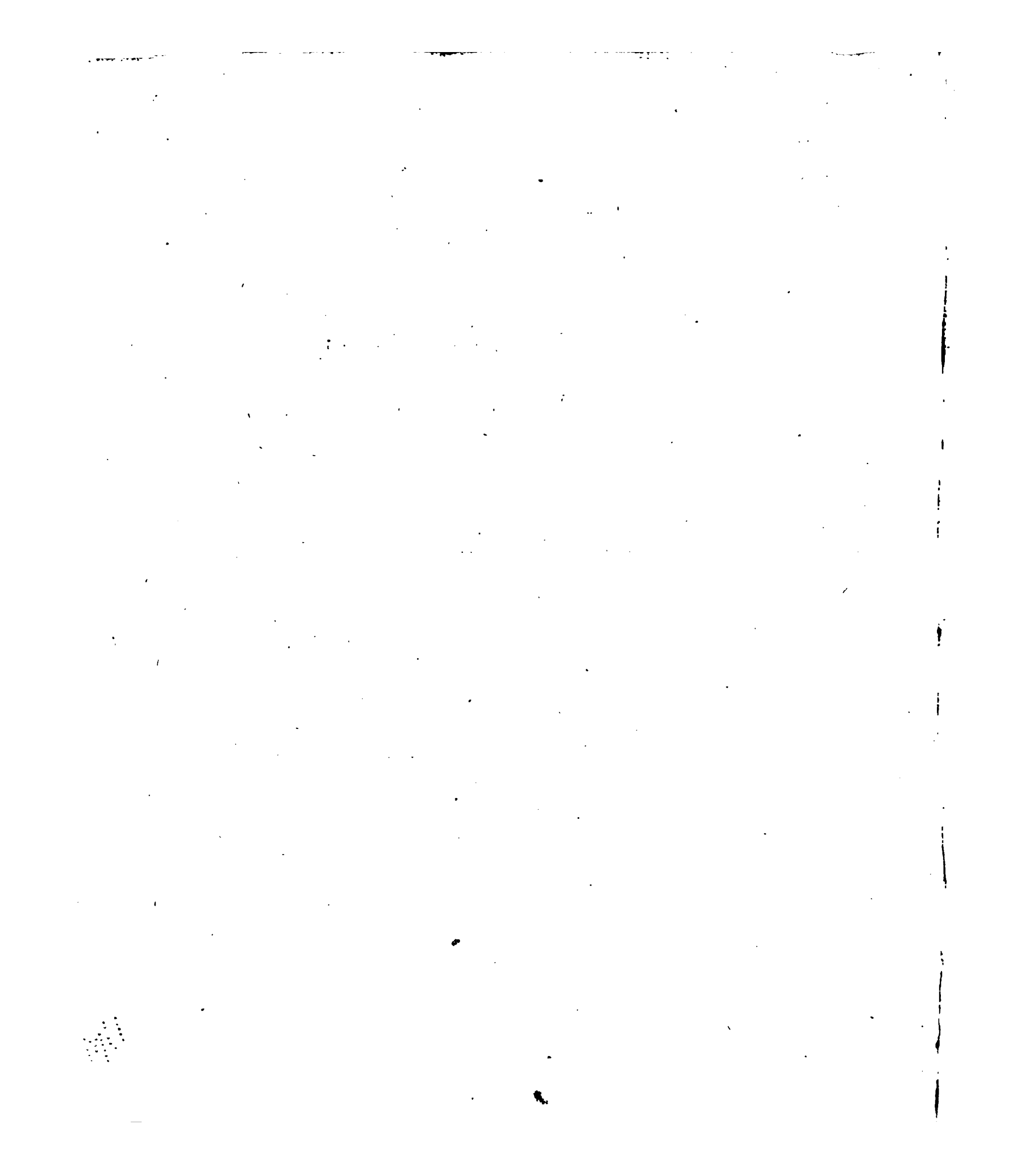
VOL. I.

By JOHN LANDEN, F. R. S.

---

L O N D O N,

Printed for the AUTHOR; and sold by J. NOURSE, Bookfeller to  
HIS MAJESTY. 1780.







afford many new geometrical and philosophical Improvements of considerable importance.

The new demonstration of the property of the straight Lever, and some other articles, in the first Memoir, I conceive, may probably give pleasure and satisfaction to some readers, who may be particularly curious in such disquisitions.—But I almost assure myself, there are few geometers who will not be pleased with the discovery of the theorem in the second Memoir, which enables us to *assign the length of any arc of any conic hyperbola by means of two elliptic arcs*:—a thing which, I am inclined to believe, had not even been thought possible by former writers on the properties of those curves; and of whose use there are many remarkable instances in the following pages.

The new theorems in the third and sixth Memoirs, respecting the motion of a Pendulum, I am induced to think, will not be unacceptable to readers who understand what had before been published concerning such motion.—And I am persuaded that the new method of computation, whereby the sums of many series are obtained, in the fifth Memoir, will engage the attention of the intelligent analyst; and incite him to exercise his skill in the farther application of that method.

With regard to such matters (in these Memoirs) as have been considered by other writers, I must particularly acknowledge, that my fourth Memoir is on a subject upon which the *celebrated Mr. D'ALEMBERT*, as I have lately found, has written at some length, in his *Opuscules Mathematiques*: yet, I flatter myself, that the critical reader will find, in that Memoir, a great number of new and interesting articles, amongst which are some very remarkable

P R E F A C E.

markable inferences respecting the gyration of certain bodies there specified.

We may understand by what is said in the fifth Tome of the *Opuscles* just now named, that, after the perusal of what had been written on the subject, a doubt remained with some mathematician, whose name is not there mentioned,—*whether there be any solid, besides the sphere, in which any line whatever, passing through its center of gravity, will be a permanent axis of rotation?*—Upon that point (which some perhaps may think is not very satisfactorily explained by Mr. D'ALEMBERT\*) I have touched a little in the *Philosophical Transactions* for the year 1777; but, in the Memoir here alluded to, I have, it is presumed, so fully explained the matter as to obviate or remove every doubt concerning it.

In the seventh, eighth, and ninth Memoirs, in which the motion of a projectile is considered, the reader will find some propositions before considered by many authors; nevertheless I persuade myself, that what I have written respecting those propositions will not be deemed trite and uninteresting.—Moreover, there are (in those Memoirs) some *new* researches concerning the motion of bodies, which the ingenious mathematician may possibly find not unworthy of his regard.

It may be observed, that the common doctrine of centripetal forces comprehends only a part of what is necessary to be understood in order to determine in general the path of a projectile and its motion therein, from a knowledge of its velocity and direction at any given time.

\* He has not given one instance of a body having the remarkable property in question!

and

and of the force or forces acting thereon whilst moving;—there being innumerable cases in which such force or forces may not continually urge the body towards a certain center as the said common doctrine supposes.—The deficiency of that doctrine in the books published on that subject, I have endeavoured in some measure to supply in the three Memoirs last mentioned:—and, as a farther application of the principal theorems in those Memoirs may be requisite to explain sufficiently the general doctrine of a projectile's motion, I purpose to make such application in some subsequent Memoirs respecting propositions too intricate to be considered amongst the examples which I thought proper to be given in the Memoirs wherein those principal theorems are investigated.

Supposing the reader to know what is meant by any term commonly used in mathematical writings, I shall seldom formally define the technical terms I may make use of in these Memoirs; but, that I may not be misunderstood, I shall endeavour particularly to explain my meaning, where, by using a new term or an ambiguous old one, I apprehend some explanation is necessary:—my design being to treat the subjects I may take upon me to consider—perspicuously, but not prolixly;—in the prosecution of which, I shall lay down the principles requisite in computations, and shew how from thence conclusions may be readily deduced, as well in such propositions as are generally reckoned abstruse as in such as are esteemed more easily investigable.

J. LANDEN.

---

---

C O N T E N T S

O F

V O L U M E I.

- M**EM. I. *Of the Mechanic Powers, so far as relates to Equilibriums* - - Page 1
- MEM. II. *Of the Ellipsis and Hyperbola* - 23
- MEM. III. *Of the Descent of a Body in a Circular Arc* - - - - 37
- MEM. IV. *Of the centrifugal Force of the Particles of a Body, arising from its Rotation about a certain Axis passing through its Center of Gravity* - 41
- MEM. V. *A new Method of obtaining the Sums of certain Series* - - - - 67
- MEM. VI. *A remarkable new Property of the Cycloid discovered, which suggests a new Method of regulating the Motion of a Clock* - - - 119
- MEM. VII. *Of the Motion of a Body, keeping always in the same given Plane, whilst acted on by any Force, or Forces, urging it continually to change its Direction in that Plane* - - - - 123
- MEM. VIII. *Of the Motion of a Body in (or upon) a Spherical Surface; in (or upon) which it is retained by*

C O N T E N T S.

<i>by some Force urging it towards the Center of the Sphere, whilst it is continually impelled by some other Force, or Forces, to change its Direction in (or upon) that Surface</i>	- - - - -	158
MEM. IX. <i>Of the Motion of a Body in any variable Plane</i>	- - - - -	173

MATHE-

---

---

M A T H E M A T I C A L  
M E M O I R S.

---

M E M O I R I.

*Of the Mechanic Powers, so far as relates to Equilibriums.*

**W**RITERS on the mechanic powers have, generally, founded their demonstrations of the properties of those powers, on a principle which has been objected to, as obscure and unnatural. For, in treating of Equilibriums (where no moving bodies act on each other, or are any way concerned in the enquiry), they have neglected the proper principles of that doctrine, and have borrowed a foreign, less evident, one from a consideration of motion. They infer from the doctrine of motion, that

“ as those bodies are equipollent in the congress and re-  
“ flexion, whose velocities are reciprocally as their innate  
“ forces: so, in the use of mechanic instruments, those  
“ agents are equipollent and mutually sustain each the  
“ contrary pressure of the other, whose velocities, esti-

B

“ mated

“mated according to the determination of the forces, are reciprocally as the forces.” This, properly understood, is indeed true; and being admitted, renders the business of the writer on the subject of those instruments very easy: yet, as it is not a clear and natural inference, but rather a theorem, wanting a demonstration, assumed as a principle; and many have expressed a dissatisfaction at the manner in which this subject is usually treated; it may be of use to consider the matter in a different light, and to build our demonstrations on principles more natural and evident. Such, I presume, are those upon which, without any regard to the doctrine of motion, I purpose to establish the fundamental parts of this doctrine.

As we shall all along, in this memoir, have frequent occasion to consider the tensions of strings drawn over pulleys, it will best suit our purpose to begin with the explaining the properties of those instruments.

1. The tension of a string, by what means soever it be stretched, is said to be equal to that weight which would stretch it just as much, being fastened to one end of it and sustained by it, whilst the string itself is suspended at rest, in a vertical position, by its other end.

2. Any string, which is no where fastened to any thing but at its ends, will, in every part of it, be equally stretched by any force acting upon it, supposing it void of gravity.

3. If a string ABCDEFG be drawn over the immovable pulleys B, C, D, E, F; and one end of it (G) being fastened to the fixed point G, a weight A be fastened to the other end; the string will in every part be equally stretched, and the parts AB, BC, CD, EF, FG, of the string being all supposed perpendicular to the horizon, the  
point:

Plate I.  
Fig. 1.

point G will be pulled upwards by a force equal to (A) the string's tension; and, the reaction of the point G being equal to (A), the action thereon, it evidently follows, that each of the lower pulleys (C, E) will be pulled upwards, and each of the upper pulleys (B, D, F) pulled downwards, by twice that force.

4. The reaction of the point G being equal to A, the action thereon, and the reaction of each of the pulleys being equal to  $2A$ ; if, instead of fastening the end G to an immoveable obstacle, a weight equal to A were suspended by it; or if any pulley, instead of being fastened to an immoveable obstacle, were pulled against the string ABCDEFG by a suspended weight equal to  $2A$ , the resistance opposing the action of that string would be the same, and the string would be stretched as before, and remain at rest. Fig. 2.  
Fig. 3.

5. Any of the lower pulleys, with or without the end G, being fastened to one weight only; it follows, that if that weight be to the weight A as the number of parts (BC, CD, &c.) of the string passing between the greater of the two weights and the upper pulleys is to unity, the string will have the same tension as if the pulleys and end G were all fastened to the immoveable obstacle; and it will likewise remain without motion, its action being equally resisted in both cases. Fig. 4.  
Fig. 5.

Thus it appears, that, by drawing a string over several pulleys, it may be made to sustain various weights, whilst its tension remains the same; and that, by such means, a very great weight may be made to rest in equilibrio with a much lesser one.

6. If a string ABCD be drawn over the pulleys B and C, and, one end of it being fastened to a fixed point D, Fig. 6.

B 2

a weight



a weight  $A$  be suspended at the other end; the pulley  $B$  being an immoveable one, and  $C$  fastened to one end  $c$  of another string  $cde$ , which is drawn over an immoveable pulley  $d$ , and has its other end fastened to a fixed point  $e$ ; and the parts  $AB$ ,  $BC$ ,  $CD$ ,  $cd$ ,  $de$ , of the strings being all perpendicular to the horizon; the tension of the string  $ABCD$  will be equal to  $A$ , and the pulley  $C$  will be pulled upwards by twice that force. Consequently the string  $cde$  opposing the ascent of that pulley ( $C$ ) will have its tension equal to  $2A$ , and the pulley  $d$  will be pulled upwards by a force equal to  $4A$ .

**Fig. 7.** Therefore, according to what has been before observed, if instead of the pulley  $d$  being fastened to an immoveable obstacle, a weight equal to  $4A$  be suspended by it, the action of the string  $cde$  will be resisted as before; and the weights  $A$  and  $4A$ , it is evident, will rest in equilibrio.

The application of this method of reasoning to any combination of pulleys (whether one, two, or more strings be employed) is so easy, that I think it unnecessary to insist any longer on this head.

**Fig. 8.**  $AB$  being an inflexible rod (considered without weight) on which three forces act, at right angles thereto, by means of three parallel strings  $Aa$ ,  $Bb$ ,  $Cc$ , fastened at any three points  $A$ ,  $B$ ,  $C$ , of the rod, two of them on one side opposing the third on the other side, keeping the rod at rest: I now propose to investigate (without any regard to the doctrine of motion) the ratio of the two forces acting on the same side of the rod; whose sum, it is plain, must be equal to the other force acting on the contrary side.

**Fig. 9.** 7. If the weight  $W$  be suspended from the middle of the inflexible horizontal rod  $BD$ , which is itself suspended by

by two parallel strings AB, CD; these strings will be equally stretched, and each will bear half the weight W. And CD being fastened at C, and AB drawn over the pulley A, the weight N appending to this string, requisite to keep an equilibrium with W, must be equal to half W.

8. The weight W being suspended from any point P of Fig. 10. the inflexible horizontal rod BD, which is itself suspended by two parallel strings AB, CD; if the tension of the string AB be denoted by T, the tension of CD will be expressed by  $W - T$ ; and the ratio of T to  $W - T$  will be the same, let W be what it will; for the tension of each string will be increased or diminished in the same proportion in which W shall be increased or diminished.

9. If BQ be equal to DP, any force acting at Q at right angles to BD will affect the tensions of the strings AB, CD respectively, in the same proportion as that in which the tensions of the strings CD, AB would be respectively affected by the application of any force at P in a parallel direction.

10. Therefore, if a string QR, parallel to the other three, having one end fastened to the rod at Q, be drawn over the pulley R by such a weight N hung to the other end of it, that the same be just sufficient, with the string CD, to sustain the rod, with the weight W appending thereto, in its former position; the string AB then having no tension, the weight which it bore being now borne by the string QR, together with part of the weight which before was borne by CD; T, the decrease of tension in AB, will be to  $N - T$  the decrease of tension in CD, as  $W - T$  to T. Consequently,  $W - T$  will be to W as T to N.

DP

DP is here considered as less than half BD, that consideration being sufficient for our purpose.

Fig. 11.

11. If DP be a submultiple of BP, that is, if,  $n$  being some integer,  $n \times BQ$  be equal to BP; let  $BQ$ ,  $Q'Q''$ ,  $Q''Q'''$ , &c. be each equal to DP. Then applying strings  $Q'R'$ ,  $Q''R''$ , &c. successively to the points  $Q'$ ,  $Q''$ , &c. of the lever, as described in the preceding article, and denoting the respective tensions of those strings as they are applied by  $N'$ ,  $N''$ , &c. we shall have.

$$W - T : W :: T : N' = \frac{TW}{W - T},$$

$$W - N' : W :: N' : N'' = \frac{N'W}{W - N'} = \frac{TW}{W - 2T},$$

$$W - N'' : W :: N'' : N''' = \frac{N''W}{W - N''} = \frac{TW}{W - 3T},$$

&c.

&c.

&c.

And, continuing the operation to the necessary length, we have  $N^{n-1} = \frac{TW}{W - n-1 \times T}$ , which (by art. 7.) will be equal to half  $W$ : whence  $n \times T = W - T$ , or  $T : W - T :: 1 : n$ .

Thus it appears, that, in case BP be any multiple of PD, the tensions of the strings AB, CD (which are denoted by  $T$  and  $W - T$ ) will be as 1 to  $n$ , that is, as DP to BP.

12. From what has been said it evidently follows, that, DP being any distance whatever, if BQ, instead of being equal thereto (as hitherto supposed) be equal to  $q$ , any submultiple of  $r - q$ ,  $r$  being the whole length of the rod BD;  $r - q$  will be to  $q$ , as  $T$ , the decrease of tension in AB, to  $N - T$ , the decrease of tension in CD, upon applying a string QR at Q as above. Whence  $r - q : r :: T : N$ .

13. We

13. We will now see what will be the ratio of the tensions of the strings AB and CD, when BP is not a multiple of DP. First let us suppose DP and BP commensurable, and equal to  $m \times p$  and  $n \times p$  respectively,  $p$  being a common measure of BP, DP. Then conceiving a string, as QR, applied successively to the points whose distances from B are  $p, 2p, 3p, \&c.$  till it comes to a point equidistant from P with the point D, and determining the several successive values of  $N', N'', \&c.$  (the tensions of  $Q'R', Q''R'', \&c.$ ) in terms of T (the tension of AB) and known quantities, the last of them will (by art. 7.) be equal to  $\frac{W}{2}$ : from which equation the ratio of T to  $W - T$  will appear. Thus, applying the string  $Q'R'$  to the point  $Q'$ , whose distance from B is  $p$ , we shall, by the preceding article, have  $r - p : r :: T : \frac{rT}{r-p} = N'$ , the weight then borne by  $Q'R'$ . Applying the string  $Q''R''$  to the point whose distance from B is  $2p$ , we shall have  $r - 2p : r - p :: \frac{rT}{r-p} : \frac{rT}{r-2p} = N''$ , the weight borne by  $Q''R''$  in that situation. And it is obvious, that the weight which the string QR will bear, when applied to the point whose distance from B is  $n - m \times p$  (that is, when its distance from P is equal to DP) will be  $\frac{rT}{r - n - m \times p} = \frac{m + n \times T}{2m}$ , which (by art. 7.) must be equal to  $\frac{W}{2}$ . Consequently T will be to W as  $m$  to  $m + n$ , and T to  $W - T$  as  $m$  to  $n$ . Therefore, T being the weight borne by the string AB before the application of  $Q'R', Q''R'', \&c.$  and  $W - T$  the weight borne by the string CD at the same

Fig. 12.

same time, it is manifest, that the tensions of those strings (AB, CD) when they sustain any weight W, appending to the point P, will be reciprocally as their distances from that point, though DP be not a submultiple of BP, if those distances be commensurable.

Fig. 13.

14. Suppose now that DP and BP are incommensurable: and if the tension of AB be not to the tension of CD as DP to BP, let it be as DP to *b*P greater or less than BP. Let *Dm*, less than *Bb*, be a submultiple of DP: and *Dn* being the least multiple of *Dm* which exceeds BD, if *b*P be supposed greater than BP; or the greatest multiple of *Dm* which is less than BD, if *b*P be supposed less than BP; *Bn* will in either case be less than *Dm*, and *n* will fall between B and *b* in both cases. Now, by what has been proved, if the string AB were at *n*, its tension to that of CD would be as DP to *n*P: therefore, if *d* be put for the difference of the tensions of AB, when applied at B, and when applied at *n*; and *t* be its tension when at B, and *T* that of CD at the same time; we shall have  $t \mp d : T \pm d :: DP : nP$ , and  $\frac{t \mp d}{T \pm d} \times nP = DP$ . But by hypothesis,  $t : T :: DP : bP$ , and  $\frac{t}{T} \times bP = DP$ ; and consequently  $\frac{t \mp d}{T \pm d} \times nP$  should be  $= \frac{t}{T} \times bP$ . This, it is plain, is impossible: for  $\frac{t \mp d}{T \pm d}$  is less or greater than  $\frac{t}{T}$ , according as *n*P is less or greater than *b*P. The ratio of the tensions of the strings can then be no other than reciprocally as their distances from P.

15. Instead of any force whatever acting on a body by other means than that of a string, we may conceive another  
another

another force to be substituted, which pulling by a string in the same direction, the same effect shall be produced; and forces being measured by the effects they produce, this force must be esteemed equal to that in whose stead it may be so substituted. Therefore, whatever we infer concerning the ratio of the tensions of strings, by means whereof any body is kept at rest, the same may be inferred of the ratio of any other forces acting in the same directions, and producing the same or an equal effect. Consequently, a straight inflexible rod being kept at rest by three forces acting thereon, at right angles thereto, whether those forces act by means of strings or otherwise, the two forces opposing the third (which must necessarily act between them) must be reciprocally as the distances of their points of action from the point at which the third force acts.

Hence the properties of straight levers, the power and weight acting at right angles thereto, are evident.

16. AB being a lever of the first kind, F the fulcrum, and M and N two weights suspended from A and B respectively, and sustaining each other in equilibrio with AB in a horizontal position; the lever may be considered as kept at rest by three forces M, N, and  $M + N$  acting at A, B, and F respectively in the manner above described. For the reaction or resistance of the fulcrum is of the same effect with respect to sustaining M and N as an active force equal to those two would be of, if it opposed them by pulling or pressing at right angles to the lever at the point F. Therefore, by the preceding article, M must be to N reciprocally as AF to BF. And in the same manner it appears, by changing the place of the fulcrum, that any forces whatever, which, acting on a straight lever of the second or third kind at right angles thereto, are equipol-

C

lent,

Fig. 14.  
Fig. 15.

lent, are reciprocally as the distances at which they act from the fulcrum.

Fig. 16.

17. Instead of considering the lever as a single straight line, let us conceive the fulcrum  $F$  to be at a point in an extended plane of any form (without weight), and that the said plane is freely moveable about  $F$  as a center. Then, if the weights  $M$  and  $N$ , acting at the points  $B$  and  $D$  in the said plane at right angles to the right line  $BFD$ , rest in equilibrio; they will, it is obvious, likewise rest in equilibrio, if, instead of acting at  $B$  and  $D$ , they act at any other points  $b$  and  $d$  respectively, in the lines  $Bb$ ,  $Dd$ ; these lines being the directions in which the forces act, whether  $B$  and  $D$  or  $b$  and  $d$  be the points of action.

Moreover, the weights  $M$  and  $N$  will also rest in equilibrio, if,  $FE$  being equal to  $FD$ , the force  $N$ , instead of acting at  $D$  in any direction  $Dd$ , act at  $E$  in the direction  $Ee$ , the angles  $FEe$ ,  $FDd$  being equal.

Now, by what is proved in the preceding article,  $BF \times M$  will be  $= DF \times N$ . Join  $Fd$ , and suppose that, instead of the weight  $N$  acting at  $d$  in the direction  $dN$ , a weight equal to  $\frac{DF \times N}{dF}$  act at  $d$  at right angles to  $dF$ . Then, by what is already proved, will this last mentioned weight rest in equilibrio with the weight  $M$  acting at  $B$  or  $b$  as before. Therefore, with regard to preserving an equilibrium with any opposing force, the weight acting at  $d$  on the lever  $dF$  in the direction  $dD$ , or  $Dd$ , is to the weight acting at the same point of the same lever at right angles thereto, as  $N$  to  $\frac{DF \times N}{dF}$ ; that is, as  $dF$  to  $DF$ , or as radius to the sine of the angle which the oblique direction makes with the ray from the fulcrum to the point of action.

18. Let

18. Let two forces, G and H, act obliquely at the points B and D, on the straight lever BFD, moveable about the fulcrum F; and BA, DC, being the directions in which those forces act, let the sines of the angles ABF, CDF, to the radius 1, be denoted by  $p'$ ,  $p''$ , respectively. Then, by what is said in the preceding article,  $p' \times G$  and  $p'' \times H$  will be the respective equivalent forces considered as acting at the same points (B and D) at right angles to the lever. And, the lever being supposed to be kept in equilibrio,  $p' \times G \times BF$  will be equal to  $p'' \times H \times DF$ . Fig. 17.

19. Considering now three forces, G, H, I, as acting at the point P in an extended plane of any form (without weight), in the directions PA, PB, PC respectively, so as to preserve an equilibrio; the ratio of those forces will be found as follows, by means of what is done above. Continue BP to any point F, and draw  $dFe$  at right angles thereto. Then, seeing that it is the same thing with regard to keeping the plane at rest, whether the forces (G, H, I) act at the points  $d$ , F,  $e$ , in the said plane respectively, or at P, as before mentioned, their directions being the same in both cases; we have, by the preceding article,  $p' \times G \times dF$  equal to  $p''' \times I \times eF$ ,  $p'$  and  $p'''$  being the sines of the angles  $AdF$ ,  $CeF$  respectively, to the radius 1. Whence it is evident, that the forces G and I must be as  $p''' \times eF$  and  $p' \times dF$ . But if  $q'$ ,  $q'''$  be the respective cosines of  $AdF$ ,  $CeF$ , or the sines of  $dPF$ ,  $ePF$ ;  $dF$  will be  $= \frac{q' \times FP}{p'}$ , and  $eF = \frac{q''' \times FP}{p'''}$ . It appears therefore, by substitution, that the force G must be to the force I as  $q'''$  to  $q'$ . Fig. 18.

And, continuing AP to E, and denoting the sine of  $ePE$  by  $q''$ , it appears by the same way of reasoning, that



the force H must be to the force I as  $q''$  to  $q'$ ; the angle BPE being manifestly equal to the angle  $dPF$ . Consequently the forces G, H, I, must be as  $q'''$ ,  $q''$ ,  $q'$ ; that is, drawing BE parallel to PC, as the sines of the angles EBP, BEP, BPE; or, which is the same thing, those forces must be respectively as the sides EP, BP, BE of the triangle BEP which are in, or parallel to, their respective directions.

Fig. 19.

20. If three forces (G, H, I) acting at any three points of a moveable plane sustain the same in equilibrium, the directions in which those forces act (unless they be parallel to each other) will intersect each other in one point only. For let the points of action be  $d, e, f$ ; and the directions  $dA, eB, fC$ . Then if these directions be supposed not to intersect each other in one point, let them be supposed to intersect at  $q, r, s$ , as in our scheme: and let the point of action of the force G be considered as at  $p$  (in  $dqr$ ), instead of  $d$ ; which change, it is plain, will not cause any alteration in the effect of the forces upon the plane. Now, the point  $p$  being between  $q$  and  $r$ , the forces H and I acting in the directions  $sr$  and  $se$ ,  $qs$  and  $sf$  will both urge the plane to turn the same way about the point  $p$ ; and consequently will cause the plane to move. Therefore the supposition is absurd, and the plane cannot remain at rest whilst three forces act upon it at different points, unless their directions (being not parallel to each other) intersect each other in only one point.

21. By means of the last two articles we may readily find the quantity and direction of one single force, which shall have the same effect as any two forces whose quantities and directions are given: and, on the contrary, one single force being given with its direction, we may find two forces, which, acting together in any proposed directions,

directions, shall be equivalent to that single force. Thus, **Fig. 20.** PA, PB denoting the quantities and directions of two forces, draw AQ, BQ parallel to PB, PA respectively; and the diagonal PQ will denote the quantity and direction of a single force equivalent to the other two. And, PQ denoting the quantity and direction of a single given force; draw at pleasure QA, PB parallel to each other; and PA, QB also parallel to each other; then will PA, PB denote the quantities and directions of two forces, which acting together will be equivalent to the said single given force.

22. The wheel AB being fixed to the cylinder Cn, **Fig. 21.** with its center C at the center of a circular section of the said cylinder, and the wheel itself in the plane of that section; and mn being a radius of another circular section of the cylinder, whose center is n; any weight P pulling at A, at right angles to the radius AC, will have the same effect in opposing the action of another weight W, pulling at m at right angles to mn, (and urging the cylinder and wheel to turn the contrary way,) let the distance of C from n be what it will. Therefore, it being manifest that, if C coincided with n, the two radii mn, CA would form a lever mnA; the weights P and W will (by art. 17.) rest in equilibrio, when the former is to the latter reciprocally as AC to mn, whether the radii, to which they act at right angles, have the same position with respect to the horizon or not, the machine being considered as only moveable about the axis (Cn) of the cylinder.

23. AB being an immoveable inclined plane, if a weight **Fig. 22.** W rest thereon in equilibrio with another weight P pulling in the direction WB, parallel to the plane, by means of a string PBW passing over the immoveable pulley B, the weight

weight  $W$  may be considered as kept at rest by three forces acting thereon, one whereof is its own gravity urging it in the direction  $Wm$  directly downwards, another the reprefure of the plane urging it in the direction  $Wn$  perpendicular to  $AB$ , and the third the tension of the string (or the weight  $P$ ) pulling it in the direction  $WB$ . Therefore,  $BC$  being a vertical line, and  $CD$  perpendicular to  $AB$ ,  $BCD$  will be a triangle having its sides in, or parallel to, the directions of those three forces. Which forces will, therefore, (by art. 19.) be respectively as  $BC$ ,  $CD$ , and  $BD$ ; that is, as  $AB$ ,  $AC$ , and  $BC$ , supposing  $AC$  parallel to the horizon.

Fig. 23. 24. The weight  $W$  resting on the immoveable inclined plane  $AB$ , in equilibrio with another weight  $P$  pulling it in the horizontal direction  $Wb$ , by means of a proper contrivance; if  $nWo$  be perpendicular to  $AB$ , and  $Wm$  perpendicular to the horizontal line  $AmoC$ ; the weight  $W$ , the reaction of the plane in the direction  $Wn$ , and the weight  $P$  will (by art. 19.) be to each other respectively as  $Wm$ ,  $Wo$ , and  $mo$ ; that is, if  $BC$  be perpendicular to  $AC$ , as  $AC$ ,  $AB$ , and  $BC$ . Therefore  $P$  must be equal to  $\frac{BC}{AC} \times W$ .

Fig. 24. 25. But if the weight  $W$  be kept at rest on the plane by  $P$  acting at the point  $E$  of a horizontal lever  $EF$ , at right angles thereto, in a direction parallel to  $AC$ ;  $F$  being the lever's fulcrum, the value of  $P$  in this case must, by the property of the lever, be to its value in the former case (considered in the last article) reciprocally as  $EF$  to  $WF$ . Consequently in this case  $P$  must be equal to  $\frac{WF \times BC}{EF \times AC} \times W$ .

26. When

26. When a weight is raised by means of a screw, one part of the screw rises with the weight, and the ascent of that part of the screw is effected by sliding its spiral threads up the spiral threads of the other part of the screw. A weight therefore, when it is sustained in equilibrio by any power acting on the screw, may be considered as resting on an inclined plane, whose inclination to the horizon is the same as that of a particle of the spiral whereon it rests. Which inclination will be the same as that of AB to AC, if AC be the circumference of a circular section of the cylinder on whose superficies the spiral is described; and BC, a perpendicular on AC, be the distance of two contiguous threads of the spiral measured according to the cylinder's length. Fig. 25.

Moreover, the screw being worked by a horizontal lever whose fulcrum is at the axis thereof, and the whole force of a weight raised or sustained by this instrument being incumbent on the spiral threads, and the distance of every point of those threads from the axis being equal to the radius of a circular section of the cylinder whereon they are described; if that radius be denoted by  $r$ , the weight will, with regard to the lever, act at the distance  $r$  from the fulcrum. Therefore, the screw being loaded with the weight  $W$ , and  $d$  being the distance from the screw's axis at which a power  $P$  acts on the lever, in a horizontal direction at right angles thereto, sustaining  $W$  in equilibrio;  $P$  must, by the last article, be equal to  $\frac{r \times BC}{d \times AC} \times W$ ,  $d$  and  $r$  being written instead of  $EF$  and  $WF$ . But  $c$  being the circumference of a circle whose radius is unity,  $AC$  will here be equal to  $c \times r$ . Therefore  $P$  will be equal to

BC

$\frac{BC}{c \times d} \times W$ , or  $P$  to  $W$  as  $BC$  to  $(c \times d)$  the circumference of a circle whose radius is  $d$ .

We have here supposed the spiral threads of the screw to be cut but a very little depth into the cylinder they are described upon. But it appearing by our demonstration that, if  $d$  and  $BC$  remain the same,  $P$  will have the same advantage with regard to sustaining  $W$ , let  $r$  be what it will; it is plain, that, how deep soever the threads of the screw be cut,  $P$  must be to  $W$  in the ratio above assigned.

27. If we would enquire when there will be an equilibrium between a power pressing against the back of a wedge, and a force acting against its sides in opposition to that power, we must first consider in what direction the wedge is resisted by that opposing force; for the question cannot be answered in general terms.

Some writers suppose the wedge resisted by a force acting at right angles to its sides; others consider the resisting force as acting in a direction parallel to the back of the wedge: whence it is, that their conclusions are different.

Fig. 26.

A wedge  $ABC$ , pressed downwards by the force of a weight  $P$ , being resisted by two forces  $W$  and  $W$  (equal to each other) pressing perpendicularly against its sides, and sustaining  $P$  in equilibrio; those three forces will be as  $AB$ ,  $AC$ , and  $BC$  respectively: and consequently  $P$  will be to  $(2W)$  the sum of the other two forces, or the whole resistance opposing the progress of the wedge, as  $(AB)$  the back of the wedge to  $(AC + BC)$  the sum of its sides, or as half the back of the wedge to one of its sides. For,  $DE$ ,  $FG$  being perpendicular to  $AC$ ,  $BC$  respectively,  
draw

draw  $DF$  parallel to  $GF$ , and  $EF$  parallel to the vertical line  $CP$ , a perpendicular on  $AB$ . Then  $EF$  being put to denote the force  $P$  pressing the wedge directly downwards,  $DE$ ,  $DF$  will (by art. 19.) denote the equal forces pressing against  $AC$  and  $BC$ . But the angles  $ACP$ ,  $CEF$  will be equal, and consequently  $DEF$  equal to  $CAP$ . And,  $DE$ ,  $DF$  being equal, the angles  $DEF$ ,  $DFE$  will be equal to each other, and each to each of the angles  $CAB$ ,  $CBA$ . Therefore the isocles triangle  $DEF$  will be similar to  $CAB$ , and the forces will be in the ratio we have assigned.

In making the experiment, represented by our scheme, the weight of the wedge, it is obvious, must be considered as part of  $P$ . And part of the weight of the sliders  $Q$  and  $R$ , which rest on the inclined planes  $S$  and  $T$ , acting conjunctly with the wedge in opposing the forces  $W$  and  $W$ , urging the said sliders against its sides; an equivalent weight must be added to  $W$ : which equivalent weight must (by art. 23.) be to the weight of the respective slider as half  $EF$  to  $DE$ ; or, which is the same thing, as  $AP$  to  $AC$ .

28. If the forces opposing the progress of the wedge, instead of acting as supposed in the last article, act in directions parallel to its back; the rest being as in that article, the ratio of the power urging the wedge forward to the equipollent resistance may be determined as follows. The two forces pressing the sides of the wedge acting by means of sliders as represented by our scheme, the wedge will press the end  $E$  of the slider  $Q$  in the direction  $ED$ , and  $\frac{ED}{EF} \times P$  will (by the preceding article) denote the force of that pressure. Moreover, (by art. 21.) the efficacy of that pressure

D

pressure

Fig. 27.

pressure in a direction parallel to HD will be expressed by  $\frac{DH}{EF} \times P$ ; being to the said pressure in the direction ED, as DH to DE; DH being drawn parallel to AB, to which QE is now supposed parallel. Now W, the force urging the slider Q against AC, being equal to  $\frac{DH}{EF} \times P$ ; (as it is evident it must be, to resist the pressure of the wedge, so as to sustain it in equilibrio;) the equal force W urging the slider R against the other side of the wedge will likewise be expressed by  $\frac{DH}{EF} \times P$ . Therefore, the power P will be to 2W, the whole resistance opposing the progress of the wedge, as P to  $2 \times \frac{DH}{EF} \times P$ ; that is, as EF to twice DH; or, which is the same thing, as (AP) half the back of the wedge to (CP) its height.

In cleaving of wood, the resistance opposing the force of the mallet (supposing the sides of the wedge perfectly polished, and its edge a line without breadth) is the attraction of cohesion of the particles of the wood which are about to be separated; and this being a kind of pressive force acting against the sides of the wedge, it is extremely absurd to attempt to compare it with the percussive force of a mallet, as some writers have done. For the greatest finite pressive force must give way to the least percussive one, and there cannot possibly be an equilibrium between two such different forces. Any percussive force acting on a moveable body generates a finite quantity of motion in an indefinitely small particle of time; but the time will be finite in which any given pressive force whatever, acting on the same body, can generate or destroy the same quantity of motion. Therefore, a body being urged in a certain

certain direction by any pressive force whatever, and in the contrary direction by any percussive one, the pressive force will be some finite time in destroying the quantity of motion which the percussive one generated in an instant. Consequently how great soever the pressive force may be, and how small soever the percussive one, the body will be moved (at least for some short time) by this last force. Indeed, after the stroke is given, the pressive force may quickly prevail and force back the body which the impulse of the other force had driven forward. And so it would frequently be with respect to cleaving wood, if the sides of the wedge were perfectly smooth. For, after the stroke of the mallet, the wedge, unless its weight were equivalent to the force of the attraction of the parts of the wood about to be separated, would presently be forced back from the place whereto the mallet had driven it. And it is chiefly the roughness of the sides of the wedge, and of the parts of the wood in contact therewith, which, in that operation keeps the wedge from receding. It is that roughness too, and the bluntness of the edge, which sometimes prevent the wedge from being moved by the stroke of the mallet. For, were it not obstructed by such roughness and bluntness, it would, according to what we just now observed, be always driven forward even by the least percussive force.

Having established the doctrine of equilibriums on its own proper principles, and explained the same so far as relates to an equipollence between the power and weight in the use of the instruments commonly called the Mechanic Powers; I shall conclude this Memoir with an article or two, farther explaining the doctrine particularly inculcated above; and moreover exemplifying the appli-



cation of the theorem borrowed from the doctrine of motion, to which recourse is commonly had in treating of equilibriums.

Fig. 28.

29. If the wedge (or inclined plane)  $ABC$ , placed with one of its sides ( $BC$ ) on a polished horizontal plane  $BD$ , be urged in the direction  $BCD$  by a weight  $P$  acting at the angular point  $C$  by means of a string passing over the immoveable pulley  $D$ , whilst the weight  $W$  appending to the string  $ENW$  (having one of its ends fastened to the fixed point  $E$ , and the other to the said weight  $W$ ) presses against the other side ( $AC$ ) of the wedge:  $P$  and  $W$  will rest in equilibrio when,  $MO$  being a vertical line, and  $MN$ ,  $NO$  at right angles to  $AC$ ,  $MW$  respectively,  $P$  is to  $W$  as  $NO$  to  $MW$ , let the length of the string  $EW$  be what it will, and the point  $E$  where it will between the vertical line  $WM$  and the side  $CA$ , or their continuations. For then  $W$  will be kept at rest by its weight, the pressure of the plane  $AC$ , and the tension of the string  $ENW$ ; which (by art. 19.) will be to each other respectively as  $MW$ ,  $MN$ , and  $NW$ . The pressure of  $W$ , in a direction parallel to  $MN$ , will therefore be equal to  $\frac{MN}{MW} \times W$ ; and (by art. 21.) the efficacy of that pressure, in a direction parallel to  $ON$ , will be expressed by  $\frac{NO}{MW} \times W$ ; which must, in the case of an equilibrium, be equal to  $P$ , let the weight of the wedge be what it will. Consequently  $P$  will in that case be to  $W$  as  $\frac{NO}{MW} \times W$  to  $W$ ; that is, as  $NO$  to  $MW$ .

If the string  $ENW$  be parallel to  $AC$ , and  $AF$  be a perpendicular from  $A$  on  $BC$ ;  $NW$  will be equal to  $AF$

$\frac{AF}{AC} \times MW$ , and NO equal to  $\frac{AF \times CF}{AC^2} \times MW$ : and therefore, in this case, NO to MW (that is, P to W) as  $AF \times CF$  to  $AC^2$ . Mr. Ferguson (in his Lectures) erroneously asserts that, in this case, P will be to W as AF to AC; and he endeavours to support his assertion by experiments, but deduces therefrom (by not rightly considering the effect of friction in the different experiments) a conclusion most egregiously absurd, purporting that, in two cases of equilibriums, the powers will be equal, whereas in truth (their ratio being as CF to AC) the power in one case may be ten thousand times greater than in the other case!

When, supposing the power and weight (in any machine) to be moved, the ascent of one and descent of the other (estimated vertically) do not always retain an invariable ratio to each other, it will be wrong, in applying the theorem regarding the ratio of the velocities mentioned at the beginning of this memoir, to take the ratio of the contemporaneous ascent and descent at the end of any certain time whatever for the ratio of the velocities of the power and weight. For instance, in the case of our wedge and weights, whilst P descends in a vertical line, W will describe a circular arc WXY about the center E, and QX drawn from AC to the said arc being parallel to the horizon, and Wg perpendicular thereto, the descent of P will be to the ascent of W as QX to Wg, which being a variable ratio, no use can be made of it in ascertaining the ratio of the velocities of P and W, unless we consider the arc WX as indefinitely small. Now, considering that arc as an indefinitely small particle of the arc WXY, Wg a tangent to that arc at W, CG parallel to that tangent, and

AgG.

$AgG$  parallel to  $BFGD$ ,  $Wg$  will be to  $QX$  as  $AF$  to  $AG$ , that is, (in the case of an equilibrium) as  $P$  to  $W$ . But (the triangles  $ACG$ ,  $MNW$  being similar)  $AF$  is to  $AG$  as  $NO$  to  $MW$ . It is evident therefore that the conclusion thus deduced agrees with that which is deduced above from the proper principles of equilibriums, whether  $EW$  be parallel to  $AC$  or not.

30. If, instead of the weight  $W$  being fastened to the string  $EW$ , it be perforated and moveable upon a string passing through the perforation and fastened at  $H$  and  $K$ , the rest being as before, when the power  $P$  descends the weight  $W$  will describe an ellipsis (whose foci are  $H$  and  $K$ ); to which let  $Wg$  be a tangent at  $W$ , (that is, let  $Wg$  be perpendicular to the line bisecting the angle  $HWK$ ;) and let  $CG$  be parallel to  $Wg$ : then, in the case of an equilibrium,  $P$  will be to  $W$  as  $AF$  to  $AG$ ; as appears by properly applying the theorem before mentioned, regarding the velocities of  $P$  and  $W$  when moved: —and the same conclusion is deducible from the proper principles of equilibriums, the line bisecting the angle  $HWK$  being the direction of the force on the weight  $W$  arising from the tension of the string  $HWK$ , and the effect of that weight on the wedge the same as if it were fastened to the end of a string coinciding with, and having its other end fastened to a fixed point somewhere in that bisecting line.

---

---

## M E M O I R II.

### *Of the Ellipsis and Hyperbola.*

**S**OME of the theorems given by mathematicians for the calculation of fluents by means of elliptic and hyperbolic arcs requiring, in the application thereof, the difference to be taken between an arc of a hyperbola and its tangent; and such difference being not directly attainable when such arc and its tangent both become infinite, as they will do when the *whole* fluent is wanted, although such fluent be at the same time finite; those theorems therefore in that case fail, a computation thereby being then impracticable without some farther help.

The supplying that defect I considered as a point of some importance in geometry, and therefore I earnestly wished and endeavoured to accomplish that business; my aim being to ascertain, by means of such arcs as above mentioned, the *limit* of the difference between the hyperbolic arc and its tangent, whilst the point of contact is supposed to be carried to an infinite distance from the vertex of the curve, seeing that, by the help of that *limit*, the computation would be rendered practicable in the case wherein, without such help, the before-mentioned theorems fail. The result of my endeavours respecting that  
point

point appears in this Memoir : which, amongst other matters, contains the investigation of a general theorem for finding the length of any arc of any conic hyperbola by means of two elliptic arcs. A discovery (first published by me in the *Philos. Transact.* for 1775,) whereby we are enabled to bring out very elegant conclusions in many interesting enquiries, as well mechanical as purely geometrical.

Fig. 29.

1. Suppose the curve ADEF be a conic *hyperbola*, whose semi-transverse axis AC is =  $m$ , and semi-conjugate =  $n$ . Let CP, perpendicular to the tangent DP, be called  $p$ ; and put  $f = \frac{m^2 - n^2}{2m}$ ,  $x = \frac{p^2}{m}$ . Then will DP - AD be

= the fluent of  $\frac{-\frac{1}{2}m^{\frac{1}{2}}x^{\frac{1}{2}}\dot{x}}{\sqrt{n^2 + 2fx - x^2}}$ ,  $p$  and  $x$  being each =  $m$  when AD is = 0. For, denoting the semidiameter CD by  $r$ , and its semi-conjugate diameter by  $s$ , we have (by the nature of the curve)  $r^2 - s^2 = m^2 - n^2 = 2fm$ , and  $ps = mn$ . Whence  $s^2 = r^2 - 2fm = \frac{m^2n^2}{p^2} = \frac{mn^2}{x}$ ; and consequently  $r^2 = \frac{2fmx + mn^2}{x}$ , and  $DP = \sqrt{r^2 - p^2} = \frac{\sqrt{mn^2 + 2fmx - mx^2}}{x^{\frac{1}{2}}}$ . Hence the fluxion of DP is found

$$= -\frac{m^{\frac{1}{2}}n^2\dot{x} + m^{\frac{1}{2}}x^{\frac{1}{2}}\dot{x}}{2x^{\frac{3}{2}}\sqrt{n^2 + 2fx - x^2}}$$

Now it is obvious that the fluxion of the curve AD is to  $\dot{r}$  as  $r$  to  $\sqrt{r^2 - p^2}$ ; therefore the fluxion of AD is =  $\frac{r\dot{r}}{\sqrt{r^2 - p^2}}$ , which by substitution appears to be =

-  $m$

$\frac{-m^{\frac{1}{2}}n^{\frac{1}{2}}z}{2z^{\frac{3}{2}}\sqrt{n^2+2fz-z^2}}$ . Consequently the difference of the fluxions of DP and AD is  $= \frac{-\frac{1}{2}m^{\frac{1}{2}}z^{\frac{1}{2}}z}{\sqrt{n^2+2fz-z^2}}$ .

2. Suppose the curve adefg to be a quadrant of an *ellipsis*, whose semi-transverse axis cg is  $= \sqrt{m^2+n^2}$ , and semi-conjugate ac  $= n$ . Let ct be perpendicular to the tangent dt, and let the abscissa cp be  $= n \times \frac{z}{m}$ . Then will the said tangent dt be  $= m \times \frac{mz-z^2}{n^2+mz}$ ; and the fluxion thereof will be found  $= \frac{1}{2}mn^2z^{-\frac{1}{2}}z \times \frac{m-z}{n^2+mz}$   
 $-\frac{\frac{1}{2}m^{\frac{1}{2}}z^{\frac{1}{2}}z}{\sqrt{n^2+2fz-z^2}}$ .

3. In the expression  $\frac{y^{q-1}j}{(a+by)^r \times (c+dy)^s}$ , let  $\frac{c+dy}{a+by}$  be supposed  $= z$ . Then will  $\frac{ay-c}{d-by}$  be  $= y$ , and the proposed expression will be  $= \frac{(ad-bc)^{1-r-s} \times z^{-s}z}{(az-d)^{1-r} \times (d-bz)^{1+s-r-s}}$ .

4. Taking, in the last article,  $r$  and  $s$  each  $= \frac{1}{2}$ ,  $q = \frac{1}{2}$ ,  $a = -d = \frac{n^2}{m}$ ,  $b = 1$ , and  $c = n^2$ , we have

$$\frac{y^{\frac{1}{2}}j}{\left(\frac{n^2}{m}+y\right)^{\frac{1}{2}} \times \left(n^2-\frac{n^2}{m}y\right)^{\frac{1}{2}}} \left( = \frac{m^{\frac{1}{2}}n^{-1}y^{\frac{1}{2}}j}{\sqrt{n^2+2fy-y^2}} \right) = -mnz^{-\frac{1}{2}}z$$

$\times \frac{m-z}{n^2+mz}$ . It appears therefore, that,  $y$  being  $= n^2 \times$

$$\frac{m-z}{n^2+ mz} - \frac{\frac{1}{2}m^{\frac{1}{2}}y^{\frac{1}{2}}j}{\sqrt{n^2+2fy-y^2}} - \frac{\frac{1}{2}m^{\frac{1}{2}}z^{\frac{1}{2}}z}{\sqrt{n^2+2fz-z^2}} \text{ is } = \frac{1}{2}mn^2z^{-\frac{1}{2}}z$$

$$\times \frac{(m-z)^{\frac{1}{2}}}{(n^2+mz)^{\frac{1}{2}}} - \frac{\frac{1}{2}m^{\frac{1}{2}}z^{\frac{1}{2}}z}{\sqrt{n^2+2fz-z^2}}; \text{ which, by Art. 2. is = the}$$

*fluxion of the tangent dt.*

Consequently, taking the fluents by Art. 1, and correcting them properly, we find

$$DP - AD + FR - AF = L + dt.$$

CP, in fig. 29, being  $= m^{\frac{1}{2}}z^{\frac{1}{2}}$ ; cp, in fig. 30,  $= n \times \left(\frac{z}{m}\right)^{\frac{1}{2}}$ ;

CR, perpendicular to the tangent FR,  $= m^{\frac{1}{2}}y^{\frac{1}{2}}$ ;

$$DP - AD = \text{the fluent of } \frac{-\frac{1}{2}m^{\frac{1}{2}}z^{\frac{1}{2}}z}{\sqrt{n^2+2fz-z^2}};$$

$$FR - AF = \text{the fluent of } \frac{-\frac{1}{2}m^{\frac{1}{2}}y^{\frac{1}{2}}j}{\sqrt{n^2+2fy-y^2}};$$

and L the *limit* to which the difference DP - AD, or FR - AF, approaches upon carrying the point D, or F, from the vertex A *ad infinitum*.

5. Suppose  $y$  equal to  $z$ , and that the points D and F then coincide in E, the points d and p being at the same time in e and q respectively. Then cv being perpendicular to the tangent ev, that tangent will be a *maximum* and equal to  $cg - ac = \sqrt{m^2 + n^2} - n$ ; the tangent EQ (in the hyperbola) will be  $= \sqrt{m^2 + n^2}$ ; the abscissa BC  $= m \times \sqrt{1 + \frac{n}{\sqrt{m^2 + n^2}}}$ ; the ordinate BE  $= n \times \sqrt{\frac{n}{\sqrt{m^2 + n^2}}}$ ; and it appears, that L is  $= 2EQ - 2AE - cv = n + \sqrt{m^2 + n^2} - 2AE$ . Thus the *limit* which I proposed to ascertain is investigated,  $m$  and  $n$  being any right lines what-

whatever. Another expression for such limit will be found in a subsequent article in this Memoir.

6. The whole fluent of  $\frac{\frac{1}{2}m^{\frac{1}{2}}z^{\frac{1}{2}}z}{\sqrt{n^2 + 2fz - z^2}}$ , generated whilst  $z$  from 0 becomes  $= m$ , being equal to  $L$ ; and the fluent of the same fluxion (supposing it to begin when  $z$  begins) being in general equal to  $L + AD - DP = FR - AF - dt$ ; it appears that  $k$  being the value of  $z$  corresponding to the fluent  $L + AD - DP$ ,  $\frac{mn^2 - n^2k}{n^2 + mk}$  will be the value of  $z$  corresponding to the fluent  $L + AF - FR$ , and  $FR - AF$  will be the part generated whilst  $z$  from  $\frac{mn^2 - n^2k}{n^2 + mk}$  becomes  $= m$ . It follows therefore, that the tangent  $dt$ , together with the fluent of  $\frac{\frac{1}{2}m^{\frac{1}{2}}z^{\frac{1}{2}}z}{\sqrt{n^2 + 2fz - z^2}}$  generated whilst  $z$  from 0 becomes equal to any quantity  $k$ , is equal to the fluent of the same fluxion generated whilst  $z$  from  $\frac{mn^2 - n^2k}{n^2 + mk}$  becomes  $= m$ ;  $cp$  being taken  $= n \times \left(\frac{k}{m}\right)^{\frac{1}{2}}$ .

Suppose  $k = \frac{mn^2 - n^2k}{n^2 + mk}$ ; its value will then be  $\frac{n}{m}\sqrt{m^2 + n^2} - \frac{n^2}{m}$ . Consequently the fluent of  $\frac{\frac{1}{2}m^{\frac{1}{2}}z^{\frac{1}{2}}z}{\sqrt{n^2 + 2fz - z^2}}$  generated whilst  $z$  from 0 becomes  $= \frac{n}{m}\sqrt{m^2 + n^2} - \frac{n^2}{m}$ , together with the quantity  $\sqrt{m^2 + n^2} - n$ , is equal to the fluent of the same fluxion generated whilst  $z$  from  $\frac{n}{m}\sqrt{m^2 + n^2} - \frac{n^2}{m}$  becomes  $= m$ : and these two parts of the whole fluent

E 2 being



being denoted by M and N respectively; M will be =  $n - AE$ , and  $N = \sqrt{m^2 + n^2} - AE$ .

7. The fluent of  $\frac{\frac{1}{2}m^{\frac{1}{2}}z^{\frac{1}{2}}z}{\sqrt{n^2 + 2fz - z^2}}$  being =  $L + AD - DP$ ,  
 the fluent of  $\frac{\frac{1}{2}m^{\frac{1}{2}}z^{\frac{1}{2}}z}{\sqrt{n^2 + 2fz - z^2}} + DP - AD - L$  will be = 0.  
 Therefore, the fluent of  $\frac{\frac{1}{2}m^{\frac{1}{2}}z^{\frac{1}{2}}z}{\sqrt{n^2 + 2fz - z^2}}$  + the fluent of  
 $\frac{\frac{1}{2}m^{-\frac{1}{2}}n^2z^{-\frac{1}{2}}z}{\sqrt{n^2 + 2fz - z^2}}$  being = the fluent of  $\frac{1}{2}z^{-\frac{1}{2}}z \times \left(\frac{n^2 + mz}{m - z}\right)^{\frac{1}{2}}$ ,  
 it is obvious, that the fluent of  $\frac{\frac{1}{2}m^{-\frac{1}{2}}n^2z^{-\frac{1}{2}}z}{\sqrt{n^2 + 2fz - z^2}}$  is =  $DP -$   
 $AD - L$  + the fluent of  $\frac{1}{2}z^{-\frac{1}{2}}z \times \left(\frac{n^2 + mz}{m - z}\right)^{\frac{1}{2}}$  =  $DP - AD - L$   
 + the *elliptic arc* dg (Fig. 30.) whose abscissa cp is  
 =  $n \times \left(\frac{z}{m}\right)^{\frac{1}{2}}$ .

Consequently, putting E for  $\frac{1}{4}$  of the periphery of that  
 ellipsis, it appears that the *whole fluent* of  $\frac{\frac{1}{2}m^{-\frac{1}{2}}n^2z^{-\frac{1}{2}}z}{\sqrt{n^2 + 2fz - z^2}}$ ,  
 generated whilst z from 0 becomes = m, is equal to  $E - L$   
 =  $E + 2AE - n - \sqrt{m^2 + n^2}$ .

8. By taking, in Art. 3. q, r, and s, each =  $\frac{1}{2}$ ; and  
 $a = -d = \frac{n^2}{m}$ ,  $b = 1$ , and  $c = n^2$ ; we find, that, if y be  
 =  $\frac{mn^2 - n^2z}{n^2 + mz}$ ,  $\frac{z^{-\frac{1}{2}}z}{\sqrt{n^2 + 2fz - z^2}} + \frac{y^{-\frac{1}{2}}y}{\sqrt{n^2 + 2fy - y^2}}$  will be = 0.

It is obvious therefore, that the fluent of  $\frac{z^{-\frac{1}{2}}z}{\sqrt{n^2 + 2fz - z^2}}$ ,  
 gene-

generated whilst  $z$  from  $o$  becomes equal to any quantity  $k$ , is equal to the fluent of the same fluxion, generated whilst  $z$  from  $\frac{mn^2 - n^2k}{n^2 + mk}$  becomes  $= m$ .

Now, supposing  $k = \frac{mn^2 - n^2k}{n^2 + mk}$ , its value will be  $\frac{n}{m} \sqrt{m^2 + n^2} - \frac{n^2}{m}$ . Consequently the fluent of  $\frac{z^{-\frac{1}{2}} dz}{\sqrt{n^2 + 2fz - z^2}}$ , generated whilst  $z$  from  $o$  becomes  $= \frac{n}{m} \sqrt{m^2 + n^2} - \frac{n^2}{m}$ , is equal to *half* the fluent of the same fluxion, generated whilst  $z$  from  $o$  becomes  $= m$ ; which *half fluent* is known by the preceding article.

9. It appears by Art. 4. that  $\frac{\frac{1}{2} m^{\frac{1}{2}} y^{\frac{1}{2}} dy}{\sqrt{n^2 + 2fy - y^2}} + \frac{\frac{1}{2} m^{\frac{1}{2}} z^{\frac{1}{2}} dz}{\sqrt{n^2 + 2fz - z^2}}$  is  $= -$  the fluxion of the tangent  $dt$ ; and it appears by the last article, that  $\frac{\frac{1}{2} m^{-\frac{1}{2}} n^2 y^{-\frac{1}{2}} dy}{\sqrt{n^2 + 2fy - y^2}} + \frac{\frac{1}{2} m^{-\frac{1}{2}} n^2 z^{-\frac{1}{2}} dz}{\sqrt{n^2 + 2fz - z^2}}$  is  $= 0$ ;  $mn^2 - n^2y - n^2z - myz$  being  $= 0$ . Therefore, by addition, we have  $\frac{1}{2} y^{-\frac{1}{2}} dy \times \frac{n^2 + my}{m - y}^{\frac{1}{2}} + \frac{1}{2} z^{-\frac{1}{2}} dz \times \frac{n^2 + mz}{m - z}^{\frac{1}{2}}$   $= -$  the fluxion of the tangent  $dt$ . Consequently, by taking the correct fluents, we find the tangent  $dt$  ( $=$  the tangent  $fw$ )  $=$  the arc  $ad$   $-$  the arc  $fg$ ; the abscissa  $cp$  being  $= n \times \frac{z^{\frac{1}{2}}}{m}$ , the abscissa  $cr = n \times \frac{y^{\frac{1}{2}}}{m}$ , and their relation expressed by the equation  $n^6 - n^4 u^2 - n^4 v^2 - m^2 u^2 v^2 = 0$ ,  $u$  and  $v$  being put for  $n \times \frac{z^{\frac{1}{2}}}{m}$  and  $n \times \frac{y^{\frac{1}{2}}}{m}$  respectively. Moreover the tangents  $dt$ ,  $fw$ , will each be  $= \frac{m^2 uv}{u^2}$ , and  $ct \times cw = cv^2 = ac \times cg$ .

If for the semi-transverse axis  $cg$  we substitute  $h$  instead of  $\sqrt{m^2 + n^2}$ , the relation of  $u$  to  $v$  will be expressed by the equation  $n^6 - n^4 u^2 - n^4 v^2 - h^2 - n^2 \times u^2 v^2 = 0$ , and  $dt (= fw)$  will be  $= \frac{b^2 - n^2}{n^3} \times uv$ .

If  $u$  and  $v$  be respectively put for  $fr$  and  $dp$ , their relation will be expressed by the equation  $h^6 - h^4 u^2 - h^4 v^2 + h^2 - n^2 \times u^2 v^2 = 0$ , and  $dt (= fw)$  will be  $= \frac{b^2 - n^2}{b^3} \times uv$ .

10. Suppose  $y$  equal to  $x$ ; (that is,  $v = u$ ;) and that the points  $d$  and  $f$  coincide in  $e$ . In which case the tangent  $dt$  will be a *maximum*, and  $= cg - ac$ . It appears then that the *arc*  $ae$  — the *arc*  $eg$  is  $= cg - ac$ . Consequently, putting  $E$  for the quadrantal arc  $ag$ , we find that

$$\text{the arc } ae \text{ is } = \frac{E + b - n}{2}$$

$$\text{the arc } eg \text{ } = \frac{E - b + n}{2}$$

There are, I am aware, some other parts of the arc  $ag$  whose lengths may be assigned by means of the whole length ( $ag$ ) with right lines; but to investigate such other parts is not to my present purpose.

11. Taking  $m$  and  $n$  each  $= 1$ ; that is,  $ac (= AC) = 1$ , and  $cg = \sqrt{2}$ ; let the arc  $ag$  be then expressed by  $e$ : put  $c$  for *one fourth* of the periphery of the circle whose radius is  $1$ ; and let the *whole fluents* of  $\frac{\frac{1}{2}z^{\frac{1}{2}}z}{\sqrt{1-z^2}}$  and  $\frac{\frac{1}{2}z^{-\frac{1}{2}}z}{\sqrt{1-z^2}}$ , generated whilst  $z$  from  $0$  becomes  $= 1$ , be denoted by  $F$  and  $G$  respectively. Then, by what is said above,  $F + G$  will be  $= e$ ; and, by Part X. of my *Math. Lucubrat.* it appears

pears that  $F \times G$  is  $= \frac{1}{2}c$ . From which equations we find  $F = \frac{1}{2}e - \frac{1}{2}\sqrt{e^2 - 2c}$ , and  $G = \frac{1}{2}e + \frac{1}{2}\sqrt{e^2 - 2c}$ .

But  $m$  and  $n$  being each  $= 1$ ,  $L$  is  $= F$ ; therefore  $1 + \sqrt{2} - 2AE$ , the value of  $L$  from Art. 5, is, in this case,  $= \frac{1}{2}e - \frac{1}{2}\sqrt{e^2 - 2c}$ . Consequently, in the equilateral hyperbola, the arc  $AE$ , whose abscissa  $BC$  is  $= \sqrt{1 + \frac{1}{\sqrt{2}}}$ , will be  $= \frac{1}{2} + \frac{1}{\sqrt{2}} - \frac{1}{4}e + \frac{1}{4}\sqrt{e^2 - 2c}$ , by what is said in the article last mentioned. Hence the rectification of that arc may be effected by means of the circle and ellipsis!

12. By substituting  $a - b$ ,  $2 \times \overline{ab}^{\frac{1}{2}}$ , and  $\overline{a - b}^2 - t^2$ , Fig. 29. for  $m$ ,  $n$ , and  $mz$  respectively, in Art. 1. it appears, that, if (in the hyperbola) the semi-transverse axis  $AC$  be  $= a - b$ , the semi-conjugate  $= 2 \times \overline{ab}^{\frac{1}{2}}$ , and the perpendicular  $CP = \overline{a - b}^2 - t^2$ ; the difference  $(DP - AD)$  between the tangent  $DP$  and the arc  $AD$  will be equal to the fluent of  $\frac{\overline{a - b}^2 - t^2}{\overline{a + b}^2 - t^2} \times t$ .

13. It is well known, that, in any ellipsis whose semi-transverse axis is  $h$ , and semi-conjugate  $n$ , if  $x$  be the abscissa, measured from the center upon the transverse axis, and  $Q$  the arc between the conjugate axis and the ordinate corresponding to  $x$ ,  $\frac{\overline{b^2 - gx^2}}{\overline{b^2 - x^2}} \times x$  will be  $= Q$ ,  $g$  being  $= \frac{b^2 - n^2}{b^2}$ .

Hence, by substituting  $a + b$ ,  $2 \times \overline{ab}^{\frac{1}{2}}$ , and  $\frac{a + b}{a - b} \times t$ , for

Fig. 31. for  $h$ ,  $n$ , and  $x$  respectively; it appears, that, in the ellipsis  $aed$  whose semi-transverse axis  $cd$  is  $= a + b$ , semi-conjugate  $ca = 2 \times \overline{ab}^{\frac{1}{2}}$ , and abscissa  $cb$  (corresponding to the ordinate  $be$ )  $= \frac{a+b}{a-b} \times t$ , the arc  $ae$  (denoted by  $Q$ ) will be equal to the fluent of  $\frac{\overline{a+b}^2 - t^2}{a-b}^{\frac{1}{2}} \times \dot{t}$ .

Fig. 32. 14. In the ellipsis  $aed$ , the semi-transverse axis  $cd$  being  $= a$ , the semi-conjugate  $ca = b$ , and the abscissa  $cb$  (corresponding to the ordinate  $be$ )  $= x$ ; if  $ep$ , the tangent at  $e$ , intercepted by a perpendicular ( $cp$ ) drawn thereto from the center  $c$ , be denoted by  $t$ ;  $gx \times \frac{a^2 - x^2}{a^2 - gx^2}^{\frac{1}{2}}$  (as is well known) will be  $= t$ ,  $g$  being  $= \frac{a^2 - b^2}{a^2}$ .

Whence we have  $x^2 = \frac{a^2 g + t^2}{2g} - \frac{\overline{a^2 - b^2}^2 - 2 \times a^2 + b^2 \times t^2 + t^4}{2g}$ .

From which equation, by taking the fluxions, we have

$$\begin{aligned} x \dot{x} &= \frac{t \dot{t}}{2g} + \frac{\overline{a^2 + b^2} \times t \dot{t} - t^3 \dot{t}}{2g \sqrt{a^2 - b^2}^2 - 2 \times a^2 + b^2 \times t^2 + t^4} \\ &= \frac{t \dot{t}}{2g} + \frac{\overline{a^2 + b^2} \times t \dot{t} - t^3 \dot{t}}{2g \sqrt{a-b}^2 - t^2 \times \overline{a+b}^2 - t^2}. \end{aligned}$$

But  $\dot{R}$ , the fluxion of the arc  $ae$ , being  $= \frac{a^2 - gx^2}{a^2 - x^2}^{\frac{1}{2}} \times \dot{x}$ , according to the preceding article, it follows that  $\frac{gx \dot{x}}{t}$  is  $= \dot{R}$ . It is obvious therefore that

$\dot{R}$  is

$$\begin{aligned}
 R \text{ is } &= \frac{i}{2} + \frac{\overline{a^2 + b^2} \times ti - t^2 i}{2\sqrt{\overline{a-b}^2 - t^2} \times \overline{a+b}^2 - t^2} \\
 &= \frac{i}{2} + \frac{\overline{a-b}^2 - t^2 \times i}{4\sqrt{\overline{a-b}^2 - t^2} \times \overline{a+b}^2 - t^2} + \frac{\overline{a+b}^2 - t^2 \times i}{4\sqrt{\overline{a-b}^2 - t^2} \times \overline{a+b}^2 - t^2} \\
 &= \frac{1}{2}t + \frac{1}{4} \times \frac{\overline{a-b}^2 - t^2}{\overline{a+b}^2 - t^2} \times i + \frac{1}{4} \times \frac{\overline{a+b}^2 - t^2}{\overline{a-b}^2 - t^2} \times i.
 \end{aligned}$$

From whence, by taking the fluents, according to Art. 12. and 13. we find  $R = ae$  (Fig. 32.)  $= \frac{1}{2}t + \frac{DP-AD}{4}$  (Fig. 29.)  $+ \frac{Q}{4}$ . Consequently the hyperbolic arc AD is  $= DP + Q - 4R + 2t$ .

Thus, beyond my expectation, I find, that the *hyperbola* may in general be rectified by means of two *ellipses*!

If  $p'e'$ ,  $p''e''$  be equal tangents to the ellipsis  $ae'e''d$ ; the arc  $ae'$  (denoted by R) will (by Art. 9.) be equal to the arc  $de''$  + the tangent  $p''e''$ , or  $p'e'$  (denoted by  $t$ ). Therefore, substituting for  $t$  its value found by this last equation, it appears that AD is  $= DP + Q - 2R - 2 \times de'' = DP + Q + 2 \times e'e'' - 2E''$ , and  $DP - AD = 2E'' - 2 \times e'e'' - Q$ ;  $E''$  being put for the quadrantal arc  $ad$ . It is observable that, when  $t$  is  $= a - b$ ,  $e''$  (by Art. 5.) coincides with  $e'$  and  $e'e''$  is  $= 0$ ;  $Q$  (by Art. 13.) = the quadrantal arc  $ad$ ; and (by Art. 12) DP and AD both become infinite. Consequently writing E for that quadrantal arc ( $ad$ ), and L for the *limit* of the difference  $DP - AD$  whilst the point of contact (D) is supposed to be carried to an infinite distance from the vertex A of the hyperbola, we find  $2E'' - E = L$ .

15. Exterminating  $a$ ,  $b$ , and  $t$ , by means of the equations  $a - b = m$ ,  $2 \times \overline{ab}^{\frac{1}{2}} = n$ , and  $\overline{a - b}^2 - t^2 = mx$ ; and writing  $f$  for  $\frac{m^2 - n^2}{2m}$  as in Art. 1. it appears that

(V) the fluent of  $\frac{\frac{1}{2}m^{\frac{1}{2}}z^{\frac{3}{2}}}{\sqrt{n^2 + 2fx - z^2}}$  is  $= 2 \times e'e'' - de$  generated whilst  $z$  increases from 0. Moreover the fluent of

$\frac{1}{2}m^{\frac{1}{2}}z^{-\frac{1}{2}} \times \frac{\overline{\frac{n^2}{m} + z}}{m - z}^{\frac{1}{2}}$ , or of its equal  $\frac{\frac{1}{2}m^{-\frac{1}{2}}n^2z^{-\frac{1}{2}} + \frac{1}{2}m^{\frac{1}{2}}z^{\frac{3}{2}}}{\sqrt{n^2 + 2fx - z^2}}$ ,

being  $= de$ , as observed in Art. 7. we find by subtraction

(W) the fluent of  $\frac{\frac{1}{2}m^{-\frac{1}{2}}n^2z^{-\frac{1}{2}}}{\sqrt{n^2 + 2fx - z^2}} = 2 \times \overline{de - e'e''}$ .

The semi-transverse axis of the ellipsis  $aed$  being now  $= \sqrt{m^2 + n^2}$ , the semi-conjugate  $= n$ ; the abscissa  $cb = \frac{m^2 + n^2}{m} \times \overline{m - z}^{\frac{1}{2}}$ , and the ordinate  $be = n \times \frac{z}{m}^{\frac{1}{2}}$ .

The semi-transverse axis ( $a$ ) of the ellipsis  $ae'e'd = \frac{\sqrt{m^2 + n^2}}{2} + \frac{m}{2}$ ; the semi-conjugate ( $b$ )  $= \frac{\sqrt{m^2 + n^2}}{2} - \frac{m}{2}$ ; the tangents  $e'p'$ ,  $e''p''$ , intercepted by perpendiculars ( $cp'$ ,  $cp''$ ) drawn thereto from the center  $c$ , each  $= m^{\frac{1}{2}} \times \overline{m - z}^{\frac{1}{2}}$ ; and the abscissa ( $cb'$ , or  $cb''$ ) on  $cd$ , corresponding to the point  $e'$ , or  $e''$ , of the curve, is determined by the expression

$$\frac{\sqrt{\overline{m^2 + n^2}^{\frac{1}{2}} + m - z} \mp \sqrt{z^2 + \frac{n^2}{m}z}}{2^{\frac{1}{2}} \times \overline{m^2 + n^2}^{\frac{1}{2}}} \times a.$$

From what is done above, many other new theorems for the calculation of fluents are deducible: the most remarkable

markable of which, I intend to insert in the Appendix to be annexed to these Memoirs.

16. Taking  $m$  and  $n$  each equal to 1, we find by the above theorem (marked W) the whole fluent of  $\frac{\frac{1}{2}z - \frac{1}{2}z}{\sqrt{1-z^2}}$   $= 2 \times \overline{e-d}$ ;  $e$  denoting one fourth of the periphery of an ellipsis whose semi-transverse and semi-conjugate axes are  $\sqrt{2}$  and 1; and  $d$  one fourth of the periphery of another ellipsis whose semi-transverse and semi-conjugate axes are  $\frac{1}{\sqrt{2}} + \frac{1}{2}$  and  $\frac{1}{\sqrt{2}} - \frac{1}{2}$ . But, by Art. 11. the same fluent is found  $= \frac{e}{2} + \frac{1}{2}\sqrt{e^2 - 2c}$ ,  $c$  denoting  $\frac{1}{4}$  of the periphery of the circle whose radius is 1. Consequently  $2 \times \overline{e-d}$  being  $= \frac{e}{2} + \frac{\sqrt{e^2 - 2c}}{2}$ , we have, from that equation,  $d = \frac{1}{2}e - \frac{1}{4}\sqrt{e^2 - 2c}$ ; or  $e = \frac{1}{2}d + \frac{1}{2}\sqrt{d^2 - 2c}$ . Thus it appears that the periphery of one certain ellipsis may be found by means of the periphery of another ellipsis and the periphery of a circle!

Before Mr. MACLAURIN published his excellent Treatise of Fluxions, some very eminent mathematicians imagined, that the *elastic curve* could not be constructed by the quadrature or rectification of the conic sections. But that gentleman has shewn, in that treatise, that the said curve may, in every case, be constructed by the rectification of the hyperbola and ellipsis; and he has observed, that, by the same means, we may construct the curve along which, if a heavy body descended, it would recede equally in equal times from a given point. Which last mentioned



curve Mr. JAMES BERNOULLI constructed by the rectification of the elastic curve, and Mr. LEIBNITZ and Mr. JOHN BERNOULLI by the rectification of a geometrical curve of a higher kind than the conic sections. It is observable, that Mr. MACLAURIN's method of construction, just now adverted to, though very elegant, is not without a defect. The difference between the hyperbolic arc and its tangent being necessary to be taken, the method (for the reason mentioned at the beginning of this Memoir) always fails when some principal point in the figure is to be determined; the said arc and its tangent then both becoming infinite, though their difference be at the same time finite. The contents of this Memoir, properly applied, will evince, that both the *elastic curve* and the *curve of equable recess from a given point* (with many others) may be constructed by the rectification of the ellipsis only, without failure in any point.

---



---

M E M O I R III.

*Of the Descent of a Body in a Circular Arc.*

1. **L**ET  $lpqn$  be a semi-circle perpendicular to the horizon, whose highest point is  $l$ , lowest  $n$ , and center  $m$ . Let  $ps$ ,  $qt$ , parallel to the horizon, meet the diameter  $lmn$  in  $s$  and  $t$ ; and let the radius  $lm$  (or  $mn$ ) be denoted by  $r$ ; the height  $ns$  by  $d$ ; and the distance  $st$  by  $x$ . Then, putting  $h$  for  $(16\frac{1}{r^2}$  feet) the space a heavy body, descending freely from rest, falls through in one second of time; and supposing a pendulum, or other heavy body, descending by its gravity from  $p$ , along the arc  $pqn$ , to have arrived at  $q$ ; the fluxion of the time of

Fig. 33.

descent will (as is well known) be  $= \frac{\frac{1}{2}rb^{-\frac{1}{2}}x^{-\frac{1}{2}}}{\sqrt{2dr-d^2-2r-dx-x^2}}$

The fluent whereof, or the time of descent from  $p$  to  $q$ , is, by Art. 15. of the preceding Memoir,  $= \frac{2r}{h^{\frac{1}{2}} \times 2r-d}$

$\times \overline{de - e'e''} : m$ , in the theorem referred to in that Memoir, being taken  $= d^{\frac{1}{2}}$ ,  $n = \sqrt{2r-d}$ ,  $z = \frac{x}{d^{\frac{1}{2}}}$ ; and accordingly

the axes  $\sqrt{m^2 + n^2}$  and  $n$  equal to  $\sqrt{2r}$  and  $\sqrt{2r-d}$ , and

the axes  $\frac{\sqrt{m^2 + n^2}}{2} + \frac{m}{2}$  and  $\frac{\sqrt{m^2 + n^2}}{2} - \frac{m}{2}$  equal to  $\frac{r}{2} + \frac{d^{\frac{1}{2}}}{2}$

and

and  $\left(\frac{r}{2}\right)^{\frac{1}{2}} - \frac{d}{2}$ ;  $cb$  (Fig. 31.) =  $\frac{2r}{d} \times \overline{d-x}^{\frac{1}{2}}$ , and  $e'p', e''p''$ , (Fig. 32.) each =  $\overline{d-x}^{\frac{1}{2}}$ .

Hence it appears, that the *whole* time of descent from  $p$  to  $n$  is =  $\frac{2r}{b^{\frac{1}{2}} \times 2r-d} \times \overline{E-E''}$ ; when, in Fig. 31. and 32. the semi-axes are taken according to the values of  $m$  and  $n$ , just now specified:  $E$  and  $E''$  denoting the quadrantal arcs  $aed$ ,  $aed$ , respectively.

2. If  $pqn$  be a quadrant; that is, if  $d$  be =  $r$ ; the *whole* time of descent from  $p$  to  $n$  will be =  $\frac{2}{b^{\frac{1}{2}}} \times \overline{E-E''}$ , by what is found in the preceding article. Which time will also be =  $\frac{1}{2b^{\frac{1}{2}}} \times \overline{E + \sqrt{E^2 - 2c}}$  (=  $1.31102877 \times \left(\frac{r}{b}\right)^{\frac{1}{2}}$ ),  $c$  being  $\frac{1}{4}$  of the periphery of the circle whose radius is  $r$ ; as appears by writing  $rx$  for  $x$  in  $\left(\frac{\frac{1}{2}rb^{-\frac{1}{2}}x^{-\frac{1}{2}}}{\sqrt{r^2-x^2}}\right)$  the fluxion of the time of descent, and referring to Art. 11. of the preceding Memoir for the *whole* fluent of  $\left(\frac{\frac{1}{2}r^{\frac{1}{2}}x^{-\frac{1}{2}}}{b^{\frac{1}{2}}\sqrt{1-z^2}}\right)$  the resulting expression:  $E$  here denoting the quadrantal arc of the ellipsis, whose semi-transverse and semi-conjugate axes are  $\overline{2r}^{\frac{1}{2}}$  and  $r^{\frac{1}{2}}$ ; and  $E''$  the quadrantal arc of another ellipsis, whose semi-transverse and semi-conjugate axes are  $\left(\frac{r}{2}\right)^{\frac{1}{2}} + \frac{1}{2}r^{\frac{1}{2}}$  and  $\left(\frac{r}{2}\right)^{\frac{1}{2}} - \frac{1}{2}r^{\frac{1}{2}}$ .

3. The semi-axes of the ellipsis (Fig. 31.), of which  $E$  denotes the quadrantal arc, being  $\sqrt{2r}$  and  $\sqrt{2r-d}$ ;  
and

and the semi-axes of the ellipsis (Fig. 32.) of which  $E''$  denotes the quadrantal arc, being  $\frac{r}{2} + \frac{d}{2}$  and  $\frac{r}{2} - \frac{d}{2}$ : those ellipses become circles when  $d$  is = 0, and  $E$  and  $E''$  are then quadrantal arcs of the circles whose radii are  $\sqrt{2r}$  and  $\frac{r}{2}$  respectively. Therefore,  $c$  being  $\frac{1}{4}$  the periphery of the circle whose radius is  $r$ , the last mentioned quadrantal arcs are equal to  $2 \times \frac{c}{\sqrt{2r}}$  and  $\frac{c}{\sqrt{2r}}$  respectively. Consequently, it appears by substitution, that the *limit* of  $\left(\frac{2r}{b^{\frac{1}{2}} \times 2r - d} \times E - E''\right)$  the *whole* time of descent from  $p$  to  $n$ , taking  $d$  less and less, is =  $\frac{c}{\sqrt{2br}}$ : which may be considered as the time of descent in a very small arc; and  $\frac{2^{\frac{1}{2}}c}{\sqrt{br}}$ , or its equal  $\frac{2r}{b} \times q$ , may accordingly be considered as the time of vibration of a pendulum describing such small arc in its descent and a similar one in its ascent,  $q$  being the quadrantal arc of the circle whose radius is 1.

4. Having joined  $lp$ ,  $pt$ , make the angle  $lpv$  equal to the angle  $ltp$ , and draw  $rv$  parallel to the horizon, intersecting the circle in  $r$ , and the diameter  $lmn$  in  $v$ . Then the pendulum or other heavy body, descending by its gravity from  $p$  along the arc  $pqrn$ , will pass over the arcs  $pq$  and  $rn$  exactly in equal times: and therefore,  $qt$  and  $rv$  coinciding when  $lt$  is equal to  $lp$ , it follows, that the time of descent from  $p$  to  $q$  will then be precisely equal to *half* the time of descent from  $p$  to  $n$ .

For-

For  $m$ ,  $n$ , and  $x$  being as specified in Art. 1.  $\frac{mn^2 - n^2x}{n^2 + mx}$   
 is  $= \frac{2dr - d^2 - 2rx + dx}{2r - d + x}$ , which by our construction is  $= sv$ .

Therefore, by Art. 8. of the preceding Memoir, the parts of the fluent of the fluxion in Art. 1. (of this Memoir) corresponding to the times of descent from  $p$  to  $q$  and from  $r$  to  $n$  are equal.

5. If  $pn$  be an arc of  $120^\circ$ ,  $q$  and  $r$  will coincide in the point of the arc  $90^\circ$  above  $n$ ; and, by the preceding article, the descending body will be just as long in passing over the quadrantal arc between  $q$  and  $n$  as in passing over the first  $30^\circ$  between  $p$  and  $q$ . Moreover,  $t$  and  $v$  then coinciding in the center  $m$ , it is obvious, that the vertical descent in the *first half* of the time (of the whole descent) will be equal to *half* the vertical descent in the *other half* of the time.

6. In general, the vertical descent in the *first half* of the time will be to the vertical descent in the *other half* of the time, as  $\sqrt{2r - d}$  to  $\sqrt{2r}$ . It is therefore manifest, that, in the circle, the vertical descents corresponding to those two equal parts of the time (of the whole descent) cannot in any case be equal, as they always are in the cycloid.

---



---

M E M O I R IV.

*Of the centrifugal Force of the Particles of a Body, arising from its rotation about a certain Axis passing through its center of Gravity.*

1. **L**ET  $p$  be a particle of matter firmly connected Fig. 34.  
with the plane  $DOEFQG$ , in which the line  $OCQ$  is situated; and,  $pq$  being a perpendicular from  $p$  to the said plane, let the distance  $pq$  be denoted by  $u$ ; also, the line  $ql$  being at right angles to  $O/CQ$ , let the distance  $pl$  be denoted by  $h$ . Then, the said plane with the particle  $p$  being made to revolve about  $O/CQ$  as an axis, with the angular velocity  $e$ , measured at the distance  $a$  from the said axis, the velocity of  $p$  will be  $= \frac{he}{a}$ , and its centrifugal force from  $l$  will (by a well known theorem) be  $= \frac{he^2}{a^2} \times p$ . Whence, by resolving that force into two others, one in the direction  $qp$ , and the other in a direction parallel to  $lq$ , it appears that, in consequence of the said centrifugal force of  $p$ , the point  $l$  of the plane  $DOEFQG$  will be urged in a direction at right angles to that plane by a force  $= \frac{ue^2}{a^2} \times p$ , let the distance  $lq$  be what it will.

G

2. The

2. The particle  $p$  being connected with the plane  $DOEFQG$  as mentioned in the preceding article, and the distance  $Cl$  being denoted by  $v$ ; if  $p$  be urged directly from the said plane by a force  $= fu \times p$ , the efficacy of that force to turn the said plane about the line  $HCI$ , therein drawn at right angles to  $OCQ$ , will (by the property of the lever) be equivalent to the force  $\frac{fuv \times p}{g}$  acting on the said line  $OCQ$ , at right angles to the said plane, at the distance  $g$  from the point  $C$ .

Moreover it is obvious, that, *cæteris paribus*, the efficacy will be the same, let the distance of  $q$  from  $l$  be what it will.

Fig. 35.

Let  $q$  coincide with  $l$ ; and let  $Ck$  be a line in the plane  $Clp$  continued, (which plane will be at right angles to the plane  $DOEFQG$ ); also,  $pk$  being at right angles to  $Ck$ , let those lines  $pk$  and  $Ck$  be denoted by  $w$  and  $x - k$  respectively,  $k$  and  $x$  respectively denoting the distances of the points  $C$  and  $k$  from some given point  $V$  in the line passing through those two points. Then, the sine and cosine of the angle  $kCO$  (to the radius 1) being respectively denoted by  $m$  and  $n$ , the force  $\frac{fuv \times p}{g}$  will be

$$= \frac{f \times p}{g} \times mn \times \overline{w^2 - x - k^2} - \overline{m^2 - n^2} \times \overline{w \cdot x - k}.$$

Consequently, if each particle of any solid body, through which a line  $HCI$  and a plane  $DOEIFQGH$  may be conceived to pass, be urged from that plane by a force expressed by  $fu \times p$  as above; the force which, acting on the line  $OCQ$  at the distance  $g$  from  $C$ , would be equivalent to the efficacy of all the forces acting on the several particles of that body, to turn the same about the  
line

line HCI, will be obtained by computing the sum of all the forces  $\frac{f \times p}{g} \times mn \times \overline{w^2 - x - k^2 - m^2 - n^2} \times \overline{w \cdot x - k}$  acting on the said body.

The computation of such equivalent force will, in many cases, be abridged by observing, that, if  $pk$  be continued to  $p''$  so that  $kp''$  be  $= kp$ , the efficacy of the force on the particle  $p''$ , to turn the body about the line HCI in opposition to the force on the particle  $p$ , will be represented by the equivalent force  $\frac{f \times p''}{g} \times mn \times \overline{x - k^2 - w^2 - m^2 - n^2} \times \overline{w \cdot x - k}$  acting on the line OCQ at the distance  $g$  from C; and that therefore the efficacy of the two forces on  $p$  and  $p''$ , to turn the body about HCI, will be represented by the equivalent force  $\frac{2f \times p}{g} \times mn \times \overline{w^2 - x - k^2}$  acting on the line OCQ, at right angles to the plane DOEIFQGH, at the distance  $g$  from C.

3. The body being of any such shape that the section thereof  $hi$ , passing through  $p$  and  $p''$  at right angles to the line  $Ck$ , is a circle whose center is  $k$ ; and every other section thereof, parallel to the said section  $hi$ , a circle whose center is in the line passing through C and  $k$ ; the ordinates corresponding to the abscissæ  $kp$ ,  $kp''$ , in the said circular section  $hi$ , will each be parallel to that diameter (HCI) of the circular section passing through C about which the body will be urged to turn, C being the center of gravity of the body\*: and each of those ordinates will be

\* In a spheroid, cylinder, cone, or any other body conformable, in regard to sections, to this under consideration, (which is called a *solid of revolution*,



be =  $\sqrt{y^2 - w^2}$ ,  $y$  being the radius of such section. Therefore, writing  $2x' \sqrt{y^2 - w^2} \times w'$  instead of  $p$ , it follows that  $\frac{4Af}{g} \times mnx' \times \frac{y^2 - x - k^2}{4} \cdot y^2$ , the whole fluent of  $\frac{4f\sqrt{y^2 - w^2}}{g} \times mnx' \times \frac{y^2 - x - k^2}{4} \times w$ , generated whilst  $w (= kp = kp'')$  from  $o$  becomes equal to the radius  $y$ , (both  $x$  and  $y$  being considered as invariable,) will express the value of (E) the force which, acting on the line  $OCQ$  at the distance  $g$  from  $C$ , would be equivalent to the force of all the particles in the said section, whose thickness is denoted by the indefinitely small quantity  $x'$ ; the distance  $Ck$  being denoted by  $x - k$ , and  $A$  being put for (.78539) the area of a quadrant of a circle whose radius is 1.

Fig. 37. 4. In the spheroid whose proper axis is  $2b$ , and equatorial diameter  $2r$ ; taking  $k$  equal to  $b$ , we have  $y^2 = \frac{r^2}{b^2} \times \overline{2bx - x^2}$ , and  $\frac{y^2 - x - k^2}{4} \cdot y^2 = \frac{r^2}{4b^2} \times \overline{2bx - x^2} - \frac{r^2}{b^2} \times \overline{x - b^2} \times \overline{2bx - x^2} = \frac{r^2 \times r^2 + 4b^2}{4b^2} \times \overline{4b^2x^2 - 4bx^3 + x^4} - r^2 \times \overline{2bx - x^2}$ . Consequently, the fluent of  $\frac{r^2 \times r^2 + 4b^2}{4b^2} \times \overline{4b^2x^2 - 4bx^3 + x^4} \times x - r^2 \times \overline{2bx - x^2} \times x$ , generated whilst

volution, because it may be conceived to be generated by the revolution of some line about the proper axis  $Ck$ ,) the (sum of the forces arising from the rotation about  $OCQ$ , to turn it about a diameter at right angles to  $HCI$  in the circular section whose center is  $C$ , it is obvious, will be = 0; it being manifest that the force of any one particle  $p$ , urging it to turn about such diameter, will be counter-balanced by the force of another particle acting in the opposite direction.

whilst

whilst  $x$  from  $o$  becomes  $= 2b$ , being  $= \frac{4br^2}{15} \times \overline{r^2 - b^2}$ ,  
 we find  $\frac{16Afb r^2}{15g} \times mn \times \overline{r^2 - b^2} = \frac{fmn}{5g} \times \overline{r^2 - b^2} \times S$  the  
 the value of (E) the force which, acting at the distance  $g$   
 from C the center of the spheroid, would be equivalent to  
 the efficacy of the forces acting as above on all the particles  
 of the spheroid to turn it about a diameter of its equator,  
 S being  $(= \frac{16Abr^2}{3})$  the mass or content of the spheroid.

5. In the half of the spheroid mentioned in the pre-  
 ceding article, (cut off at the equator,) taking  $k$  equal to  
 $\frac{5b}{8}$ , we have  $\frac{y^4}{4} - \overline{x - k^2} \cdot y^2 = \frac{r^4}{4b^2} \times \overline{2bx - x^2} - \frac{r^2}{b^2} \times x - \frac{5b^2}{8}$   
 $\times \overline{2bx - x^2}$ . Consequently, the fluent of  $\frac{y^4}{4} - \overline{x - k^2} \cdot y^2 \times x$ ,  
 generated whilst  $x$  from  $o$  becomes  $= b$ , being  $= \frac{br^2}{480} \times$   
 $\overline{64r^2 - 19b^2}$ , we obtain  $\frac{fmn}{320g} \times \overline{64r^2 - 19b^2} \times M$  for the  
 value of (E) the force which, acting at the distance  $g$   
 from the center of gravity of the hemispheroid, would be  
 equivalent to the efficacy of the forces acting as above on  
 all the particles of the hemispheroid to turn it about a  
 diameter of the circular section in which the said center  
 is situated, M being the mass or content of the hemi-  
 spheroid.

6. In the parabolic conoid, the equation of whose gene-  
 rating curve is  $px = y^2$ ; if the height be  $= b$ , and V be  
 at the vertex,  $k$  will be  $= \frac{2b}{3}$ , and  $\frac{y^4}{4} - \overline{x - k^2} \cdot y^2 = \frac{p^2 x^2}{4} - px$   
 $\times x$

$\times x - \frac{2b}{3}$ . Consequently, the fluent of  $\frac{p^2 x^2}{4} - px \times x - \frac{2b}{3}$ 
 $\times \dot{x}$ , generated whilst  $x$  from 0 becomes  $= b$ , being  $= \frac{pb^2}{36}$ 
 $\times \overline{3p - b}$ , the force E will be  $= \frac{fmn}{18g} \times \overline{3pb - b^2} \times M =$ 
 $\frac{fmn}{18g} \times \overline{3r^2 - b^2} \times M$ ,  $r$  being the radius of the base, and
  $M (= 2Apb^2)$  the mass or content of the body.

7. In the solid consisting of two parabolic conoids joined
 together at their bases; if the dimensions of each conoid
 be denoted as in the preceding article, and  $k$  be taken  $= 0$ ,
 $y^2$  will be  $= p \cdot \overline{b - x}$ , and  $\frac{y^4}{4} - \overline{x - k}^2 \cdot y^2 = \frac{p^2 \cdot \overline{b - x}^2}{4} - \overline{b - x}$ .
 Consequently, twice the fluent of  $\frac{p^2 \cdot \overline{b - x}^2}{4} - px \cdot \overline{b - x} \times \dot{x}$ ,
 generated whilst  $x$  from 0 becomes  $= b$ , being  $= \frac{pb^2}{6}$ 
 $\times \overline{p - b}$ , the force E will be  $= \frac{fmn}{6g} \times \overline{r^2 - b^2} \times N$ , the
 mass of the double conoid being denoted by N.

8. In the cone, the radius of whose base is  $= r$ , and
 perpendicular height  $= b$ ; if V be at the vertex,  $k$  will
 be  $= \frac{3b}{4}$ ; and,  $y$  being  $= \frac{rx}{b}$ ,  $\frac{y^4}{4} - \overline{x - k}^2 \cdot y^2$  will be  $= \frac{r^4 x^4}{4b^4}$ 
 $- \frac{r^2 x^2}{b^2} \times x - \frac{3b}{4}$ . Consequently, by taking the fluent of
 $\frac{r^4 x^4}{4b^4} - \frac{r^2 x^2}{b^2} \times x - \frac{3b}{4}$   $\times \dot{x}$ , (generated whilst  $x$  from 0 be-
 comes  $= b$ ), and multiplying it by  $\frac{4A fmn}{g}$ , (as in the pre-
 ceding

ceding examples,) we find the force  $E = \frac{Afmnr^2b}{20g} \times \overline{4r^2 - b^2}$   
 $= \frac{3fmn}{80g} \times \overline{4r^2 - b^2} \times M, \left(\frac{4Ar^2b}{3}\right)$  the content of the body  
 being denoted by M.

9. In the solid consisting of two cones joined together at their bases; if the dimensions of each cone be denoted as in the preceding article, the force E, by proceeding as above, will be found  $= \frac{fmn}{20g} \times \overline{3r^2 - 2b^2} \times N$ , the content of the double cone being denoted by N.

10. In the cylinder whose length is  $b$  and diameter  $2r$ ; taking  $k$  equal to  $\frac{b}{2}$ , we have  $\frac{r^2}{4} - \overline{x - k}^2 \cdot y^2 = r^2 \times \overline{\frac{r^2}{4} - x - \frac{b}{2}}$ ,  $y$  being  $= r$ . Consequently, the fluent of  $\overline{\frac{r^2}{4} - x - \frac{b}{2}}$   $\times x$ , generated whilst  $x$  from 0 becomes  $= b$ , being  $= \frac{r^2b}{4} - \frac{b^3}{12}$ , we find, by multiplying that fluent by  $\frac{4Afmnr^2}{g}$ , the value of  $E = \frac{Afmnr^2b}{3g} \times \overline{3r^2 - b^2} = \frac{fmn}{12g} \times \overline{3r^2 - b^2} \times M$ , the content of the cylinder being denoted by M.

These equivalent forces are distinguished by the name of *motive* forces; the correspondent *accelerative* forces are computed in the following manner.

11. The body being a spheroid whose center is C, and whose proper axis PN is  $= 2b$  and equatorial diameter AB  $= 2r$ ; let F be the accelerative force of a particle of the

Fig. 37.

the body at the distance  $g$  from the axis about which the body is urged to turn, which axis is a diameter of its equator when the motive force is such as we have considered above. Denote  $Ck$  and  $ki$  by  $x - k$  and  $y$  as above; and let the abscissa  $kp$ , and its correspondent ordinate (parallel to the last mentioned axis) in the circle whose radius is  $ki$  be denoted by  $w$  and  $t$  respectively. Then, considering the body as urged to turn about that diameter of its equator which is at right angles to  $AB$ , the accelerative force of every particle in the said ordinate will be  $= \frac{\sqrt{w^2 + x - k^2}}{g}$   $\times F$ , and the motive force of all the particles in the same ordinate will be  $= \frac{\sqrt{w^2 + x - k^2}}{g} \times Ft w' x' = \frac{\sqrt{w^2 + x - k^2}}{g} \times F w' x' \sqrt{y^2 - w^2}$ : to which (by the property of the lever) a motive force  $= \frac{w^2 + x - k^2}{g^2} \times F w' x' \sqrt{y^2 - w^2}$ , acting at the distance  $g$  from the center  $C$ , at right angles to a ray therefrom, would be equivalent. Therefore, considering  $x$  and  $y$  as invariable, and  $w$  only as variable,  $\frac{4F x'}{g^2} \times$  the whole fluent of  $w \sqrt{y^2 - w^2} \times \sqrt{w^2 + x - k^2}$  will denote a force which, acting at the distance  $g$  from  $C$ , would be equivalent to the motive force of all the particles in the circular section  $hi$ , whose radius is  $ki$  and thickness the indefinitely small quantity  $x'$ . Which fluent is  $= A \times \frac{y^2}{4} + x - k^2 \cdot y^2 = \frac{A r^2}{b^2} \times \sqrt{2bx - x^2} \times \frac{r^2}{4b^2} \times \sqrt{2bx - x^2 + x - b^2}$  in our spheroid,  $k$  being taken equal to  $CN = b$ . Consequently

quently  $\frac{4Ar^2F}{b^2g^2} \times$  the whole fluent of  $x \times \sqrt{2bx - x^2} \times \frac{r^2}{4b^2} \times \sqrt{2bx - x^2 + x - b}$ , generated whilst  $x$  from 0 becomes  $= 2b$ , will denote a motive force which, acting at the distance  $g$  from  $C$ , at right angles to a ray therefrom, would be equivalent to the whole motive force urging the spheroid to turn as above mentioned. Such equivalent force will therefore be  $= \frac{F}{5g^2} \times \sqrt{r^2 + b^2} \times S$ : and this being put  $= \frac{fmn}{5g} \times \sqrt{r^2 - b^2} \times S$ , (the value of the same force found in Art. 4.) we find  $F = fgm n \times \frac{r^2 - b^2}{r^2 + b^2}$ ; which will be  $= \frac{gmne^2}{a^2} \times \frac{r^2 - b^2}{r^2 + b^2}$ , if  $f$  be taken  $= \frac{e^2}{a^2}$ , agreeably to the computation in Art. 1. and 2.

Or  $F$  will be denoted by  $\frac{cd}{r} \times \frac{r^2 - b^2}{r^2 + b^2}$ ; if  $r$  be to  $e$  as  $m$  to  $d$ , and as  $n$  to  $c$ ;  $a$  and  $g$  being each  $= r$ .

12. It is evident from what is done in the preceding article, that the circle  $hi$ , whose radius is  $ki (= y)$ , being the section of any solid of revolution whose proper axis coincides with  $Ck$ , if  $C$  be the center of gravity of the body, and  $Ck$  be  $= x - k$ ,  $\frac{4AFx'}{g^2} \times \frac{y^2}{4} + \sqrt{x - k} \cdot y^2$  will denote a force which, acting at the distance  $g$  from  $C$ , as above described, would be equivalent to the motive force of all the particles in the said circular section, whose radius is  $y$  and thickness the indefinitely small quantity  $x'$ . It follows then, from Art. 4. and what is here said, that  $\frac{4AF}{g^2} \times$

H

the

the fluent of  $x \times \frac{y^2}{4} + x - k^2 \cdot y^2$  will be  $= \frac{4Afmn}{g}$   $\times$  the  
 fluent of  $x \times \frac{y^2}{4} - x - k^2 \cdot y^2$ : whence we have the theorem

$$F = fgm n \times \frac{\text{the fluent of } x \times \frac{y^2}{4} - x - k^2 \cdot y^2}{\text{the fluent of } x \times \frac{y^2}{4} + x - k^2 \cdot y^2};$$

the value of  $k$ , and the said fluents, being so taken as to comprehend the whole of the body under consideration, according to the examples given above.

13. Computing by the theorem just now investigated, it appears that,

$$\text{in the hemi-spheroid, } F \text{ will be } = \frac{cd}{r} \times \frac{64r^2 - 19b^2}{64r^2 + 19b^2};$$

$$\text{in the parabolic conoid, } F = \frac{cd}{r} \times \frac{3r^2 - b^2}{3r^2 + b^2};$$

$$\text{in the double conoid, } F = \frac{cd}{r} \times \frac{r^2 - b^2}{r^2 + b^2};$$

$$\text{in the cone, } F = \frac{cd}{r} \times \frac{4r^2 - b^2}{4r^2 + b^2};$$

$$\text{in the double cone, } F = \frac{cd}{r} \times \frac{3r^2 - 2b^2}{3r^2 + 2b^2};$$

$$\text{in the cylinder, } F = \frac{cd}{r} \times \frac{3r^2 - b^2}{3r^2 + b^2};$$

$a$  and  $g$  being each taken  $= r$ ,  $f = \frac{e^2}{r^2}$ ; 1 to  $e$  as  $m$  to  $d$ , and as  $n$  to  $c$ : and reference being had to the respective articles above for the quantities denoted by  $e$ ,  $m$ ,  $n$ ,  $b$ , and  $r$ .

14. Seeing that the force which we have computed above will be  $= 0$ , when,

in

in the hemi-spheroid, the height ( $b$ ) is  $= \frac{8r}{\sqrt{19}}$ ;

in the parabolic conoid, the height ( $b$ ) is  $= 3^{\frac{1}{2}}r$ ;

in the double conoid, the half length ( $b$ ) is  $= r$ ;

in the cone, the height ( $b$ ) is  $= 2r$ ;

in the double cone, the half length ( $b$ ) is  $= \frac{3}{2}r$ ;

in the cylinder, the length ( $b$ ) is  $= 3^{\frac{1}{2}}r$ ;

it is manifest, that each of those bodies will (with respect to its own particles) undisturbedly revolve about any axis whatever passing through its center of gravity, as will a sphere! And it is obvious that, by means of what is done above, other bodies (being solids of revolution, or frustums of such solids) of various forms, may be found having the like property! No more being requisite thereto than that the dimensions of the body be so proportioned that the fluent of  $x \times \frac{z^2 - x^2 - kl^2}{4} \cdot y^2$  be  $= 0$ ;  $k$  being taken equal to the distance from the center of gravity of the body to one end of its proper axis, and  $x$  being considered as flowing from 0, till it becomes equal to the whole length of that axis.

15. Any axis about which a body will so undisturbedly revolve, I call a *permanent axis of rotation*.

## S E C O N D P A R T.

Having hitherto only shewn how to compute the values of E and F when the body is a solid of revolution, I purpose now to shew how those values may be computed when the body is not such a solid; and likewise how the



forces  $E''$  and  $F''$ , at right angles to the direction of the forces  $E$  and  $F$ , may be computed: which forces  $E''$  and  $F''$ , though always  $= 0$  in a solid of revolution, are not generally  $= 0$  in other bodies.

Fig. 34 16. By resolving the force  $\frac{h^2}{a^2} \times p$ , mentioned in Art. I.

into two others, as specified in that article, we found that, in consequence of the centrifugal force of the particle  $p$ , (revolving as there described,) the point  $l$  of the plane  $DOEFQG$  will be urged in a direction at right angles to that plane by a force  $= \frac{u^2}{a^2} \times p$ , let the distance  $lq$  be what it will. Now it is farther observable, that, at the same time, the same point ( $l$ ) will, in consequence of the same centrifugal force, be urged in the direction  $lq$  by a force  $= \frac{u'^2}{a^2} \times p$ , let the distance  $pq$  (denoted by  $u$ ) be what it will,  $u''$  being put for  $\sqrt{h^2 - u^2} = lq$ . And it follows, from the property of the lever, that,  $Cl$  being denoted by  $v$  (as before), if the same point  $l$  be urged in the said direction  $lq$  by a force  $= fu'' \times p$ , the efficacy of that force, to turn the said plane edgeways about an axis passing through the point  $C$  at right angles to  $HCl$ , will be equivalent to the force  $\frac{fu''v}{g} \times p$  acting on the line  $OCQ$ , in a direction parallel to  $lq$ , at the distance  $g$  from the point  $C$ .

17.  $Ck$  being a line in the plane  $O/CQk\rho''$  at right angles to the plane  $DOEFQG$ , let the plane  $k\beta\gamma\delta\gamma''\delta''$ , or a section of the body at right angles to  $Ck$ , be conceived to pass through the

the points  $k, \beta, p, p'', \gamma, \delta, \gamma'', \delta''$ ; and let  $\gamma\beta p\gamma'', \delta k\beta\delta''$  (drawn in the last mentioned plane) be at right angles to each other. Call the abscissa  $\delta\beta$  and its correspondent ordinates  $\beta\gamma, \beta\gamma''$ ;  $x, b', b''$ , respectively: and call the distances  $k\delta, k\delta''$ , and  $\beta p$ ;  $a', a''$ , and  $y$  respectively. Also,  $VCV''$  being the other dimension of the body, measured upon the line  $VCkV''$ , let the distances  $CV, CV''$ , and  $Vk$  be called  $k', k''$ , and  $z$  respectively. Moreover, the sine and cosine of the angle  $kCl$  (to the radius 1) being respectively denoted by  $m$  and  $n$  (as before), let the sine and cosine of the angle  $\beta k p''$  (to the same radius) be denoted by  $s$  and  $t$  respectively;  $p q, p'' l$  being each perpendicular to the plane  $DOEFQG$ , in which the points  $l$  and  $q$  are situated; and  $pp''$  parallel to  $lq$ . Then will

$$\begin{aligned} v & \text{ be } = m t \overline{x - a'} + m s \overline{y - b'} + n \overline{z - k'}, \\ u & = n t \overline{x - a'} + n s \overline{y - b'} - m \overline{z - k'}, \\ u'' & = \overline{t y - b'' - s x - a''}. \end{aligned}$$

It appears therefore, by substitution, that,  $E$  and  $E''$  being respectively put for the values of the motive forces which, acting at the distance  $g$  from  $C$  on the line  $OCQ$ , in directions parallel to  $lp'', lq$  respectively, would be equivalent to the force of all the particles of the body, to turn it about a line passing through the point  $C$  in a plane at right angles to  $OCQ$ ;

$$\begin{aligned} E & \left( = \frac{f}{g} \times \text{the sum of all the } uv \times p \right) \text{ will be} \\ & = \frac{f}{g} \times mn \times \overline{A''t^2 + B''s^2 + 2Ast - K''} + n^2 - m^2 \times \overline{Bs + Ks}, \end{aligned}$$

$$\begin{aligned} E'' & \left( = \frac{f}{g} \times \text{the sum of all the } u''v \times p \right) \\ & = \frac{f}{g} \times m \times \overline{A.t^2 - s^2 + D'st + n \times \overline{Bt - Ks}}; \end{aligned}$$

A being

A being = the sum of all the  $p \times \overline{x - a' . y - b'}$ ,

B = the sum of all the  $p \times \overline{y - b' . z - k'}$ ,

K = the sum of all the  $p \times \overline{x - a' . z - k'}$ ,

A'' = the sum of all the  $p \times \overline{x - a'}^2$ ,

B'' = the sum of all the  $p \times \overline{y - b'}^2$ ,

K'' = the sum of all the  $p \times \overline{z - k'}^2$ ,

and  $D' = B'' - A''$ .

Which sums may commonly be computed in the following manner.

1st. To find the value of A; take the fluent of  $\overline{x - a' . y - b' . y}$ , generated whilst  $y$  from 0 becomes  $= b' + b''$ , considering  $y$  only as variable: which fluent is  $= \frac{1}{2} . x - a' . b'^2 - b''^2$ .

2dly. Having substituted, in that fluent, the values of  $b'$  and  $b''$  in terms of  $x$ ; multiply by  $x$  and take the fluent, generated whilst  $x$  from 0 becomes  $= a' + a''$ , considering  $x$  only as variable.

3dly. Having substituted, in the second fluent, the values of  $a'$  and  $a''$  in terms of  $z$ ; multiply by  $z$  and take the fluent, generated whilst  $z$  from 0 becomes  $= k' + k''$ , considering  $z$  only as variable.

Then will such third fluent express the sum of all the  $p \times \overline{x - a' . y - b'}$ : and in the same manner the values of B, K, A'', B'', K'' are to be computed.

18. It is obvious that, if  $b'$  be every where  $= b''$ ,  $(\frac{1}{2} . b'^2 - b''^2)$  the whole fluent of  $\overline{y - b' . y}$ , generated whilst  $y$  from 0 becomes  $= b' + b''$ , will be  $= 0$ ; therefore A and B will then be each  $= 0$ .

19. If

19. If the line joining the point  $k$  and the center of gravity of the respective section  $k\beta\gamma\delta\gamma''\delta''$  be always at right angles to the line  $\delta k\beta\delta''$ ; the whole fluent of  $\overline{x.x - a.b' + b''}$ , generated whilst  $x$  from  $o$  becomes  $= a' + a''$ , will, by the property of such center, be  $= 0$ : therefore  $K$  will, in such case, be  $= 0$ , whether  $b'$  be  $= b''$  or not.

20. When all the ordinates  $\gamma\beta\gamma'$  can be bisected by a right line (not coincident with  $\delta\beta\delta''$ ),  $b' - b''$  will be  $= q' + q''x - a'$ ,  $q'$  and  $q''$  being some invariable quantities: therefore,  $\frac{1}{2}b' - b''$  (the whole fluent of  $y - b'.y$ ) being  $= \frac{1}{2}q'.b' + b'' + \frac{1}{2}q''.x - a'.b' + b''$ ;  $B$ , by what is said in the preceding article, will be  $= 0$ , if the point  $k$  be situated as described in that article and  $q'$  be  $= 0$ .

If  $q''$  be  $= 0$ , (i. e. if the bisecting line be parallel to  $\delta k\beta\delta''$ ), and  $k$  be so situated;  $A$  and  $K$  will each be  $= 0$ .

Such remarks as these in the last three articles serve to abridge the expressions for the values of  $E$  and  $E''$ .

21. If  $A$ ,  $B$ , and  $K$  be each  $= 0$ , and  $A'' = B'' = K''$ ;  $A''t^2 + B''s^2$  will be  $= A''$ ; and both  $E$  and  $E''$  will vanish, let  $m$  and  $s$  be what they will: therefore, in that case, any line whatever passing through  $C$ , the center of gravity of the body, will be a permanent axis of rotation.

22. If,  $A$ ,  $B$ , and  $K$  being each  $= 0$ ,  $A''$  be  $= B''$ ;  $E''$  will vanish, let  $m$  and  $s$  be what they will; and  $E$  will vanish when  $m$  is  $= 1$ : therefore, in such case, any line whatever passing through  $C$ , in a plane at right angles to  $Ck$ , will be a permanent axis of rotation. Moreover  $E$  being

being also made to vanish by taking  $m = 0$ ,  $Ck$  will also be a permanent axis of rotation.

23. If,  $A$ ,  $B$ , and  $K$  being each  $= 0$ ,  $A''$  be  $= K''$  and  $s = 0$ ;  $E$  and  $E''$  will both vanish, let  $m$  be what it will: therefore, in such case, any line whatever passing through  $C$ , in the plane  $\delta k \beta \delta'' C$ , will be a permanent axis of rotation. And, as  $E$  and  $E''$  will also both vanish when  $m$  and  $s$  are each  $= 1$ ; the body will also have a permanent axis of rotation (passing through  $C$ ) perpendicular to the said plane  $\delta k \beta \delta'' C$ , in which the other such axes will be situated.

24. If,  $A$ ,  $B$ , and  $K$  being each  $= 0$ ,  $B''$  be  $= K''$  and  $s = 1$ ;  $E$  and  $E''$  will both vanish, let  $m$  be what it will: therefore, in such case, any line whatever passing through  $C$ , in a plane  $Ck$  at right angles to the plane  $\delta k \beta \delta'' C$ , will be a permanent axis of rotation. And, as  $E$  and  $E''$  will also both vanish when  $m$  is  $= 1$  and  $s = 0$ , the body will also have a permanent axis of rotation (passing through  $C$ ) parallel to  $\delta k \beta \delta''$ .

25. If,  $B$  and  $K$  being each  $= 0$ ,  $A''$  be  $= B''$ ,  $A'' - K'' \pm A = 0$ , and  $s^2 = t^2$ ;  $E$  and  $E''$  will both vanish let  $m$  be what it will: therefore any line whatever passing through  $C$ , in a certain plane, will then be a permanent axis of rotation. Which plane will pass through  $C$  and  $k$ , and make an angle of  $45^\circ$  with the plane  $\delta k \beta \delta'' C$ , so that  $s$  shall be  $= \sqrt{\frac{1}{2}}$  or  $= -\sqrt{\frac{1}{2}}$ , according as the upper or lower of the two signs prefixed to the term  $A$  takes place. And, as  $E$  and  $E''$  will also both vanish when  $m$  is  $= 1$  and  $s =$

$s^2 = t^2$ , the body will also have a permanent axis of rotation (passing through C) perpendicular to the plane in which the other such axes will be situated.

26. If B and K be each = 0, E and E' will both vanish if  $m$  be = 0: and they will also both vanish if  $m$  be = 1, and  $D'st - A.\overline{s^2 - t^2} = 0$ . Now  $\frac{s}{t}$  being, by this last equation,  $= \frac{D'}{2A} + \frac{\sqrt{D'^2 + 4A^2}}{2A}$ , or  $= \frac{D'}{2A} - \frac{\sqrt{D'^2 + 4A^2}}{2A}$ ; (which can neither become equal, nor imaginary values;) and the sum of the two arcs whose tangents are those values of (the tangent)  $\frac{s}{t}$ , being (as may be easily proved) =  $90^\circ$ ; it is evident that the body will, at least, have *three* permanent axes of rotation; whereof Ck will be one, and the other two will each pass through C and be perpendicular to Ck and to each other.

Moreover, B and K being each = 0, E and E' will both vanish when  $A''t^2 + B''s^2 - K'' + 2Ast$  is = 0, and  $D'st + A.\overline{t^2 - s^2} = 0$  at the same time: in which case  $m$  may express the sine of any angle whatever. Therefore, at first sight it seems possible, that, when A, A'', B'', and K'' have a certain relation amongst themselves, the two values of  $\frac{s}{t}$  in the one of those equations may correspond to

two values of  $\frac{s}{t}$  in the other of those equations; and that, in such case, the body may have an infinite number of permanent axes of rotation (passing through C) in *two* different planes. But, by exterminating  $\frac{s}{t}$  by means of the two last equations, it appears that, for those to be true

I

equations,

equations,  $D''$  must be  $= D'''$ , or  $A^2 = D''D'''$ ;  $D''$  being put for  $K' - A'$ , and  $D'''$  for  $K'' - B''$ . If  $D''$  be  $= D'''$ ,  $A'$  will be  $= B'$ , and  $D' \pm A = 0$ ; and our case the same as that considered in the preceding article. If  $A^2$  be  $= D'D'''$ ;  $\frac{s}{t}$ , in the equation  $D'st + A.\overline{t^2 - s^2} = 0$ , will be  $= -\frac{D''}{A}$ , or  $= \frac{D'''}{A}$ : at the same time  $\frac{s}{t}$  in the other equation ( $A't^2 + B's^2 - K' + 2Ast = 0$ ) will have only one value  $= -\frac{A}{D''}$ . Consequently,  $-\frac{A}{D''}$  being  $= -\frac{D''}{A}$ , but not  $= \frac{D'''}{A}$ , the values of  $E$  and  $E'$  can only vanish together when,  $m$  being  $= 0$ ,  $\frac{s}{t}$  is of any value whatever; or, when  $\frac{m}{n}$  being of any value whatever,  $\frac{s}{t}$  is  $= -\frac{D''}{A}$ ; or,  $m$  being  $= 1$ ,  $\frac{s}{t}$  is  $= \frac{D'''}{A}$ . Therefore it is evident, that, besides the single permanent axis of rotation determined by the equations  $m = 1$  and  $\frac{s}{t} = \frac{D'''}{A}$ , the body (in the case here considered) can only have an infinite number of such axes in *one* certain plane determined by the equation  $\frac{s}{t} = -\frac{D''}{A}$  ( $= -\frac{A}{D''}$ ): to which last mentioned plane, the said single axis will be perpendicular; the sum of the two arcs whose tangents are  $\frac{D''}{A}$  and  $\frac{D'''}{A}$  being  $90^\circ$ .

27. Supposing the general values of  $E$  and  $E'$  each  $= 0$ , and exterminating  $m$  and  $n$  by means of those two equations, we get

P

$$P \frac{s^3}{t^3} + Q \frac{s^2}{t^2} + R \frac{s}{t} + S = 0;$$

$$P \text{ being } = A^2B + AD'''K - BK^2,$$

$$Q = 2B^2K - A^2K - K^3 - D'D'''K - AB.D' + D',$$

$$R = 2BK^2 - A^2B - B^3 + BD'D' + AK.D' - D''',$$

$$S = A^2K + ABD' - B^2K.$$

Now, by our equation so found, it appears that  $\frac{s}{t}$  will,

at least, have *one* real value; and consequently that  $\frac{m}{n}$ ,

which according to our supposition is  $= \frac{Bt - Ks}{A s^2 - t^2 - D'st}$ , will

also, at least, have *one* real value. Therefore, in enquiring concerning the number of permanent axes of rotation in any body, we may suppose OCQ to coincide with Ck; and then, that the values of E and E' may each vanish upon taking  $m = 0$ , our expressions for those values will become

$$E = \frac{f}{g} \times mn \times \frac{A't^2 + B's^2 - K' + 2Ast}{A s^2 - t^2 - D'st},$$

$$E' = \frac{f}{g} \times m \times \frac{A.t^2 - s^2 + D'st}{A s^2 - t^2 - D'st};$$

B and K being necessarily each = 0.

28. Consequently, considering Ck as a permanent axis of rotation, it appears, by art. 26. that any body whatever will, at least, have *three* such axes, situated as described in that article :

And, by art. 22. that, if A and D' be each = 0, any line passing through C, in a plane at right angles to Ck, will be a permanent axis of rotation :

1 2

Also,



Also, by art. 23. 24. 25. and 26. that, if A and D' be each = 0; or A and D'' each = 0; or D' and D' ± A each = 0; or A² - D''D''' = 0; any line passing through C, in a certain plane Ckℓ, (determined as shewn above,) will be a permanent axis of rotation; the body, at the same time, having such an axis (passing through C) perpendicular to the said plane:

Moreover, by art. 21. that, if A, D', and D'' be each = 0; any line whatever passing through the center of gravity (C) will be a permanent axis of rotation.

29. The accelerative forces F and F', corresponding to the motive forces E and E', will, after what has been said, be readily found: it being now obvious enough, that

$$\begin{aligned}
 E \text{ will be} &= \frac{F}{g^2} \times \text{the sum of all the } p \times \overline{v^2 + u^2} \\
 &= \frac{F}{g^2} \times \overline{A't^2 + B's^2 + K' + 2Ast} \\
 &= \frac{f}{g} \times \overline{mn.A't^2 + B's^2 - K' + 2Ast + n^2 - m^2.Bs + Kt}, \\
 E'' \dots \dots &= \frac{F''}{g^2} \times \text{the sum of all the } p \times \overline{v^2 + u'^2} \\
 &= \frac{F''}{g^2} \times \overline{m^2.A'' + B'' + n^2.A''s^2 + B''t^2 + K'' - 2Ast + 2mn.Bs + Kt} \\
 &= \frac{f}{g} \times \overline{m.D'st + A.t^2 - s^2 + n.Bt - Kt}.
 \end{aligned}$$

Whence the values of F and F'' may be immediately obtained.

30. Two or three examples will, I presume, sufficiently explain the method of computing by the theorems above investigated.

*Example*

*Example 1.* Let the body be a *parallelepipedon*, whose dimensions are  $a$ ,  $b$ , and  $k$ . Let the section  $k\beta\gamma\delta\gamma''\delta''$  be parallel to one face thereof, whose length is  $a$  and breadth  $b$ : and,  $Ck$  being conceived to pass through the middle point (or center of gravity) of the said section  $k\beta\gamma\delta\gamma''\delta''$ , and of every other section parallel thereto; let the line  $\delta k\beta\delta''$  divide the section wherein it is drawn (which will be a parallelogram) into two equal parallelograms, so that the length and breadth of each shall be  $a$  and  $\frac{b}{2}$ .

Then, by our remarks above,  $A$ ,  $B$ , and  $K$  will each be  $= 0$ ; and consequently

$$E = \frac{f}{g} \times mn \cdot \overline{A''t^2 + B''s^2 - K''}, \quad E'' = \frac{f}{g} \times mD'st;$$

where  $A''$  will be  $= \frac{a^3bk}{12}$ \*,  $B'' = \frac{ab^3k}{12}$ , and  $K'' = \frac{abk^2}{12}$ .

If  $a$  be  $= b$ ,  $A''$  will be  $= B''$ ,  $D' = 0$ ,  $E = \frac{f}{g} \times mn \cdot \overline{A'' - K''}$ , and  $E'' = 0$ . In which case, (that is when the body is a square prism,) any line passing through the center of gravity of the body, in a plane to which the permanent axis of rotation  $Ck$  is perpendicular, will be such an axis.

If  $a$  be  $= b = k$ ,  $E$  and  $E''$  will each be  $= 0$ , let  $m$  and  $s$  be what they will. It appears therefore, that, in a *cube*, (as in a sphere,) any line whatever passing through the

\* The value of  $A''$  is found in the following manner.

1st. Take the fluent of  $\overline{y \cdot x - a^2}$ : which fluent is  $= 2\sqrt{y \cdot x - a^2}$ .

2dly. Take the fluent of  $2\sqrt{y \cdot x - a^2}$ : which fluent is  $= \frac{4a^3\sqrt{y}}{3} = \frac{a^3b}{12}$ ,

$a'$  being  $= \frac{1}{2}a$ ,  $\sqrt{y} = \frac{1}{2}b$ .

3dly. Take the fluent of  $\frac{a^3b\sqrt{y}}{12}$ : which fluent is  $= \frac{a^3bk}{12} = A''$ .

In like manner the values of  $B''$  and  $K''$  are found.

center

center of gravity of the body will be a permanent axis of rotation.

*Example 2.* Let the body be a triangular prism, whose ends are isosceles triangles; the base of each of which triangles being =  $b$ , the perpendicular thereon from the angle made by the two equal sides (in each) =  $k$ , and the length of the body =  $a$ .

Then, conceiving the section  $k\beta\gamma\delta\gamma''\delta''$  (whereof  $k$  is the middle point) to be parallel to that side of the prism whose length is  $a$  and breadth  $b$ , and supposing the line  $\delta k\beta\delta''$  to divide such section (which will be a parallelogram) lengthwise into two equal parallelograms; A, B, and K will each be = 0; and consequently  $E = \frac{f}{g} \times mn \cdot \overline{A''t^2 + B''s^2 - K''}$ , and  $E'' = \frac{f}{g} \times mD'st$ ; as in the preceding example: but now  $A''$  will be =  $\frac{a^2bk}{24}$ \*,  $B'' = \frac{ab^2k}{48}$ , and  $K'' = \frac{abk^2}{36}$ .

If  $b$  be =  $2\frac{1}{2}a$ , any line passing through the center of gravity of the body, in a plane to which the permanent axis of rotation  $Ck$  is perpendicular, will be such an axis.

\* The value of  $A''$  is computed in the following manner.

1st. Take the fluent of  $\sqrt{b \cdot x - a^2}$ : which fluent is =  $2b \cdot \sqrt{b \cdot x - a^2}$ .

2dly. Take the fluent of  $2b \cdot x \cdot \sqrt{b \cdot x - a^2}$ : which fluent is =  $\frac{4a^2b}{3} = \frac{a^2bx}{12k}$ ;

$$a^2 \text{ being } = \frac{1}{2}a, \quad b = \frac{bx}{2k}$$

3dly. Take the fluent of  $\frac{a^2bx}{12k}$ : which fluent is =  $\frac{a^2bk}{24} = A''$ .

In like manner the values of  $B''$  and  $K''$  are computed.

If

If  $k$  be  $= \frac{3}{2} a$ ; any line passing through the center of gravity of the body, in a plane bisecting the angle made by the two equal sides of the triangle at each end of the prism, will be a permanent axis of rotation; and the body will also have such an axis (passing through its center of gravity) perpendicular to the said bisecting plane.

If  $k$  be  $= \frac{3}{2} b$ ; that is, if the ends of the prism be equilateral triangles; any line passing through the center of gravity of the body, in a plane parallel to the planes of the said triangles, will be a permanent axis of rotation: and the body (as is very obvious) will also have such an axis passing through the center of gravity of the triangle at each end of the prism.

If  $b$  be  $= 2\frac{1}{2}a$  and  $k = \frac{3}{2}a$ ,  $E$  and  $E''$  will each be  $= 0$ , let  $m$  and  $s$  be what they will: therefore, in such a prism, (whose ends will be equilateral triangles,) any line whatever passing through the center of gravity of the body, will be a permanent axis of rotation.

*Example 3.* Let the body be a pyramid; whose base is a parallelogram, the length and breadth whereof are  $a$  and  $b$ ; and the perpendicular height of the body  $= k$ .

Then, conceiving the section  $k\beta\gamma\delta\gamma''\delta''$  (whereof  $k$  is the middle point) to be parallel to the base, and supposing the line  $\delta k\beta\delta''$  to divide such section (which will be a parallelogram) lengthwise into two equal parallelograms;  $A$ ,  $B$ , and  $K$  will each be  $= 0$ ; and the values of  $E$  and  $E''$  will be expressed as in the preceding examples:

amples: but here  $A''$  will be  $= \frac{a^3 b k}{60}$ \*,  $B'' = \frac{a b^3 k}{60}$ , and  $K'' = \frac{a b k^3}{80}$ .

If  $k$  be  $= \frac{2a}{\sqrt{3}}$ ; any line passing through the center of gravity of the body, in a plane passing through the vertex of the pyramid and bisecting its base lengthwise, will be a permanent axis of rotation: and the body will also have such an axis (passing through its center of gravity) perpendicular to the said bisecting plane.

If  $a$  be  $= b$ ; that is, if the body be a square pyramid of any height whatever; any line passing through the center of gravity of the body, in a plane to which the permanent axis of rotation  $Ck$  is perpendicular, will be such an axis.

If  $b$  be  $= a$ , and  $k = \frac{2a}{\sqrt{3}}$ ,  $E$  and  $E''$  will each be  $= 0$ , let  $m$  and  $s$  be what they will: therefore it appears, that, in a square pyramid whose height is to the side of its base as 2 to  $\sqrt{3}$ , any line whatever passing through the center of gravity of the body will be a permanent axis of rotation.

\* The value of  $A''$  is computed as follows.

1st. Take the fluent of  $\sqrt{x-a}^2$ : which fluent is  $= 2\sqrt{x-a}$ .

2dly. Take the fluent of  $2b'\sqrt{x-a}^2$ : which fluent is  $= \frac{4a'^3\sqrt{x}}{3} = \frac{a^3 b x^{\frac{3}{2}}}{12k^3}$ ;

$$a' \text{ being } = \frac{ax}{2k}, \sqrt{x} = \frac{bx}{2k}.$$

3dly. Take the fluent of  $\frac{a^3 b x^{\frac{3}{2}}}{12k^3}$ : which fluent is  $= \frac{a^3 b k}{60} = A''$ .

The values of  $B''$  and  $K''$  are computed in the same manner.

It

It likewise appears by computation, that the *tetrahedron* and the *octahedron* have each the last mentioned property.

31. That the axis of rotation of any revolving body, under no restraint in regard to its rotatory motion, will always pass through the center of gravity of the body, is assumed above as a well-known truth. Indeed it is a very obvious truth: for it appears by what is said above concerning the centrifugal force of a particle of a revolving body, that *the sum of all the  $p \times \overline{x - a'}$  and the sum of all the  $p \times \overline{y - b'}$*  must each be  $= 0^*$ : or else the joint centrifugal force of all the particles of the body would give motion to that point thereof about which the body is, by that same force, urged to turn; that is, the body would be moved entirely out of its place by an internal force arising from its own rotatory motion, without being acted on by any external force to give it a progressive motion; which is absurd. But, by the property of the center of gravity, when *the sum of all the  $p \times \overline{x - a'}$  and the sum of all the  $p \times \overline{y - b'}$*  are each  $= 0$ , our point C is that center.

\* By what is said above, the centrifugal force of the particle  $p$  urges the point  $l$ , in the direction  $lp''$ , with a force

$$fp_u = fp \times \overline{nt.x - a' + ns.y - b' - mx - k'}$$

and in the direction  $lq$ , (at right angles to  $lp''$ ), with a force

$$fp_{u''} = fp \times \overline{t.y - b' - s.x - a'}$$

which equations, upon supposing (as we may)  $Ck$  coincident with  $CO$ , and  $lk\beta d''$  parallel to  $lp''$ , become  $fp_u = fp \times \overline{x - a'}$  and  $fp_{u''} = fp \times \overline{y - b'}$ ;  $m$  and  $s$  being then each  $= 0$ , and  $n$  and  $t$  each  $= 1$ . Therefore, that the joint centrifugal force of all the particles of the body may not give a progressive motion to it, *the sum of all the  $p \times \overline{x - a'}$  and the sum of all the  $p \times \overline{y - b'}$*  must each be  $= 0$ .

K

32. It

32. It is observable, that our method of computing the forces E, E', F, and F" holds true when the body is restrained from revolving freely about an axis passing through its center of gravity, and is made to revolve about any point C which is not that center; excepting such inferences as are expressly derived from the property of such center.

When the force compounded of the two forces E and E', (or F and F'') which we have been computing, is not = 0, it will disturb the rotatory motion of the body so as to cause it to change its axis of rotation every instant, and endeavour to revolve about a new one: the compound motion arising from such perturbation, and likewise from the perturbation caused by the action of an external force, I have, in some measure, explained in the *Philos. Transact.* for the year 1777; and I intend to treat of the same subject more fully in some subsequent Memoir.

MEMOIR

---



---

## M E M O I R V.

*A new Method of obtaining the Sums of certain Series.*

1. **T**HE fluxion of the circular arc  $z$ , whose radius is 1 and cosine  $x$ , being  $= \frac{-\dot{x}}{\sqrt{1-x^2}}$ ;  $n\dot{z}$  is  $= \frac{-n\dot{x}}{\sqrt{1-x^2}} = \frac{-\dot{c}}{\sqrt{1-c^2}}$ ,  $c$  being the cosine of  $(nz)$   $n$  times the arc  $z$ . Whence it appears that  $\frac{\dot{c}}{n}$  is  $= \frac{s\dot{x}}{\sqrt{1-x^2}}$ ;  $s$  denoting  $\sqrt{1-c^2}$ , the sine of  $nz$ .

Moreover  $\dot{s}$  is  $= \frac{-c\dot{c}}{\sqrt{1-c^2}} = nc\dot{z} = \frac{-nc\dot{x}}{\sqrt{1-x^2}}$ . Hence  $\frac{\dot{s}}{n}$   $= \frac{-c\dot{x}}{\sqrt{1-x^2}}$ .

2.  $\frac{n\dot{x}}{\sqrt{x^2-1}}$  being assumed  $= \frac{\dot{y}}{\sqrt{y^2-1}}$ ; we get, by taking the fluents,

$$\text{Log. of } x + \sqrt{x^2-1} \Big|^n = \text{Log. of } y + \sqrt{y^2-1},$$

$$\text{or } x + \sqrt{x^2-1} \Big|^n = y + \sqrt{y^2-1},$$

supposing  $x=1$  when  $y$  is  $= 1$ .

K 2

Whence



$$\text{Whence } y = \frac{x + \sqrt{x^2 - 1} + x - \sqrt{x^2 - 1}}{2},$$

$$\text{and } \sqrt{1 - y^2} = \frac{x + \sqrt{x^2 - 1} - x - \sqrt{x^2 - 1}}{2\sqrt{-1}}.$$

$$\text{But } \frac{nx}{\sqrt{x^2 - 1}} \text{ being } = \frac{j}{\sqrt{y^2 - 1}}, \quad \frac{nx}{\sqrt{1 - x^2}} \text{ will be } = \frac{j}{\sqrt{1 - y^2}};$$

whereof the equation of the fluents is

$n \times \text{circ. arc, rad. 1, cosine } x = \text{circ. arc, rad. 1, cosine } y;$   
 where  $x$  is  $= 1$ , when  $y$  is  $= 1$ , agreeable to the supposition we made above when we took the fluents of the assumed equation by logarithms. Therefore, if  $A$  be put for the least arc whose cosine is  $y$ , and  $C$  for the whole circumference, to the radius 1;  $y$  being the cosine of  $A$ ,  $A + C$ ,  $A + 2C$ ,  $A + 3C$ , &c.  $x$  will be the cosine of  $\frac{A}{n}$ ,  $\frac{A+C}{n}$ ,  $\frac{A+2C}{n}$ , &c. Consequently the values of  $\sqrt{1 - y^2}$  and  $y$ , found above, are respectively equal to the sine and cosine of  $(nx)$   $n$  times the arc whose cosine is  $x$ .

3. Let  $s', s'', s''', \&c.$   $c', c'', c''', \&c.$  denote the sines and cosines of  $x, 2x, 3x, \&c.$  respectively. Then,

if  $x$  be  $= 1$ ;

$s', s'', s''', s^{iv}, \&c.$  will be 0, 0, 0, 0, &c.

$c', c'', c''', c^{iv}, \&c.$  . . . . . 1, 1, 1, 1, &c.

if  $x$  be  $= 0$ ;

$s', s'', s''', s^{iv}, \&c.$  will be 1, 0,  $-1$ , 0, &c.

$c', c'', c''', c^{iv}, \&c.$  . . . . . 0,  $-1$ , 0, 1, &c.

if  $x$  be  $= -1$ ;

$s', s'', s''', s^{iv}, \&c.$  will be 0, 0, 0, 0, &c.

$c', c'', c''', c^{iv}, \&c.$  . . . . .  $-1$ , 1,  $-1$ , 1, &c.

if

if  $x$  be  $= \frac{1}{2}$ ;

$s', s'', s''', s^{iv}, s^v, s^vi, \&c.$  will be  $\frac{\sqrt{3}}{2}, \frac{\sqrt{3}}{2}, 0, -\frac{\sqrt{3}}{2}, -\frac{\sqrt{3}}{2}, 0, \&c.$

$c', c'', c''', c^{iv}, c^v, c^vi, \&c. \dots \dots \frac{1}{2}, -\frac{1}{2}, -1, -\frac{1}{2}, \frac{1}{2}, 1, \&c.$

if  $x$  be  $= -\frac{1}{2}$ ;

$s', s'', s''', s^{iv}, s^v, s^vi, \&c.$  will be  $\frac{\sqrt{3}}{2}, -\frac{\sqrt{3}}{2}, 0, \frac{\sqrt{3}}{2}, -\frac{\sqrt{3}}{2}, 0, \&c.$

$c', c'', c''', c^{iv}, c^v, c^vi, \&c. \dots \dots -\frac{1}{2}, -\frac{1}{2}, 1, -\frac{1}{2}, -\frac{1}{2}, 1, \&c.$   
&c. &c.

4. The Log. of  $1 + u$  being  $= u - \frac{u^2}{2} + \frac{u^3}{3} - \frac{u^4}{4} \&c.$  we  
from thence, by substituting  $\frac{1}{u}$  instead of  $u$ , have

Log. of  $1 + \frac{1}{u} = \text{Log. of } 1 + u - \text{Log. of } u = u^{-1} - \frac{u^{-2}}{2} + \frac{u^{-3}}{3} \&c.$

and, by subtraction, get

Log. of  $u = u - u^{-1} - \frac{u^2 - u^{-2}}{2} + \frac{u^3 - u^{-3}}{3} - \frac{u^4 - u^{-4}}{4} \&c.$

Now, if  $u$  be supposed  $= x + \sqrt{x^2 - 1}$ ,  $\frac{u}{x}$  will be  $= \frac{x}{\sqrt{x^2 - 1}}$

$= \frac{x}{\sqrt{-1} \times \sqrt{1 - x^2}} = \frac{-x\sqrt{-1}}{\sqrt{1 - x^2}}$ ; and the fluent of  $\frac{u}{x}$ , or

Log. of  $u$ ,  $= z\sqrt{-1}$ ,  $z$  denoting the circular arc whose radius is 1 and cosine  $x$ .

Therefore, by substitution, we have, after dividing by  $2\sqrt{-1}$ ,

the series  $\frac{x + \sqrt{x^2 - 1} - x - \sqrt{x^2 - 1}}{2\sqrt{-1}} - \frac{(x + \sqrt{x^2 - 1})^2 - (x - \sqrt{x^2 - 1})^2}{2 \cdot 2\sqrt{-1}}$   
 $+ \frac{(x + \sqrt{x^2 - 1})^3 - (x - \sqrt{x^2 - 1})^3}{3 \cdot 2\sqrt{-1}} \&c.$

$$\text{or its equal } s' - \frac{s''}{2} + \frac{s'''}{3} - \frac{s^{iv}}{4} \&c. = \frac{x}{2};$$

except when (the cofine)  $x$  is  $= -1$ .

*Example.* Let  $x$  be  $= \frac{1}{2}$ , and let  $a$  denote one fourth of the periphery of the circle whose radius is 1. Then,  $s', s'', s''', \&c.$  being  $\frac{\sqrt{3}}{2}, \frac{\sqrt{3}}{2}, 0, -\frac{\sqrt{3}}{2}, -\frac{\sqrt{3}}{2}, 0, \&c.$  as observed in the preceding article; we have

$$\frac{\sqrt{3}}{2} \times 1 - \frac{1}{2} + \frac{1}{4} - \frac{1}{5} + \frac{1}{7} - \frac{1}{8} \&c. = \frac{a}{3}.$$

Whence it appears that  $\frac{1}{1.2} + \frac{1}{4.5} + \frac{1}{7.8} \&c.$  is  $= \frac{2a}{3^{\frac{1}{2}}}$ .

5. Putting  $a$  to denote the same quadrantal arc as in the preceding example,  $2a - x$  will denote the arc whose cofine is  $= -x$  ( $= -c'$ ); and the fines of  $x, 2x, 3x, 4x, \&c.$  being  $s', s'', s''', s^{iv}, \&c.$  the fines of  $2a - x, 2 \times 2a - x, 3 \times 2a - x, 4 \times 2a - x, \&c.$  will be  $s', -s'', s''', -s^{iv}, \&c.$  respectively. Therefore, by substitution, it appears that

$$s' + \frac{s''}{2} + \frac{s'''}{3} + \frac{s^{iv}}{4} \&c. \text{ is } = a - \frac{x}{2},$$

except when  $x$  is  $= 1$ :

and, from this and the theorem in the preceding article, it follows that

$$s'' + \frac{s^{iv}}{2} + \frac{s^{vi}}{3} \&c. \text{ is } = a - x,$$

$$\text{and } s' + \frac{s'''}{3} + \frac{s^v}{5} \&c. = \frac{a}{2};$$

except when  $x$  is  $= 1$ , or  $= -1$ .

*Example.*

*Example.* Taking  $x = \frac{1}{2}$ ; the correspondent values of  $s, s'', s''', \&c.$  in art. 3. being properly substituted in our theorem, we find

$$\frac{\sqrt{3}}{2} \times 1 - \frac{1}{5} + \frac{1}{7} - \frac{1}{11} + \frac{1}{13} - \frac{1}{17} \&c. = \frac{a}{2} :$$

$$\text{whence } \frac{1}{1.5} + \frac{1}{7.11} + \frac{1}{13.17} + \frac{1}{19.23} \&c. = \frac{a}{4\sqrt{3}}.$$

6. Denoting the sine and cosine of  $qx$  by  $s^{(q)}$  and  $c^{(q)}$  respectively; the fluents of  $\frac{s^{(q)} x}{\sqrt{1-x^2}}$  and  $\frac{-c^{(q)} x}{\sqrt{1-x^2}}$  will, by art. 1. be respectively equal to  $\frac{c^{(q)}}{q}$  and  $\frac{s^{(q)}}{q}$ .

Therefore, from our equation

$$s' - \frac{s''}{2} + \frac{s'''}{3} - \frac{s^{iv}}{4} \&c. = \frac{x}{2}, \text{ (found in art. 4.)}$$

by multiplying one side by  $\frac{x}{\sqrt{1-x^2}}$  and the other by its equal  $-x$ , and taking the fluents, we get

$$c' - \frac{c''}{2^2} + \frac{c'''}{3^2} - \frac{c^{iv}}{4^2} \&c. - p'' = -\frac{x^2}{4};$$

$p''$  denoting the series  $1 - \frac{1}{2^2} + \frac{1}{3^2} - \frac{1}{4^2} \&c.$

and the equation being so adjusted that each side shall vanish when  $x$  (or  $c'$ ) is  $= 1$  and  $z = 0$ .

Whence, by taking  $x = 0$ , and  $z = a$ ; ( $c', c'', c''', c^{iv}, c^v, \&c.$  being, according to art. 3. equal to  $0, -1, 0, 1, 0, \&c.$  respectively;) we have

$$\frac{1}{2^2} - \frac{1}{4^2} + \frac{1}{6^2} - \frac{1}{8^2} \&c. - p'' = -\frac{a^2}{4}.$$

But

But  $\frac{1}{2^2} - \frac{1}{4^2} + \frac{1}{6^2} - \frac{1}{8^2} \&c.$  is evidently  $= \frac{p''}{4}$ : consequently  $\frac{p''}{4} - p'' (= -\frac{3p''}{4})$  is  $= -\frac{a^2}{4}$ , and  $p'' = \frac{a^2}{3}$ .

Let  $P''$  denote  $1 + \frac{1}{2^2} + \frac{1}{3^2} + \frac{1}{4^2} \&c.$

and  $Q'' \dots\dots 1 + \frac{1}{3^2} + \frac{1}{5^2} + \frac{1}{7^2} \&c.$

Then  $Q'' + \frac{P''}{4}$  being manifestly  $= P''$ , and  $Q'' - \frac{P''}{4} = p''$ ; it follows that  $P''$  is  $= \frac{2a^2}{3}$ , and  $Q'' = \frac{a^2}{2}$ .

Thus, with great ease, are the sums of those series obtained: and with equal facility our method may be pursued in finding the sums of a great number of other series; the obtaining of which sums has generally been considered as a business of some difficulty.

7. To proceed with perspicuity, let us put

$$P''' = 1 + \frac{1}{2^3} + \frac{1}{3^3} + \frac{1}{4^3} \&c.$$

$$p''' = 1 - \frac{1}{2^3} + \frac{1}{3^3} - \frac{1}{4^3} \&c.$$

$$Q''' = 1 + \frac{1}{3^3} + \frac{1}{5^3} + \frac{1}{7^3} \&c.$$

$$q''' = 1 - \frac{1}{3^3} + \frac{1}{5^3} - \frac{1}{7^3} \&c.$$

$$P^{iv} = 1 + \frac{1}{2^4} + \frac{1}{3^4} + \frac{1}{4^4} \&c.$$

$$p^{iv} = 1 - \frac{1}{2^4} + \frac{1}{3^4} - \frac{1}{4^4} \&c.$$

$$Q^{iv} = 1 + \frac{1}{3^4} + \frac{1}{5^4} + \frac{1}{7^4} \&c.$$

$$q^{iv} = 1 - \frac{1}{3^4} + \frac{1}{5^4} - \frac{1}{7^4} \&c.$$

&c.

&c.

Then,

Then, from our equation

$$c' - \frac{c''}{2^2} + \frac{c'''}{3^2} - \frac{c^{iv}}{4^2} \&c. - p'' = -\frac{x^2}{4},$$

by multiplying one side by  $\frac{-x}{\sqrt{1-x^2}}$  and the other by its

equal  $x$ , and taking the fluents, we get

$$s' - \frac{s''}{2^2} + \frac{s'''}{3^2} - \frac{s^{iv}}{4^2} \&c. - p''x = -\frac{x^2}{12}:$$

and, by repeating the operation, we from thence obtain

$$c' - \frac{c''}{2^4} + \frac{c'''}{3^4} - \frac{c^{iv}}{4^4} \&c. + \frac{p''x^2}{2} - p^{iv} = \frac{x^4}{48}.$$

And so we may proceed as far as we please, obtaining a new series and its sum at each successive operation.

8. From the last equation but one, ( $s', s'', s''', s^{iv}, s^v$ , &c. being respectively equal to 1, 0, -1, 0, 1, &c. when  $x$  is = 0 and  $x = a$ .) we find

$$1 - \frac{1}{3^2} + \frac{1}{5^2} - \frac{1}{7^2} \&c. - ap'' = -\frac{a^2}{12};$$

or,  $ap''$  being found =  $\frac{a^2}{3}$ ,

$$q''' = 1 - \frac{1}{3^2} + \frac{1}{5^2} - \frac{1}{7^2} \&c. = \frac{a^2}{4}.$$

9. The equation  $c' - \frac{c''}{2^4} + \frac{c'''}{3^4} - \frac{c^{iv}}{4^4} \&c. + \frac{p''x^2}{2} - p^{iv} = \frac{x^4}{48}$ , upon taking  $x = 0$  and  $x = a$ , becomes

$$\frac{1}{2^4} - \frac{1}{4^4} + \frac{1}{6^4} - \frac{1}{8^4} \&c. + \frac{a^2 p''}{2} - p^{iv} = \frac{a^4}{48}.$$

But  $\frac{1}{2^4} - \frac{1}{4^4} + \frac{1}{6^4} - \frac{1}{8^4} \&c.$  is manifestly =  $\frac{p^{iv}}{2^4}$ :

consequently  $\frac{p^{iv}}{16} + \frac{a^2 p''}{2} - p^{iv} (= \frac{a^4}{6} - \frac{15p^{iv}}{16})$  is =  $\frac{a^4}{48}$

$$\text{and } p^{iv} = \frac{7a^4}{45}$$

L

More-

Moreover,  $Q^{iv} + \frac{P^{iv}}{2^4}$  being  $= P^{iv}$ , and  $Q^{iv} - \frac{P^{iv}}{2^4} = p^{iv}$ ; it follows that  $P^{iv}$  is  $= \frac{8a^4}{45}$ , and  $Q^{iv} = \frac{a^4}{6}$ .

10. The arc whose cosine is  $x$  being denoted by  $z$ ,  $2a - z$  will denote the arc whose cosine is  $-x$ ; and the cosines of  $z$ ,  $2z$ ,  $3z$ , &c. being  $c'$ ,  $c''$ ,  $c'''$ , &c. the cosines of  $2a - z$ ,  $2 \times 2a - z$ ,  $3 \times 2a - z$ , &c. will be  $-c'$ ,  $c''$ ,  $-c'''$ ,  $c^{iv}$ , &c. Therefore, by substituting accordingly in the theorem deduced in art. 6. it appears that

$$c' + \frac{c''}{2^2} + \frac{c'''}{3^2} \text{ \&c. is } = \frac{2a^2}{3} - az + \frac{z^2}{4}:$$

and, from this and the theorem just now mentioned, it follows that

$$c'' + \frac{c^{iv}}{2^2} + \frac{c^{vi}}{3^2} \text{ \&c. is } = \frac{2a^2}{3} - 2az + z^2,$$

$$\text{and } c' + \frac{c'''}{3^2} + \frac{c^v}{5^2} \text{ \&c. } = \frac{1}{2}a \cdot a - z.$$

from whence other theorems may be derived by our method pursued above.

11. Let  $F$  denote the fluent of  $\frac{\dot{u}}{u} \times \text{Log. of } \sqrt{1+u}$ , which from art. 4. is found

$$= u - \frac{u^2}{2^2} + \frac{u^3}{3^2} - \frac{u^4}{4^2} \text{ \&c. } - p''$$

$$\text{or } = -u^{-1} + \frac{u^{-2}}{2^2} - \frac{u^{-3}}{3^2} + \frac{u^{-4}}{4^2} \text{ \&c. } + \frac{U^2}{2} + p'';$$

$U$  denoting the Log. of  $u$ ,

and  $F$  being supposed  $= 0$  when  $u$  is  $= 1$ .

Let  $G'$  denote the fluent of  $\frac{\dot{u}}{u}F$ ,

$G''$  . . . . . the fluent of  $\frac{\dot{u}}{u}G'$ .

Then

Then will

$$G' \text{ be } = \frac{u^2}{1^2 \cdot 2} - \frac{u^3}{2^2 \cdot 3} + \frac{u^4}{3^2 \cdot 4} \&c. - p''u + p'' - R'$$

$$\text{or } = -\frac{u^{-2}}{1 \cdot 2^2} + \frac{u^{-2}}{2 \cdot 3^2} - \frac{u^{-3}}{3 \cdot 4^2} \&c. + u + p''u - p'' - 1 + R'' \\ + \frac{uU^2}{2} - uU - U;$$

$$G'' = \frac{u^2}{1^2 \cdot 2^2} - \frac{u^3}{2^2 \cdot 3^2} + \frac{u^4}{3^2 \cdot 4^2} \&c. - p''u + \overline{p'' - R'.U} + p'' - S$$

$$\text{or } = \frac{u^{-1}}{1^2 \cdot 2^2} - \frac{u^{-2}}{2^2 \cdot 3^2} + \frac{u^{-3}}{3^2 \cdot 4^2} \&c. + 3u + p''u - p'' - 3 - S \\ + \frac{uU^2}{2} - \frac{U^2}{2} - 2uU - \overline{p'' + 1 - R''.U};$$

$$p'' \text{ being put for } 1 - \frac{1}{2^2} + \frac{1}{3^2} - \frac{1}{4^2} \&c. = \frac{a^2}{3},$$

$$R' \dots \dots \dots \text{ for } \frac{1}{1^2 \cdot 2} - \frac{1}{2^2 \cdot 3} + \frac{1}{3^2 \cdot 4} - \frac{1}{4^2 \cdot 5} \&c.$$

$$R'' \dots \dots \dots \text{ for } \frac{1}{1 \cdot 2^2} - \frac{1}{2 \cdot 3^2} + \frac{1}{3 \cdot 4^2} - \frac{1}{4 \cdot 5^2} \&c.$$

$$S \dots \dots \dots \text{ for } \frac{1}{1^2 \cdot 2^2} - \frac{1}{2^2 \cdot 3^2} + \frac{1}{3^2 \cdot 4^2} - \frac{1}{4^2 \cdot 5^2} \&c.$$

where it is observable that  $R' + R''$  is  $= 2p'' - 1 = \frac{2a^2}{3} - 1$ ,  
and  $R' - R'' = S$ .

Whence, by multiplying each side of the equation arising from the two values of  $G''$  by  $\frac{u^{-\frac{1}{2}}}{2\sqrt{-1}}$ , bringing the two series together in order to form only one series, and substituting according to the method explained in art. 4. we find

$$\frac{\binom{\frac{1}{2}}{s}}{1^2 \cdot 2^2} - \frac{\binom{\frac{1}{2}}{s}}{2^2 \cdot 3^2} + \frac{\binom{\frac{1}{2}}{s}}{3^2 \cdot 4^2} - \frac{\binom{\frac{1}{2}}{s}}{4^2 \cdot 5^2} \&c. = s \times \frac{\binom{\frac{1}{2}}{s}}{3} - \frac{x^2}{2} + 3 - 2cx;$$

L 2

(\frac{1}{2})  
s,



$\left(\frac{1}{2}\right) \left(\frac{1}{2}\right) \left(\frac{1}{2}\right)$   $s, s, s, \&c.$  and  $\left(\frac{1}{2}\right) \left(\frac{1}{2}\right) \left(\frac{1}{2}\right)$   $c, c, c, \&c.$  denoting the sines and cosines of  $\frac{1}{2}z, \frac{1}{2}z, \frac{1}{2}z, \&c.$  respectively.

12. The sines of  $\frac{1}{2} \times \overline{2a - z}, \frac{1}{2} \times \overline{2a - z}, \frac{1}{2} \times \overline{2a - z}, \&c.$  being  $\left(\frac{1}{2}\right) \left(\frac{1}{2}\right) \left(\frac{1}{2}\right) \left(\frac{1}{2}\right)$   $c, -c, c, -c, \&c.$  respectively; and the cosine of  $\frac{1}{2} \times \overline{2a - z}$  equal to  $\left(\frac{1}{2}\right) s$ ; we find by substitution

$$\left(\frac{1}{2}\right) \frac{c}{1 \cdot 2^2} + \left(\frac{1}{2}\right) \frac{c}{2^2 \cdot 3^2} + \left(\frac{1}{2}\right) \frac{c}{3^2 \cdot 4^2} \&c. = 2 \left(\frac{1}{2}\right) s \cdot \overline{2a - z} + c \cdot \frac{4a^2}{3} - 2az + \frac{z^2}{2} - 3.$$

*Example.* If  $x (= c')$  be = 1,  $z$  and  $s$  will each be = 0; and  $\left(\frac{1}{2}\right) \left(\frac{1}{2}\right) \left(\frac{1}{2}\right)$   $c, c, c, \&c.$  each equal to 1. Therefore we have by our theorem

$$\frac{1}{1^2 \cdot 2^2} + \frac{1}{2^2 \cdot 3^2} + \frac{1}{3^2 \cdot 4^2} \&c. = \frac{4a^2}{3} - 3.$$

13. F denoting the fluent of  $\frac{u}{u} \times \text{Log. of } 1 + u$  as in art. 11. let

H' denote the fluent of  $u \dot{u} F$ ,

H'' . . . . . the fluent of  $\frac{u}{u} \dot{H}'$ ,

H''' . . . . . the fluent of  $u \dot{u} H''$ ,

H'''' . . . . . the fluent of  $\frac{u}{u} \dot{H}'''$ .

Then will

$$H' \text{ be } = \frac{u^2}{1^2 \cdot 3} - \frac{u^4}{2^2 \cdot 4} + \frac{u^5}{3^2 \cdot 5} \&c. - \frac{p'' u^2}{2} + \frac{p'''}{2} - R''''$$

$$\begin{aligned} \text{or} &= \frac{u^{-1}}{1 \cdot 3^2} - \frac{u^{-2}}{2 \cdot 4^2} + \frac{u^{-3}}{3 \cdot 5^2} \&c. - u + \frac{u^2}{8} + \frac{p'' u^2}{2} \\ &\quad + \frac{u^2 U^2}{4} - \frac{u^2 U}{4} + \frac{U}{4} - \frac{p''}{2} + \frac{7}{8} - R^{iv}; \\ H'' &= \frac{u^2}{1^2 \cdot 3^2} - \frac{u^4}{2^2 \cdot 4^2} + \frac{u^3}{3^2 \cdot 5^2} \&c. - \frac{p'' u^2}{4} + \frac{p''}{2} - R''' \cdot U + \frac{p''}{4} - S'' \\ \text{or} &= -\frac{u^{-1}}{1^2 \cdot 3^2} + \frac{u^{-2}}{2^2 \cdot 4^2} - \frac{u^{-3}}{3^2 \cdot 5^2} \&c. - u + \frac{3u^2}{16} + \frac{p'' u^2}{4} \\ &\quad + \frac{u^2 U^2}{8} + \frac{U^2}{8} - \frac{u^2 U}{4} - \frac{p''}{2} - \frac{7}{8} + R^{iv} \cdot U - \frac{p''}{4} + \frac{13}{16} + S''; \\ R''' \text{ being} &= \frac{1}{1^2 \cdot 3} - \frac{1}{2^2 \cdot 4} + \frac{1}{3^2 \cdot 5} \&c. \\ R^{iv} \dots &= \frac{1}{1 \cdot 3^2} - \frac{1}{2 \cdot 4^2} + \frac{1}{3 \cdot 5^2} \&c. \\ S'' \dots &= \frac{1}{1^2 \cdot 3^2} - \frac{1}{2^2 \cdot 4^2} + \frac{1}{3^2 \cdot 5^2} \&c. \end{aligned}$$

where it is observable, that

$$R''' + R^{iv} \text{ is } = \frac{1}{8}, \text{ and } R''' - R^{iv} = 2S''.$$

Whence, by multiplying each side of the equation arising from the two values of  $H''$  by  $\frac{u^{-1}}{2}$ , bringing the two series together, and substituting according to our method, we find

$$\frac{c''}{1^2 \cdot 3^2} - \frac{c'''}{2^2 \cdot 4^2} + \frac{c^{iv}}{3^2 \cdot 5^2} - \frac{c^v}{4^2 \cdot 6^2} \&c. = \frac{c' z}{4} + c' \cdot \frac{a^2}{6} - \frac{z^2}{8} + \frac{3}{16} - \frac{1}{2}.$$

Moreover, the cofines of  $2a - z$ ,  $2 \times 2a - z$ ,  $3 \times 2a - z$ , &c. being  $-c'$ ,  $c''$ ,  $-c'''$ ,  $c^{iv}$ , &c. as observed in art. 10. we get by substitution

$$\frac{c''}{1^2 \cdot 3^2} + \frac{c'''}{2^2 \cdot 4^2} + \frac{c^{iv}}{3^2 \cdot 5^2} \&c. = \frac{c'}{4} \cdot 2a - z + c' \cdot \frac{a^2}{3} - \frac{az}{2} + \frac{z^2}{8} - \frac{3}{16} - \frac{1}{2}.$$

And, from this and the preceding theorem, we have

∴

$$\frac{c''}{1^2 \cdot 3^2} + \frac{c^{iv}}{3^2 \cdot 5^2} + \frac{c^{vi}}{5^2 \cdot 7^2} \&c. = \frac{s'a}{4} + \frac{c'a}{4} \frac{a-z}{a-z} - \frac{1}{2},$$

$$\frac{c'''}{1^2 \cdot 2^2} + \frac{c^v}{2^2 \cdot 3^2} + \frac{c^{vii}}{3^2 \cdot 4^2} \&c. = 4s'a - z + c' \frac{4a^2}{3} - 4az + 2z^2 - 3.$$

*Example 1.* If  $x$  be  $= 1$ ,  $s'$  and  $z$  will each be  $= 0$ ; and  $c'$ ,  $c''$ ,  $c'''$ , &c. each  $= 1$ : therefore it appears by substitution, that

$$\frac{1}{1^2 \cdot 3^2} - \frac{1}{2^2 \cdot 4^2} + \frac{1}{3^2 \cdot 5^2} \&c. \text{ is } = \frac{a^2}{6} - \frac{5}{16},$$

$$\frac{1}{1^2 \cdot 3^2} + \frac{1}{2^2 \cdot 4^2} + \frac{1}{3^2 \cdot 5^2} \&c. = \frac{a^2}{3} - \frac{11}{16},$$

$$\frac{1}{1^2 \cdot 3^2} + \frac{1}{3^2 \cdot 5^2} + \frac{1}{5^2 \cdot 7^2} \&c. = \frac{a^2}{4} - \frac{1}{2}.$$

*Example 2.* If  $x$  be  $= 0$ ,  $s'$  will be  $= 1$ ,  $z = a$ ; and  $c'$ ,  $c''$ ,  $c'''$ , &c. respectively equal to  $0$ ,  $-1$ ,  $0$ ,  $1$ , &c. therefore, by substituting accordingly, it appears that

$$\frac{1}{1^2 \cdot 3^2} - \frac{1}{3^2 \cdot 5^2} + \frac{1}{5^2 \cdot 7^2} - \frac{1}{7^2 \cdot 9^2} \&c. \text{ is } = \frac{1}{2} - \frac{a}{4}:$$

and, from this and the theorem next above, that

$$\frac{1}{1^2 \cdot 3^2} + \frac{1}{5^2 \cdot 7^2} + \frac{1}{9^2 \cdot 11^2} \&c. \text{ is } = \frac{1}{8} \frac{a^2}{a^2 - a}.$$

14. By proceeding as in the preceding article, the values of  $H'''$  and  $H^{iv}$  will be obtained: and then, by multiplying each side of the equation arising from the two values of  $H^{iv}$  by  $\frac{u^{-2}}{2}$ , bringing the two series together, and substituting as in that article, it will appear that

'''

$$\frac{c'''}{1^2 \cdot 3^2 \cdot 5^2} - \frac{c^{iv}}{2^2 \cdot 4^2 \cdot 6^2} + \frac{c^v}{3^2 \cdot 5^2 \cdot 7^2} - \frac{c^{vi}}{4^2 \cdot 6^2 \cdot 8^2} \&c. \text{ is } = \begin{cases} \frac{3s''z}{128} - \frac{c'}{9} + \\ c'' \cdot \frac{a^2}{96} - \frac{z^2}{128} + \frac{23}{32 \cdot 32} \\ + \frac{a^2}{48} - \frac{z^2}{64} + \frac{1}{64}. \end{cases}$$

Hence, by substitution; the cofines of  $\overline{2a - z}$ ,  $2 \times \overline{2a - z}$ ,  $3 \times \overline{2a - z}$ , &c. being  $-c'$ ,  $c''$ ,  $-c'''$ ,  $c^{iv}$ , &c. and the sine of  $2 \times \overline{2a - z}$  being  $= -s''$ , as observed above; we find

$$\frac{c'''}{1^2 \cdot 3^2 \cdot 5^2} + \frac{c^{iv}}{2^2 \cdot 4^2 \cdot 6^2} + \frac{c^v}{3^2 \cdot 5^2 \cdot 7^2} \&c. = \begin{cases} \frac{3s'' \overline{2a - z}}{128} - \frac{c'}{9} + \\ c'' \cdot \frac{a^2}{48} - \frac{az}{32} + \frac{z^2}{128} - \frac{23}{32 \cdot 32} \\ + \frac{a^2}{24} - \frac{az}{16} + \frac{z^2}{64} - \frac{1}{64}. \end{cases}$$

And, from these last two theorems, we have

$$\begin{aligned} \frac{c'''}{1^2 \cdot 3^2 \cdot 5^2} + \frac{c^v}{3^2 \cdot 5^2 \cdot 7^2} + \frac{c^{vii}}{5^2 \cdot 7^2 \cdot 9^2} \&c. &= \frac{3s''a}{128} - \frac{c'}{9} + \frac{c''a}{64} \overline{a - z} + \frac{a^2 - az}{32} \\ \frac{c^{iv}}{1^2 \cdot 2^2 \cdot 3^2} + \frac{c^{vi}}{2^2 \cdot 3^2 \cdot 4^2} + \frac{c^{viii}}{3^2 \cdot 4^2 \cdot 5^2} \&c. &= \begin{cases} \frac{3s'' \overline{a - z}}{2} + \frac{2a^2}{3} - 2az + z^2 - 1 \\ + c'' \cdot \frac{a^2}{3} - az + \frac{z^2}{2} - \frac{23}{16}. \end{cases} \end{aligned}$$

*Example 1.* If  $x$  be  $= 1$ ,  $s''$  and  $z$  will each be  $= 0$ , and  $c'$ ,  $c''$ ,  $c'''$ , &c. each  $= 1$ : it appears therefore, by substitution, that

$$\begin{aligned} \frac{1}{1^2 \cdot 3^2 \cdot 5^2} - \frac{1}{2^2 \cdot 4^2 \cdot 6^2} + \frac{1}{3^2 \cdot 5^2 \cdot 7^2} \&c. \text{ is } &= \frac{a^2}{32} - \frac{673}{9 \cdot 32 \cdot 32} \\ \frac{1}{1^2 \cdot 3^2 \cdot 5^2} + \frac{1}{2^2 \cdot 4^2 \cdot 6^2} + \frac{1}{3^2 \cdot 5^2 \cdot 7^2} \&c. &= \frac{a^2}{16} - \frac{1375}{9 \cdot 32 \cdot 32}. \end{aligned}$$

$$\frac{1}{1^2 \cdot 3^2 \cdot 5^2} + \frac{1}{3^2 \cdot 5^2 \cdot 7^2} + \frac{1}{5^2 \cdot 7^2 \cdot 9^2} \&c. = \frac{3a^2}{64} - \frac{1}{9},$$

$$\frac{1}{1^2 \cdot 2^2 \cdot 3^2} + \frac{1}{2^2 \cdot 3^2 \cdot 4^2} + \frac{1}{3^2 \cdot 4^2 \cdot 5^2} \&c. = a^2 - \frac{39}{16}.$$

*Example 2.* If  $x$  be  $= 0$ ,  $s''$  will be  $= 0$ ,  $z = a$ ; and  $c'$ ,  $c''$ ,  $c'''$ , &c. respectively equal to  $0$ ,  $-1$ ,  $0$ ,  $1$ , &c. therefore it follows, that

$$\frac{1}{1^2 \cdot 2^2 \cdot 3^2} - \frac{1}{2^2 \cdot 3^2 \cdot 4^2} + \frac{1}{3^2 \cdot 4^2 \cdot 5^2} \&c. \text{ is } = \frac{7}{16} - \frac{a^2}{6}.$$

And, by addition, we have

$$\frac{1}{1^2 \cdot 2^2 \cdot 3^2} + \frac{1}{3^2 \cdot 4^2 \cdot 5^2} + \frac{1}{5^2 \cdot 6^2 \cdot 7^2} \&c. = \frac{5a^2}{12} - 1:$$

and, by subtraction,

$$\frac{1}{2^2 \cdot 3^2 \cdot 4^2} + \frac{1}{4^2 \cdot 5^2 \cdot 6^2} + \frac{1}{6^2 \cdot 7^2 \cdot 8^2} \&c. = \frac{7a^2}{12} - \frac{23}{16}.$$

15. The fluent of  $\frac{u}{1+u^2}$ , beginning when  $x$  is  $= 1$ ,

$$\text{is } = u - \frac{u^3}{3} + \frac{u^5}{5} - \frac{u^7}{7} \&c. - \frac{a}{2},$$

$$\text{or } = -u^{-1} + \frac{u^{-3}}{3} - \frac{u^{-5}}{5} + \frac{u^{-7}}{7} \&c. + \frac{a}{2}.$$

From which equation, by bringing the two series together and substituting according to our method, we have

$$c' - \frac{c''}{3} + \frac{c'''}{5} - \frac{c''''}{7} \&c. = \frac{a}{2};$$

except when  $x (= c')$  is negative, or  $= 0$ .

Hence, by proceeding as in art. 7. many other theorems may be readily deduced: and more theorems may be de-

derived from the two values of the fluent of  $\frac{u}{1+u^2}$ , by proceeding as in art. 11.

16. By

16. By art. 4. we have

$$\frac{1}{2} \text{Log. } \sqrt{1+u} - \text{Log. } u = \begin{cases} u - \frac{u^2}{2} + \frac{u^3}{3} - \frac{u^4}{4} \&c. \\ + u^{-1} - \frac{u^{-2}}{2} + \frac{u^{-3}}{3} - \frac{u^{-4}}{4} \&c. \end{cases}$$

Hence, by substituting  $x + \sqrt{x^2 - 1}$  for  $u$ , according to our method pursued above, we get

$$\frac{1}{2} \text{Log. } \sqrt{1+x} + \frac{1}{2} \text{Log. } 2 = c' - \frac{c''}{2} + \frac{c'''}{3} - \frac{c^{iv}}{4} \&c.$$

And, writing  $-x$  instead of  $x$ ; and  $-c'$ ,  $-c''$ ,  $-c'''$ ,  $-c^{iv}$ ,  $\&c.$  instead of  $c'$ ,  $c''$ ,  $c'''$ ,  $c^{iv}$ ,  $\&c.$  respectively, agreeably to what is said in art. 10. we have

$$\frac{1}{2} \text{Log. } \sqrt{1-x} + \frac{1}{2} \text{Log. } 2 = -c' - \frac{c''}{2} - \frac{c'''}{3} \&c.$$

Therefore it is obvious, that

$$\frac{1}{2} \text{Log. } \frac{1}{\sqrt{1-x^2}} - \frac{1}{2} \text{Log. } 2 \text{ is } = \frac{c''}{2} + \frac{c^{iv}}{4} + \frac{c^{vi}}{6} \&c.$$

$$\text{and } \frac{1}{4} \text{Log. } \frac{1+x}{1-x} = c' + \frac{c'''}{3} + \frac{c^{v}}{5} \&c.$$

By which last theorem, and that in the preceding article, we find

$$\frac{1}{4} \text{Log. } \frac{1+x}{1-x} + \frac{a}{4} = c' + \frac{c^{v}}{5} + \frac{c^{ix}}{9} \&c.$$

$$\text{and } \frac{1}{4} \text{Log. } \frac{1+x}{1-x} - \frac{a}{4} = \frac{c'''}{3} + \frac{c^{vii}}{7} + \frac{c^{xi}}{11} \&c.$$

when  $x$  is a positive quantity.

*Example.* If  $x$  be the cosine of  $45^\circ$ ;  $c'$ ,  $-c''$ ,  $-c'''$ ,  $c^{iv}$ ,  $c^{ix}$ ,  $\&c.$  will each be  $= \frac{1}{\sqrt{2}}$ ; and consequently

$$1 - \frac{1}{5} + \frac{1}{9} - \frac{1}{13} \&c. = \frac{a + \text{Log. } 1 + \sqrt{2}}{2^{\frac{3}{2}}},$$

$$\text{and } \frac{1}{3} - \frac{1}{7} + \frac{1}{11} - \frac{1}{15} \&c. = \frac{a - \text{Log. } 1 + \sqrt{2}}{2^{\frac{3}{2}}}:$$

which series, it may be observed, are respectively equal to the fluents of  $\frac{y}{1+y^2}$  and  $\frac{y^3}{1+y^2}$ , generated whilst  $y$  from  $\alpha$  becomes equal to 1.

17. The fluent of  $\frac{u^{m-1}u}{1+u^n}$ ,

$$\text{is } = \frac{u^m}{m} - \frac{eu^{m+n}}{m+n} + \frac{e''u^{m+2n}}{m+2n} \&c. - M',$$

$$\text{or } = -\frac{u^{m-n}}{en-m} + \frac{eu^{m-n-n}}{en+n-m} - \frac{e''u^{m-n-2n}}{en+2n-m} \&c. + M'';$$

$$e' \text{ being } = \frac{e.e+1}{2}, e'' = \frac{e.e+1.e+2}{2.3}, \&c.$$

and  $M'$  and  $M''$  respectively denoting

$$\text{the series } \frac{1}{m} - \frac{e}{m+n} + \frac{e''}{m+2n} \&c.$$

$$\text{and } \frac{1}{en-m} - \frac{e}{en+n-m} + \frac{e''}{en+2n-m} \&c.$$

which are the respective fluents of  $\frac{u^{m-1}u}{1+u^n}$  and  $\frac{u^{m-n-1}u}{1+u^n}$ ,

generated whilst  $u$  from  $\alpha$  becomes equal to 1.

Whence, by multiplying by  $u^{m-2n-1}u$  and taking the fluents, we have

$$\frac{u^{en-m}}{m.en-m} - \frac{eu^{en+n-m}}{m+n.en+n-m} + \frac{e'u^{en+2n-m}}{m+2n.en+2n-m} \&c. - \frac{M'u^{en-2m}}{en-2m} + \frac{M'}{en-2m}$$

$$= \frac{u^{-m}}{m.en-m} - \frac{eu^{-m-n}}{m+n.en+n-m} + \frac{e'u^{-m-2n}}{m+2n.en+2n-m} \&c. + \frac{M''u^{en-2m}}{en-2m} - \frac{M''}{en-2m}$$

And hence, bringing the two series together, after multiplying by  $\frac{u^{m-\frac{1}{2}en}}{2\sqrt{-1}}$  and substituting  $x + \sqrt{x^2 - 1}$  for  $u$ , &c. according to our method, we find

$$\frac{M' + M'' \cdot \frac{(\frac{1}{2}en-m)}{s}}{en-2m} =$$

$$\frac{\frac{(\frac{1}{2}en)}{s}}{m.en-m} - \frac{\frac{(\frac{1}{2}en+n)}{s}}{m+n.en+n-m} + \frac{\frac{(\frac{1}{2}en+2n)}{s}}{m+2n.en+2n-m} - \frac{\frac{(\frac{1}{2}en+3n)}{s}}{m+3n.en+3n-m} \&c.$$

$\frac{(\frac{1}{2}en-m)}{s}, \frac{(\frac{1}{2}en)}{s}, \frac{(\frac{1}{2}en+n)}{s}, \frac{(\frac{1}{2}en+2n)}{s}, \&c.$  denoting the sines of  $\frac{1}{2}en - m.z, \frac{1}{2}en.z, \frac{1}{2}en + n.z, \frac{1}{2}en + 2n.z, \&c.$  respectively; of which arcs, the cosines will be denoted by  $\frac{(\frac{1}{2}en-m)}{c}, \frac{(\frac{1}{2}en)}{c}, \frac{(\frac{1}{2}en+n)}{c}, \frac{(\frac{1}{2}en+2n)}{c}, \&c.$  respectively; and the sines and cosines of other arcs, which are multiples or sub-multiples of the arc  $z$ , will be expressed in like manner.

Now it is observable that, if  $y$  be  $= \frac{1}{u}, \frac{u^{m-1}u}{1+u^n}$  will be  $= -\frac{y^{en-m-1}y}{1+y^n}$ ; therefore,  $y$  decreasing from 1 to 0 whilst  $z$  increases from 1 to infinity, it follows, that the fluent of  $\frac{u^{en-m-1}u}{1+u^n}$ , generated whilst  $u$  from 0 becomes equal to 1,



will be equal to the fluent of  $\frac{u^{m-1} \dot{u}}{1+u^n}$ , generated whilst  $u$  from 1 becomes infinite; and consequently, that  $(M' + M'')$  the sum of the fluents of  $\frac{u^{m-1} \dot{u}}{1+u^n}$  and  $\frac{u^{en-m-1} \dot{u}}{1+u^n}$ , generated whilst  $u$  from 0 becomes equal to 1, will be equal to the whole fluent of  $\frac{u^{m-1} \dot{u}}{1+u^n}$ , or of  $\frac{u^{en-m-1} \dot{u}}{1+u^n}$ , generated whilst  $u$  from 0 becomes infinite. Consequently, denoting this last mentioned whole fluent by  $M$ , our series which we found equal to

$$\frac{M' + M''}{en - 2m} \cdot \frac{(1/2)^{en-m}}{s} \text{ will be } = \frac{M}{en - 2m} \cdot \frac{(1/2)^{en-m}}{s} . .$$

18. If  $nx$  be  $= 2a$ ;  $s$ ,  $-\frac{(1/2)^{en}}{s}$ ,  $-\frac{(1/2)^{en+n}}{s}$ ,  $-\frac{(1/2)^{en+2n}}{s}$ ,  $-\frac{(1/2)^{en+3n}}{s}$ , &c. will each be equal to the sine of  $ea$ ; and  $\frac{(1/2)^{en-m}}{s}$  equal to the sine of  $\frac{en-2m}{n}a$ : therefore it appears, that.

$$\frac{M}{en - 2m} \cdot \frac{(1/2)^{en-m}}{s} \text{ will then be } =$$

$$\frac{1}{m.en-m} + \frac{s}{m+n.en+n-m} + \frac{s^2}{m+2n.en+2n-m} + \frac{s^3}{m+3n.en+3n-m} \text{ \&c.}$$

$(1/2)^{en}$  and  $\frac{(1/2)^{en-m}}{s}$  being as just now mentioned.

But this series is equal to  $\frac{N' - N''}{en - 2m}$ ,  $N'$  and  $N''$  being the fluents of  $\frac{u^{m-1} \dot{u}}{1-u^n}$  and  $\frac{u^{en-m-1} \dot{u}}{1-u^n}$  respectively, generated whilst

whilst  $u$  from 0 becomes equal to 1. Hence therefore we have the remarkable theorem

$$\text{Sine of } \frac{en - 2m}{n} a \times M = \text{fine of } ea \times \overline{N' - N''},$$

which is of considerable use in the calculation of fluents.

*Example.* If  $e$  be = 1, and  $m = rn$ ;  $s$  will be = 1,  $(\frac{1}{2}en - m)$   
 $s = c$ , and (by our Appendix)  $M = \frac{2a}{n \cdot s}$ . Consequently

$$\frac{2a}{n \cdot s} \cdot \frac{c}{s} \text{ will be } = N' - N'',$$

the difference of the fluents of  $\frac{u^{m-1}u}{1-u^n}$  and  $\frac{u^{m-1}u}{1-u^n}$ , generated whilst  $u$  from 0 becomes equal to 1; and

$$\frac{2a}{1-2r} \cdot \frac{c}{s} = \frac{1}{r \cdot 1 - r} + \frac{1}{1+r \cdot 2 - r} + \frac{1}{2+r \cdot 3 - r} + \frac{1}{3+r \cdot 4 - r} \&c.$$

$s$  and  $c$  being the sine and cosine of  $2ra$  respectively.

19. Multiplying each side of the general theorem in art. 17. by  $\frac{x}{\sqrt{1-x^2}}$ , and taking the fluents, we have

$$\frac{M}{en - 2m} \times \overline{c - 1} + S = \text{the series}$$

$$\frac{c}{m \cdot en - m} - \frac{e}{m + n \cdot en + 2n \cdot en + n - m} + \frac{e''}{m + 2n \cdot en + 4n \cdot en + 2n - m} \&c.$$

$$S \text{ being } = \frac{1}{m \cdot en - m} - \frac{e}{m + n \cdot en + 2n \cdot en + n - m} \&c.$$

20. If

20. If  $nz$  be  $= 2a$ ;  $c^{(\frac{1}{2}n)}$ ,  $-c^{(\frac{1}{2}n+n)}$ ,  $c^{(\frac{1}{2}n+2n)}$ ,  $-c^{(\frac{1}{2}n+3n)}$ , &c. will be equal to each other; therefore it follows, that

$$\frac{M}{en - 2m^2} \times \overbrace{c^{(\frac{1}{2}n-m)}} - 1 + S \text{ will then be } =$$

$$c^{(\frac{1}{2}n)} \times \frac{1}{m \cdot en \cdot en - m} + \frac{e}{m + n \cdot en + 2n \cdot en + n - m} + \frac{e^2}{m + 2n \cdot en + 4n \cdot en + 2n - m} \&c.$$

*Example.* If  $e$  be  $= 1$ , and  $m = rn$ ; the cosine  $c^{(\frac{1}{2}n-m)}$  will be  $= s^{(m)}$ , and the other cosines  $c^{(\frac{1}{2}n)}$ ,  $c^{(\frac{1}{2}n+n)}$ , &c. each  $= 0$ : consequently  $M$  being then  $= \frac{2a}{n \cdot s^{(m)}}$ ,

$$\frac{1-s^{(m)}}{s^{(m)}} \times \frac{2a}{1-2r} \text{ will be } (= S) =$$

$$\frac{1}{1 \cdot r \cdot 1 - r} - \frac{1}{3 \cdot 1 + r \cdot 2 - r} + \frac{1}{5 \cdot 2 + r \cdot 3 - r} - \frac{1}{7 \cdot 3 + r \cdot 4 - r} \&c.$$

$s^{(m)}$  being the sine of  $2ra$ .

$$21. \frac{x + \sqrt{x^2 - 1}}{2} + \frac{x - \sqrt{x^2 - 1}}{2} \text{ and } \frac{x + \sqrt{x^2 - 1}}{2\sqrt{-1}} - \frac{x - \sqrt{x^2 - 1}}{2\sqrt{-1}}$$

being respectively equal to  $c^{(\rho)}$  and  $s^{(\rho)}$ , series whose sums are

the values of  $\frac{c^{(\rho)}}{c^{(\rho)}}$ ,  $\frac{s^{(\rho)}}{c^{(\rho)}}$ ,  $\frac{c^{(\rho)}}{s^{(\rho)}}$ , and  $\frac{s^{(\rho)}}{s^{(\rho)}}$  may be directly obtained by means of the binomial theorem: and by re-

collecting what is proved in art. 1. we may, by means of those series, deduce most of the theorems investigated in the preceding articles, and many others, without deriving them

them (by substitution) from other algebraic expressions as we have done above.

By the binomial theorem,

$$\frac{\binom{n}{s}}{2^s c^s} = \frac{x + \sqrt{x^2 - 1}}{2}^n + \frac{x - \sqrt{x^2 - 1}}{2}^n$$

$$\text{is } = \frac{x + \sqrt{x^2 - 1}}{2}^{\frac{1}{2}m} - e \cdot \frac{x + \sqrt{x^2 - 1}}{2}^{\frac{1}{2}m+n} + e'' \cdot \frac{x + \sqrt{x^2 - 1}}{2}^{\frac{1}{2}m+2n} \&c.$$

$$\text{or } = \frac{x - \sqrt{x^2 - 1}}{2}^{\frac{1}{2}m} - e \cdot \frac{x - \sqrt{x^2 - 1}}{2}^{\frac{1}{2}m+n} + e'' \cdot \frac{x - \sqrt{x^2 - 1}}{2}^{\frac{1}{2}m+2n} \&c.$$

whence, by multiplying the first series by  $\frac{x + \sqrt{x^2 - 1}}{2}^{m-\frac{1}{2}m}$  and the second by  $\frac{x - \sqrt{x^2 - 1}}{2}^{m-\frac{1}{2}m}$ , and subtracting one of the products from the other, we have, after dividing by  $2\sqrt{-1}$ ,

$$\frac{\binom{m-\frac{1}{2}m}{s}}{2^s c^s} = s - e \frac{\binom{m}{s}}{s} + e'' \frac{\binom{m+2n}{s}}{s} - e''' \frac{\binom{m+3n}{s}}{s} \&c.$$

And, by adding those products together, we have, after dividing by 2,

$$\frac{\binom{m-\frac{1}{2}m}{c}}{2^s c^s} = c - e \frac{\binom{m}{c}}{c} + e'' \frac{\binom{m+2n}{c}}{c} - e''' \frac{\binom{m+3n}{c}}{c} \&c.$$

$e''$ ,  $e'''$ , &c. being as in art. 17.

From these theorems, others (of considerable use in calculations) may be easily deduced in the following manner.

22. From the equation  $\frac{\binom{m-\frac{1}{2}m}{s}}{2^s c^s} = s - e \frac{\binom{m}{s}}{s} \&c.$  by multiplying by  $\frac{x}{\sqrt{1-x^2}}$  and taking the fluents according to what is proved in art. 1. we have

$$dG + fl. \frac{\overset{(m-1)n}{x}}{\sqrt{1-x^2}} \cdot \frac{\overset{(m-1)n}{s}}{\overset{(1n)}{2c}} = \frac{\overset{(m)}{c}}{m} - \frac{\overset{(m+n)}{c}}{m+n} + \frac{\overset{(m+2n)}{c}}{m+2n} \&c.$$

which is adjusted so that the fluent is supposed to begin when  $x$  is  $= \frac{2a}{n}$ ;  $x$  (the cosine of  $z$ ) being then the greatest root of the equation  $x + \sqrt{x^2 - 1}^{1n} + x - \sqrt{x^2 - 1}^{1n} = 0$ , whereof the other roots are the cosines of  $\frac{6a}{n}$ ,  $\frac{10a}{n}$ , &c. so long as these arcs are less than  $2a$ :  $d$  denoting the cosine of  $\frac{2ma}{n}$ , and  $G$  the sum of the series  $\frac{1}{m} + \frac{e}{m+n} + \frac{e''}{m+2n} \&c.$

which is equal to the *whole* fluent of  $\frac{x^{m-1} \dot{x}}{1-x^{2n}}$ , generated whilst  $x$  from 0 becomes equal to 1. For,  $c$  being  $= d$  when  $x$  is the cosine of  $\frac{2a}{n}$ , the cosines  $\overset{(m+n)}{c}$ ,  $\overset{(m+2n)}{c}$ ,  $\overset{(m+3n)}{c}$ , &c. will be equal to  $-d$ ,  $d$ ,  $-d$ , &c. respectively.

23. When the cosine  $x$  is  $= 1$ ,  $z$  is  $= 0$ , and  $\overset{(m)}{c}$ ,  $\overset{(m+n)}{c}$ ,  $\overset{(m+2n)}{c}$ , &c. each  $= 1$ : consequently, denoting the cosine of  $\frac{2a}{n}$  by  $e$ , and putting  $H$  for the sum of the series  $\frac{1}{m} - \frac{e}{m+n} + \frac{e''}{m+2n} - \frac{e'''}{m+3n} \&c.$  which is equal to the fluent of  $\frac{x^{m-1} \dot{x}}{1+x^{2n}}$ , generated whilst  $x$  from 0 becomes equal

to

to 1; the fluent of  $\frac{x^{\frac{m-1}{2}}}{\sqrt{1-x^2}} \cdot \frac{s}{\left(\frac{1}{2}\right)^c}$ , generated whilst  $x$  from  $c$  becomes equal to 1, will be = H - dG.

24. If  $m$  be =  $\frac{1}{2}en$ ,  $s$  will be = 0, and fl.  $\frac{x^{\frac{m-1}{2}}}{\sqrt{1-x^2}} \cdot \frac{s}{\left(\frac{1}{2}\right)^c}$

= 0; therefore it appears that dG will then be = H.

In this easy manner we discover, that

the cosine of  $ea$  is to radius,

as the series  $\frac{1}{e} - \frac{e}{e+2} + \frac{e''}{e+4} - \frac{e'''}{e+6} \&c.$

to the series  $\frac{1}{e} + \frac{e}{e+2} + \frac{e''}{e+4} + \frac{e'''}{e+6} \&c.$

that is, as the fluent of  $\frac{x^{\frac{1}{2}en-1} x}{1+x^n}$  generated whilst  $x$  from 0

becomes equal to 1, is to the fluent of  $\frac{x^{\frac{1}{2}en-1} x}{1-x^n}$  generated

in the same time. Which fluent of the first written fluxion is equal to *half* the *whole* fluent of the same fluxion, generated whilst  $x$  from 0 becomes infinite\*.

\* If  $y$  be =  $\frac{x}{x}$ ,  $\frac{x^{\frac{1}{2}en-1} x}{1+x^n}$  will be =  $-\frac{y^{\frac{1}{2}en-1} y}{1+y^n}$ ; therefore,  $y$  decreasing from 1 to 0 whilst  $x$  increases from 1 to infinity, it is evident that the fluent of  $\frac{x^{\frac{1}{2}en-1} x}{1+x^n}$ , generated whilst  $x$  from 1 becomes infinite, will be equal to the fluent of the same fluxion, generated whilst  $x$  from 0 becomes equal to 1.

N

25. If

25. If  $m - \frac{1}{2}en$  be  $= f$ ,  $e$  will be  $= \frac{2m-2f}{n}$ , and  $\frac{(m-\frac{1}{2}en)^s}{\sqrt{1-x^2}}$   
 $= fx^{f-1} + f'''x^{f-3} \sqrt{x^2-1} + f^{(5)}x^{f-5} \sqrt{x^2-1}^3 \&c. = 2x^{f-1}$   
 $- \frac{f-2}{2} \cdot 2x^{f-3} + \frac{f-3 \cdot f-4}{2} \cdot 2x^{f-5} - \frac{f-4 \cdot f-5 \cdot f-6}{2 \cdot 3} \cdot 2x^{f-7}$   
 $+ \frac{f-5 \cdot f-6 \cdot f-7 \cdot f-8}{2 \cdot 3 \cdot 4} \cdot 2x^{f-9} \&c.$  till the exponent of the  
 power of  $2x$  becomes 0 or 1: which series will both  
 terminate if  $f$  be a positive integer;  $f'''$  being  $= \frac{f \cdot f-1 \cdot f-2}{2 \cdot 3}$ ,  
 $f^{(5)} = \frac{f \cdot f-1 \cdot f-2 \cdot f-3 \cdot f-4}{2 \cdot 3 \cdot 4 \cdot 5}$ , &c. Therefore it follows, that

$$dG + \text{the fluent of } \frac{x \times 2x^{f-1} - f - 2 \cdot 2x^{f-3} \&c.}{x + \sqrt{x^2-1}^{\frac{1}{2}n} + x - \sqrt{x^2-1}^{\frac{1}{2}n}}$$

$$\text{will be } = \frac{c^{(m)}}{m} - \frac{c^{(m+n)}}{m+n} + \frac{c^{(m+2n)}}{m+2n} - \frac{c^{(m+3n)}}{m+3n} \&c.$$

the fluent being supposed to begin when  $x$  is equal to the  
 cosine of  $\frac{2a}{n}$ .

But if the fluent be supposed to begin when  $x$  is  $= 1$ ,

$$H + \text{the fluent of } \frac{x \times 2x^{f-2} - f - 2 \cdot 2x^{f-3} \&c.}{x + \sqrt{x^2-1}^{\frac{1}{2}n} + x - \sqrt{x^2-1}^{\frac{1}{2}n}}$$

$$\text{will be } = \frac{c^{(m)}}{m} - \frac{c^{(m+n)}}{m+n} + \frac{c^{(m+2n)}}{m+2n} - \frac{c^{(m+3n)}}{m+3n} \&c.$$

*Example 1.* If  $e$  be  $= 1$ ,  $m = 2$ , and  $n = 2$ ;  $f$  will be  
 $= 1$ , and

$\frac{1}{2}$  Log.

$$\frac{1}{2} \text{Log. } 2 + \frac{1}{2} \text{Log. } x = \frac{c''}{2} - \frac{c^{iv}}{4} + \frac{c^{vi}}{6} \&c.$$

$$\text{or Log. } x = c'' - \frac{c^{iv}}{2} + \frac{c^{vi}}{3} \&c. - \text{Log. } 2.$$

H, in this case, being  $= \frac{1}{2} \text{Log. } 2.$

*Example 2.*  $m$  being  $= e + f$  when  $n$  is  $= 2$ , we have

$$\begin{aligned} H + 2^{f-e-1} x^{f-e} \times \frac{1}{f-e} - \frac{f-2.2x}{f-e-2} \&c. - 2^{f-e-1} \times \frac{1}{f-e} - \frac{f-2.2^{-2}}{f-e-2} \&c. \\ = \frac{c^{(e+f)}}{e+f} - \frac{c^{(e+f+2)}}{e+f+2} + \frac{c^{(e+f+4)}}{e+f+4} \&c. \end{aligned}$$

$$\text{where H is } = \frac{1}{e+f} - \frac{e}{e+f+2} + \frac{e''}{e+f+4} \&c.$$

Hence, by taking  $x = 0$ , it appears that,  $f$  being greater than  $e$ ,

$$H \text{ is } = dG + 2^{f-e-1} \times \frac{1}{f-e} - \frac{f-2.2^{-2}}{f-e-2} \&c.$$

$d$  being here the cosine of  $e + f.a$ ,

$$\text{and } G = \frac{1}{e+f} + \frac{e}{e+f+2} + \frac{e''}{e+f+4} \&c.$$

26. From the theorem  $H + F = \frac{c^{(m)}}{m} - \frac{e^{(m+n)}}{m+n}$ , &c. in the preceding article, we get, by multiplying by  $\frac{-x}{\sqrt{1-x^2}}$  and taking the fluents,

$$Hx - fl. \frac{x F}{\sqrt{1-x^2}} = \frac{c^{(m)}}{m^2} - \frac{e^{(m+n)}}{(m+n)^2} + \frac{e^{(m+2n)}}{(m+2n)^2} \&c.$$



the fluent beginning when  $x$  is  $= 1$  and  $z = 0$ ; and  $F$

denoting the fluent of  $\frac{x \times 2x|^{f-1} - f - 2.2x|^{f-3} \&c.}{x + \sqrt{x^2 - 1}|^{kn} + x - \sqrt{x^2 - 1}|^{kn}}$ .

Likewise, by multiplying the last theorem by  $\frac{x}{\sqrt{1-x^2}}$  and taking the fluents again, we get

$$-\frac{1}{2}Hx^2 - \text{fl.} \frac{x}{\sqrt{1-x^2}} \text{fl.} \frac{x F}{\sqrt{1-x^2}} + Q_u = \frac{(m)}{m^3} - \frac{(m+n)}{m+n} + \frac{(m+2n)}{m+2n} \&c.$$

$Q_u$  denoting the series  $\frac{1}{m^2} - \frac{e}{m+n} + \frac{e''}{m+2n} \&c.$

And, by repeating the operation, other theorems may be obtained.

If we consider the fluents as beginning when  $x$  is  $= c$ , we, by the like operation, obtain, from the theorem

$$dG + F = \frac{(m)}{m} - \frac{(m+n)}{m+n} \&c. \text{ in the preceding article,}$$

$$dG.z - \frac{2a}{n} - \text{fl.} \frac{x F}{\sqrt{1-x^2}} + b Q_u = \frac{(m)}{m^2} - \frac{(m+n)}{m+n} + \frac{(m+2n)}{m+2n} \&c.$$

$Q_u$  denoting the series  $\frac{1}{m^2} + \frac{e}{m+n} + \frac{e''}{m+2n} \&c.$

and  $b$  the sine of  $\frac{2ma}{n}$ , of which  $d$  is the cosine.

And from hence, by proceeding in the same manner, other theorems may be deduced.

When  $x$  is  $= 1$ ,  $z$  is  $= 0$ ; and  $s, s, s, \&c.$  each  $= 0$ : therefore it appears that

$b Q_u$

$$bQ_u \text{ is } = \frac{2adG}{n} + \text{the fluent of } \frac{x^m F}{\sqrt{1-x^2}},$$

generated whilst  $x$  from  $c$  becomes equal to 1.

*Example 1.* If  $m$  be  $= \frac{1}{2}en$ ,  $f$  and  $F$  will each be  $= 0$ ;  
and consequently  $Q_u = \frac{2adG}{bn} = \frac{2aH}{bn}$ ;

$$\text{or } \frac{2b}{an} \times \frac{1}{e^2} + \frac{e}{e+2} + \frac{e^2}{e+4} \&c.$$

$$= dG (= d \times \text{the whole fluent of } \frac{x^{\frac{1}{2}en-1}x}{1-x^{2e}})$$

$$= H = \text{the fluent of } \frac{x^{\frac{1}{2}en-1}x}{1+x^{2e}},$$

generated whilst  $x$  from 0 becomes  $= 1$ ;  
which is a very useful theorem, for computing the values  
of such fluents in numbers.

Taking  $e$  and  $m$  each  $= \frac{1}{2}$ , and  $n = 2$ ;  $b$  will be  $= d$   
 $= \frac{1}{\sqrt{2}}$ : consequently we find

$$1 + \frac{1}{2 \cdot 5^2} + \frac{1 \cdot 3}{2 \cdot 4 \cdot 9^2} + \frac{1 \cdot 3 \cdot 5}{2 \cdot 4 \cdot 6 \cdot 13^2} \&c.$$

$$= \frac{1}{4} Q_u = \frac{aG}{4} = \frac{a}{4} \times E + \sqrt{E^2 - 2a};$$

$G$  being (in this case) the whole fluent of  $\frac{x^{-\frac{1}{2}}x}{\sqrt{1-x^2}}$ ; the  
value whereof, as is shewn in the Appendix, is  $E +$   
 $\sqrt{E^2 - 2a}$ ,  $E$  denoting the quadrantal arc of an ellipsis  
whose semi-axes are  $\sqrt{2}$  and 1.

*Example*

*Example 2.* If  $e$  be  $= \frac{1}{2}$ ,  $m = \frac{1}{2}$ , and  $n = 2$ ;  $b$  will be  $= \frac{1}{\sqrt{2}}$ ,  $d = -\frac{1}{\sqrt{2}}$ ,  $c = 0$ ,  $f = 1$ ,  $\dot{F} = \frac{x}{\sqrt{2x}}$ , and  $F = \sqrt{2x}$ , supposing the fluent to begin when  $x$  begins: consequently we have

$$\begin{aligned} & \frac{1}{3^2} + \frac{1}{2 \cdot 7^2} + \frac{1 \cdot 3}{2 \cdot 4 \cdot 11^2} + \frac{1 \cdot 3 \cdot 5}{2 \cdot 4 \cdot 6 \cdot 15^2} \text{ \&c.} \\ & = \frac{1}{4} Q'' = \frac{2F'' - aG}{4} = \frac{2-a}{4} \times E - \sqrt{E^2 - 2a}; \end{aligned}$$

$G$  being now the *whole* fluent of  $\frac{x^{\frac{1}{2}}}{\sqrt{1-x^2}}$ , and  $F''$  equal to the same fluent.

*Example 3.* Taking  $e = 1$ ,  $m = 2$ , and  $n = 2$ ; we have  $f = 1$ ,  $\dot{F} = \frac{x}{2x}$ , and

$$Hx = \frac{1}{2} \text{ fl. } \frac{x}{\sqrt{1-x^2}} \cdot \text{fl. } \frac{x}{x} = \frac{x''}{2^2} - \frac{x^{iv}}{4^2} + \frac{x^{vi}}{6^2} \text{ \&c.}$$

the fluent being supposed to begin when  $x$  is  $= 1$ , and  $H$  being  $= \frac{1}{2} \text{Log. } 2 =$  the fluent of  $\frac{x^2}{1+x^2}$  generated whilst  $x$  from  $0$  becomes equal to  $1$ .

$$\begin{aligned} \text{But fl. } \frac{x}{\sqrt{1-x^2}} \cdot \text{fl. } \frac{x}{x} \text{ is } & = -x \text{Log. } x - \text{fl. } \frac{x}{x} \cdot \text{fl. } \frac{x}{\sqrt{1-x^2}} \\ & = q'' + a - x \cdot \text{Log. } x - x - \frac{x^3}{2 \cdot 3^2} - \frac{1 \cdot 3 \cdot x^3}{2 \cdot 4 \cdot 5^2} - \frac{1 \cdot 3 \cdot 5 \cdot x^3}{2 \cdot 4 \cdot 6 \cdot 7^2} \text{ \&c.} \\ q'' \text{ being } & = 1 + \frac{1}{2 \cdot 3^2} + \frac{1 \cdot 3}{2 \cdot 4 \cdot 5^2} + \frac{1 \cdot 3 \cdot 5}{2 \cdot 4 \cdot 6 \cdot 7^2} \text{ \&c.} \end{aligned}$$

Therefore

$$\frac{1}{2}z \text{Log. } 2 - \frac{1}{2}a - z. \text{Log. } x - \frac{1}{2}q'' + \frac{1}{2} \times x + \frac{x^3}{2.3^2} \&c. \text{ is } = \frac{s''}{2^2} - \frac{s^{iv}}{4^2} \&c.$$

Hence, by taking  $x = 0$ , it appears that  $q''$  is  $= a \text{ Log. } 2$ ;  $x$  being then  $= a$ ,  $a - z. \text{Log. } x = 0$ , and  $s''$ ,  $s^{iv}$ , &c. each  $= 0$ . Consequently, after substituting for  $q''$  its value so found, it follows, that

$$x + \frac{x^3}{2.3^2} + \frac{1.3x^5}{2.4.5^2} \&c. \text{ is } = a - z. \text{Log. } 2x + \frac{1}{2} \times s'' - \frac{s^{iv}}{2^2} + \frac{s^{vi}}{3^2} \&c.$$

If  $x$  be  $= \frac{1}{\sqrt{2}}$ ,  $z$  will be  $= \frac{a}{2}$ ; and  $s''$ ,  $s^{iv}$ ,  $s^{vi}$ ,  $s^{viii}$ , &c. equal to  $1$ ,  $0$ ,  $-1$ ,  $0$ , &c. respectively: therefore it is evident, that

$$\sqrt{2} \times 1 + \frac{2^{-1}}{2.3^2} + \frac{1.3.2^{-3}}{2.4.5^2} \&c. \text{ is } = \frac{1}{2}a \text{Log. } 2 + 1 - \frac{1}{3^2} + \frac{1}{5^2} - \frac{1}{7^2} \&c.$$

Which last series being known to be equal to  $\text{fl. } \frac{x}{\sqrt{1-x^2}} \cdot \text{fl. } \frac{x}{\sqrt{1+x^2}}$ , or  $= \text{fl. } \frac{x}{x} \cdot \text{fl. } \frac{x}{1+x^2}$ , each generated whilst  $x$  from  $0$  becomes equal to  $1$ ; and the series  $1 + \frac{2^{-1}}{2.3^2} \&c.$  being equal to  $\frac{1}{2^2} \text{fl. } \frac{x}{x} \cdot \text{fl. } \frac{x}{\sqrt{2x-x^2}}$ , generated in the same time; it follows, that

$$\text{fl. } \frac{x}{\sqrt{1-x^2}} \cdot \text{Log. } x + \sqrt{1+x^2} \text{ is } = \text{fl. } \frac{ix}{x} = \frac{1}{2} \text{fl. } \frac{vx}{x} - \frac{1}{2}a \text{Log. } 2,$$

$$\text{and fl. } \frac{v-2t}{x} \cdot \frac{x}{x} = a \text{Log. } 2 = \text{fl. } \frac{a-z}{x} \cdot \frac{x}{x};$$

these fluents being all generated in the time before mentioned,

tioned, and  $t$  denoting the circular arc whose radius is 1 and tangent  $x$ , and  $v$  an arc (of the same circle) whose versed sine is  $x$ .

*Example 4.* If  $e$  be  $= \frac{1}{2}$ ,  $m = 2$ , and  $n = 4$ ;  $b$  will be  $= 1$ ,  $d = 0$ ,  $c = \frac{1}{\sqrt{2}}$ ,  $f = 1$ , and  $F = \frac{x}{2^{\frac{1}{2}}\sqrt{2x^2-1}}$ : it follows therefore, that

$2^{\frac{1}{2}}$  fl.  $\frac{x}{\sqrt{1-x^2}}$  fl.  $\frac{x}{\sqrt{2x^2-1}}$ , generated whilst  $x$  from  $c$  becomes  $= 1$ ,

$$is = 1 + \frac{1}{2 \cdot 3} + \frac{1 \cdot 3}{2 \cdot 4 \cdot 5} + \frac{1 \cdot 3 \cdot 5}{2 \cdot 4 \cdot 6 \cdot 7} \&c.$$

which, by the preceding example, is  $= e$  Log. 2.

*Example 5.* If  $m$  be  $= e$  and  $n = 2$ ,  $f$  and  $F$  will each be  $= 0$ ; therefore, taking  $x$  equal to  $a$ , we have

$$Q_u = \frac{1}{e^2} - \frac{e}{e+2} + \frac{e^2}{e+4} \&c. = \frac{1}{2} a^2 H + d \times \frac{1}{e^2} + \frac{e}{e+2} + \frac{e^2}{e+4} \&c.$$

$d$  being the cosine of  $ea$ .

and  $H =$  fl.  $\frac{x^{e-1} x}{1+x^2}$ , generated whilst  $x$  from 0 becomes  $= a$ :

which fluent, by art. 24, is  $= d \times$  the whole fluent of  $\frac{x^{e-1} x}{1+x^2}$ .

When  $e$  is  $= 1$ ,  $d$  is  $= 0$ , and  $H = \frac{1}{2} a$ ; and we have, in that case,  $1 - \frac{1}{3} + \frac{1}{5} - \frac{1}{7} \&c. = \frac{1}{2} a$ , as in art. 8.

27. From

27. From the equation  $\frac{c^{(m-\frac{1}{2}en)}}{(2c)^{\frac{1}{2}en}} = c - e \frac{(m+n)}{c}$  &c. in art.

21. by multiplying by  $\frac{x^{\frac{1}{2}en}}{\sqrt{1-x^2}}$  and taking the fluents according to what is shewn in art. 1. we get

$$bG - \text{fl.} \frac{x^{\frac{1}{2}en}}{\sqrt{1-x^2}} \cdot \frac{c^{(m-\frac{1}{2}en)}}{(2c)^{\frac{1}{2}en}} = \frac{s^{(m)}}{m} - \frac{s^{(m+n)}}{m+n} + \frac{s^{(m+2n)}}{m+2n} \&c.$$

the fluent being supposed to begin when  $x$  is ( $= c$ ) the cosine of  $\frac{2a}{n}$ ; the sines  $s, s, s, s, \&c.$  being then equal to  $b, -b, b, -b, \&c.$  respectively.

When  $x$  is  $= 1$ , the sines  $s, s, s, \&c.$  are each

$$= 0: \text{ consequently } bG \text{ is } = \text{fl.} \frac{x^{\frac{1}{2}en}}{\sqrt{1-x^2}} \cdot \frac{c^{(m-\frac{1}{2}en)}}{(2c)^{\frac{1}{2}en}}, \text{ generated}$$

whilst  $x$  from  $c$  becomes equal to 1.

If the fluent be considered as beginning when  $x$  is  $= 1$ , the theorem will stand (without the term  $bG$ )

$$- \text{fl.} \frac{x^{\frac{1}{2}en}}{\sqrt{1-x^2}} \cdot \frac{c^{(m-\frac{1}{2}en)}}{(2c)^{\frac{1}{2}en}} = \frac{s^{(m)}}{m} - \frac{s^{(m+n)}}{m+n} + \frac{s^{(m+2n)}}{m+2n} \&c.$$

*Example 1.* Taking  $e = 1$ , and  $m$  and  $n$  each equal 2, we have  $x = s'' = \frac{s^{17}}{2} + \frac{s^{11}}{3} \&c.$  when  $x$  is positive: from

⊙

whence,

whence, by multiplying by  $\frac{x}{\sqrt{1-x^2}}$  and taking the fluents, other theorems may be easily deduced, as in art. 6. and 7.

*Example 2.* If  $e$  be  $= \frac{1}{3}$ ,  $m = 1$ , and  $n = 6$ ;  $b$  will be  $= \frac{1}{2}$ ,  $c = \frac{\sqrt{3}}{2}$ , and

$$\frac{1}{2}G = \text{the fluent of } \frac{2^{-\frac{1}{2}}x^{-\frac{1}{2}}}{4x^2 - 3)^{\frac{1}{2}} \times \sqrt{1-x^2}}$$

$$= s' - \frac{s^{vii}}{3 \cdot 7} + \frac{4s^{xiii}}{3 \cdot 6 \cdot 13} - \frac{4 \cdot 7s^{xix}}{3 \cdot 6 \cdot 9 \cdot 19} \&c.$$

$G$  denoting the whole fluent of  $\frac{x}{1-x^2}$ , or of  $\frac{xy^{-\frac{1}{2}}}{1-y^2}$  generated whilst  $x$ , or  $y$ , from 0 becomes equal to 1: the value whereof is assigned, in Table IV. of the Appendix, by means of a *circular* and an *elliptic* arc.

The series  $s' - \frac{s^{vii}}{3 \cdot 7} \&c.$  vanishing when  $x$  is = 1, the fluent of  $\frac{2^{-\frac{1}{2}}x^{-\frac{1}{2}}}{4x^2 - 3)^{\frac{1}{2}} \times \sqrt{1-x^2}}$ , generated whilst  $x$  from  $\frac{\sqrt{3}}{2}$  becomes equal to 1, is  $= \frac{1}{2}G$ .

*Example 3.* If  $m$  be  $= e$  and  $n = 2$ ,

$$bG = \text{fl. } \frac{x}{2^e x^e \sqrt{1-x^2}} \text{ will be } = \frac{s}{e} - \frac{s^{e+2}}{e+2} + \frac{s^{e+4}}{e+4} \&c.$$

and the *whole* fluent of  $\frac{x}{2^e x^e \sqrt{1-x^2}} = bG = b \times \text{the whole flu. of } \frac{x^{e-1} x}{1-x^2}$

$$= \frac{b}{d}$$

$= \frac{b}{d} \times$  the contemporary fluent of  $\frac{x^{e-1} \dot{x}}{1+x^2}^e$ , by art. 24.

$b$  and  $d$  being the sine and cosine of  $ea$ .

*Example 4.* Taking  $m = e + 1$  and  $n = 2$ , we have

$$bG - \text{fl.} \frac{x^{1-e} \dot{x}}{2^e \sqrt{1-x^2}} = \frac{x^{e+1}}{e+1} - \frac{x^{e+3}}{e+3} + \frac{x^{e+5}}{e+5} \&c.$$

and  $bG =$  the whole fluent of  $\frac{x^{1-e} \dot{x}}{2^e \sqrt{1-x^2}}$ .

Moreover,  $G$ , by Ex. 2. art. 25. being  $= \frac{1}{d} \times H - \frac{1}{2^e \cdot 1 - e}$ ,

it appears that

$\frac{b}{d} \times H - \frac{1}{2^e \cdot 1 - e}$  is also = the same whole fluent;

$b$  and  $d$  being here the sine and cosine of  $e + 1.a$ ,

$\therefore G =$  the whole fluent of  $\frac{x^e \dot{x}}{1-x^2}^e$

$=$  the series  $\frac{1}{e+1} + \frac{e}{e+3} + \frac{e''}{e+5} + \frac{e'''}{e+7} \&c.$

and  $H =$  the contemporary fluent of  $\frac{x^e \dot{x}}{1+x^2}^e$

$=$  the series  $\frac{1}{e+1} - \frac{e}{e+3} + \frac{e''}{e+5} - \frac{e'''}{e+7} \&c.$

Other theorems may be deduced from the general theorem in this article, by multiplying by  $\frac{x}{\sqrt{1-x^2}}$  and taking the fluents, as in the preceding article.



28. By division we have

$$\frac{1}{x - \sqrt{x^2 - 1} + y \cdot x + \sqrt{x^2 - 1}} = \frac{1}{x + \sqrt{x^2 - 1}} - \frac{y \cdot x + \sqrt{x^2 - 1}}{x + \sqrt{x^2 - 1}} + \frac{y^2 \cdot x + \sqrt{x^2 - 1}}{x + \sqrt{x^2 - 1}} \&c.$$

$$\text{and } \frac{1}{x + \sqrt{x^2 - 1} + y \cdot x - \sqrt{x^2 - 1}} = \frac{1}{x - \sqrt{x^2 - 1}} - \frac{y \cdot x - \sqrt{x^2 - 1}}{x - \sqrt{x^2 - 1}} + \frac{y^2 \cdot x - \sqrt{x^2 - 1}}{x - \sqrt{x^2 - 1}} \&c.$$

whence, by subtracting one equation from the other, and dividing by  $2\sqrt{-1}$ , we get

$$\frac{1}{1 - y} \cdot \frac{\binom{2n}{n}}{1 + y^2 + 2yc} = s - y \cdot s + y^2 \cdot s \&c.$$

and, by taking  $n$  equal to 2,

$$\frac{1}{1 - y} \cdot \frac{1 - x^2}{1 - y^2 + 4yx^2} = s - y s'' + y^2 s' \&c.$$

Hence, by multiplying by  $\frac{x}{\sqrt{1 - x^2}}$ , taking the fluents, and writing  $y^2$  instead of  $y$ , we find

$$\frac{1}{2} \text{ Circ. Arc, rad. } 1, \text{ tang. } \frac{2xy}{1 - y^2} = c'y - \frac{c''y^2}{3} + \frac{c''''y^4}{5} \&c.$$

where both  $x$  and  $y$  may be considered as variable, independent of each other.

*Example 1.* Taking  $x = \frac{1}{\sqrt{2}}$ , we have  $c', c''', c^2, c^{21}$ , &c. equal to  $\frac{1}{\sqrt{2}}, -\frac{1}{\sqrt{2}}, -\frac{1}{\sqrt{2}}, \frac{1}{\sqrt{2}}$ , &c. respectively; and

$$\frac{1}{\sqrt{2}}$$

$$\frac{x}{\sqrt{2}} \text{ Circ. Arc, rad. 1, tang. } \frac{2^{\frac{1}{2}}y}{1-y^2} = y + \frac{y^3}{3} - \frac{y^5}{5} + \frac{y^7}{7} \&c.$$

$$= \text{fl. } \frac{y}{1+y^2} + \text{fl. } \frac{y^3}{1+y^2}$$

Moreover, the sum of these two fluents, generated whilst  $y$  from 0 becomes equal to 1, being, by art. 17. equal to the *whole* fluent of  $\frac{y}{1+y^2}$ , or of  $\frac{y^3}{1+y^2}$ , generated whilst  $y$  from 0 becomes infinite; it appears that each of these *whole* fluents is  $= \frac{a}{\sqrt{2}}$ .

*Example 2.* If  $x$  be  $= \frac{1}{2}$ ;  $e, e^m, e^r, e^{11}, e^{12}, \&c.$  will be equal to  $\frac{1}{2}, -1, \frac{1}{2}, \frac{1}{2}, -1, \&c.$  respectively; and

$$\text{Circ. Arc, rad. 1, tang. } \frac{y}{1-y^2} = y + \frac{2y^3}{3} + \frac{y^5}{5} - \frac{y^7}{7} + \frac{2y^9}{9} - \frac{y^{11}}{11} \&c.$$

$$= \text{fl. } \frac{y}{1+y^2} + 2 \text{ fl. } \frac{y^3}{1+y^2} + \text{fl. } \frac{y^5}{1+y^2} :$$

or,  $\text{fl. } \frac{y^3}{1+y^2}$  being  $= \frac{1}{3} \text{ Circ. Arc, rad. 1, tang. } y^3$ ,

$$\text{Circ. Arc, rad. 1, tang. } \frac{y}{1-y^2} = \frac{1}{3} \text{ Circ. Arc, rad. 1, tang. } y^3$$

$$= y - \frac{y^7}{7} + \frac{y^{13}}{13} - \frac{y^{19}}{19} \&c. + \frac{y^3}{5} - \frac{y^{11}}{11} + \frac{y^{17}}{17} - \frac{y^{23}}{23} \&c.$$

$$= \text{fl. } \frac{y}{1+y^2} + \text{fl. } \frac{y^3}{1+y^2}$$

Moreover, the sum of these two last written fluents, generated whilst  $y$  from 0 becomes equal to 1, being, by art. 17. equal to the *whole* fluent of  $\frac{y}{1+y^2}$ , or of  $\frac{y^3}{1+y^2}$ , generated whilst  $y$  from 0 becomes infinite; it follows, that each of these *whole* fluents is  $= \frac{2a}{3}$ .

29. Adding together the first two equations in the preceding article, and dividing by 2, we have

$$\frac{1 + y \cdot c}{1 + y^2 + 2yc} \binom{n}{2} = c - c \cdot y + c \cdot y^2 \text{ \&c.}$$

and, by taking  $n$  equal to 2,

$$\frac{1 + y \cdot x}{1 - y^2 + 4yx^2} = c - c'''y + c'y^2 \text{ \&c.}$$

Hence, by multiplying by  $\frac{-x}{\sqrt{1-x^2}}$ , taking the fluents, and writing  $y^2$  instead of  $y$ , we have

$$\frac{1}{4} \text{Log.} \frac{1 + y^2 + 2y\sqrt{1-x^2}}{1 + y^2 - 2y\sqrt{1-x^2}} = s'y - \frac{s''y^3}{3} + \frac{s'y^5}{5} \text{ \&c.}$$

*Example 1.* Taking  $x = \frac{1}{\sqrt{2}}$ , we have  $s', s''', s^v, s^{vii}$ , &c. equal to  $\frac{1}{\sqrt{2}}, \frac{1}{\sqrt{2}}, -\frac{1}{\sqrt{2}}, -\frac{1}{\sqrt{2}}$ , &c. respectively, and

$$\begin{aligned} \frac{1}{2^{\frac{3}{2}}} \text{Log.} \frac{1 + y^2 + 2^{\frac{1}{2}}y}{1 + y^2 - 2^{\frac{1}{2}}y} &= y - \frac{y^3}{3} - \frac{y^5}{5} + \frac{y^7}{7} \text{ \&c.} \\ &= \text{fl.} \frac{y}{1 + y^4} - \text{fl.} \frac{y^2y}{1 + y^4}. \end{aligned}$$

But we have found, in Ex. 1. of the preceding article, that the sum of these two last written fluents is  $\frac{A}{\sqrt{2}}$ . Consequently we find

$$\text{fl.} \frac{y}{1 + y^4} = \frac{A + L}{2^{\frac{3}{2}}}; \text{ and fl.} \frac{y^2y}{1 + y^4} = \frac{A - L}{2^{\frac{3}{2}}};$$

A de-

A denoting the Circ. Arc, rad. 1, tang.  $\frac{2^{\frac{1}{2}}y}{1-y^2}$ .

and L the Log. of  $\frac{1+y^2+2^{\frac{1}{2}}y}{1+y^2-2^{\frac{1}{2}}y}$  = Log.  $\frac{1+y^2+2^{\frac{1}{2}}y}{\sqrt{1+y^2}}$ .

Example 2. If  $x$  be  $= \frac{1}{2}$ ;  $s', s'', s', s''$ , will be equal to  $\frac{\sqrt{3}}{2}$ , 0,  $-\frac{\sqrt{3}}{2}$ ,  $\frac{\sqrt{3}}{2}$ , &c. respectively; and

$$\frac{1}{2\sqrt{3}} \text{Log.} \frac{1+y^2+3^{\frac{1}{2}}y}{1+y^2-3^{\frac{1}{2}}y} = y - \frac{y^3}{5} - \frac{y^7}{7} + \frac{y^{11}}{11} \text{ \&c.}$$

$$= \text{fl.} \frac{j}{1+y^6} - \text{fl.} \frac{y^7j}{1+y^6}.$$

Hence, and from Ex. 2. of the preceding article, we have

$$\text{fl.} \frac{j}{1+y^6} = \frac{A}{2} - \frac{A''}{3} + \frac{L}{4\sqrt{3}}, \text{ and fl.} \frac{y^7j}{1+y^6} = \frac{A}{2} - \frac{A''}{3} - \frac{L}{4\sqrt{3}};$$

A being = Circ. Arc, rad. 1, tang.  $\frac{y}{1-y^2}$ ; A'' = Circ. Arc,

rad. 1, tang.  $y^3$ ; and L = Log.  $\frac{1+y^2+3^{\frac{1}{2}}y}{1+y^2-3^{\frac{1}{2}}y}$ .

30. Multiplying the first equation in art. 28. by

$\frac{x+\sqrt{x^2-1}}{x-\sqrt{x^2-1}} \frac{n}{2}$ , and the second by  $\frac{x-\sqrt{x^2-1}}{x+\sqrt{x^2-1}} \frac{n}{2}$ , we have

$$\frac{\frac{x+\sqrt{x^2-1}}{x-\sqrt{x^2-1}} \frac{n}{2}}{\frac{x-\sqrt{x^2-1}}{x+\sqrt{x^2-1}} \frac{n}{2} + y \frac{x+\sqrt{x^2-1}}{x-\sqrt{x^2-1}} \frac{n}{2}} = \frac{x+\sqrt{x^2-1}}{x-\sqrt{x^2-1}} \frac{n}{2} - y \frac{x+\sqrt{x^2-1}}{x-\sqrt{x^2-1}} \frac{n}{2} \text{ \&c.}$$

$$\text{and } \frac{\frac{x-\sqrt{x^2-1}}{x+\sqrt{x^2-1}} \frac{n}{2}}{\frac{x+\sqrt{x^2-1}}{x-\sqrt{x^2-1}} \frac{n}{2} + y \frac{x-\sqrt{x^2-1}}{x+\sqrt{x^2-1}} \frac{n}{2}} = \frac{x-\sqrt{x^2-1}}{x+\sqrt{x^2-1}} \frac{n}{2} - y \frac{x-\sqrt{x^2-1}}{x+\sqrt{x^2-1}} \frac{n}{2} \text{ \&c.}$$

Whence,

Whence, by subtracting one equation from the other, and dividing by  $2\sqrt{-1}$ , we get

$$\frac{\binom{n}{1}}{1+y^2+2yc} = \binom{n}{1} s - \binom{2n}{2} s y + \binom{3n}{3} s y^2 - \binom{4n}{4} s y^3 \&c.$$

Hence, by multiplying by  $\frac{x}{\sqrt{1-x^2}}$  and taking the fluents, we, upon adjusting the equation, find

$$\frac{1}{2} \text{Log. } 1+y^2+2yc = c y - \frac{\binom{2n}{2} c y^2}{2} + \frac{\binom{3n}{3} c y^3}{3} - \frac{\binom{4n}{4} c y^4}{4} \&c.$$

31. Writing  $-y$  instead of  $y$ , we have (from the last theorem)

$$\frac{1}{2} \text{Log. } 1+y^2-2yc = -c y - \frac{\binom{2n}{2} c y^2}{2} - \frac{\binom{3n}{3} c y^3}{3} \&c.$$

and consequently (from the last two theorems)

$$\frac{1}{4} \text{Log. } \frac{1+y^2+2yc}{1+y^2-2yc} = c y + \frac{\binom{3n}{3} c y^2}{3} + \frac{\binom{5n}{5} c y^4}{5} \&c.$$

$$\text{and } \frac{1}{4} \text{Log. } \frac{1}{(1+y^2)^2 - 4y^2 c} = \frac{\binom{2n}{2} c y^2}{2} + \frac{\binom{4n}{4} c y^4}{4} + \frac{\binom{6n}{6} c y^6}{6} \&c.$$

*Example.* If  $x$  be  $= \frac{1}{2}$ ;  $c'$ ,  $c''$ ,  $c'''$ , &c. will be as in Ex. 2. art. 28. and therefore, from the last theorem but one,  $n$  being therein taken equal to 1, we have

$$\frac{1}{2} \text{Log. } \frac{1+y^2+y}{1+y^2-y} = y - \frac{2y^3}{3} + \frac{y^5}{5} + \frac{y^7}{7} - \frac{2y^9}{9} + \frac{y^{11}}{11} \&c.$$

$= \text{fl.}$

$$= \text{fl. } \frac{y}{1-y^6} - 2 \text{ fl. } \frac{y^2}{1-y^6} + \text{fl. } \frac{y^4}{1-y^6} :$$

$$\text{or, fl. } \frac{y^2}{1-y^6} \text{ being } = \frac{1}{6} \text{ Log. } \frac{1+y^3}{1-y^3},$$

$$\begin{aligned} & \frac{1}{6} \text{ Log. } \frac{1+y^2+y}{1+y^2-y} + \frac{1}{6} \text{ Log. } \frac{1+y^3}{1-y^3} \\ & = y + \frac{y^7}{7} \&c. + \frac{y^5}{5} + \frac{y^{11}}{11} \&c. = \text{fl. } \frac{y}{1-y^6} + \text{fl. } \frac{y^2}{1-y^6}. \end{aligned}$$

32. Taking  $n$  equal to 1, we have from the preceding article and art. 28.

$$\begin{aligned} & \frac{1}{8} \text{ Log. } \frac{1+y^2+2xy}{1+y^2-2xy} + \frac{1}{4} \text{ Circ. Arc, rad. } 1, \text{ tang. } \frac{2xy}{1-y^2} \\ & = c'y + \frac{c^3y^3}{5} + \frac{c^{13}y^9}{9} \&c. \end{aligned}$$

$$\begin{aligned} \text{and } \frac{1}{8} \text{ Log. } \frac{1+y^2+2xy}{1+y^2-2xy} - \frac{1}{4} \text{ Circ. Arc, rad. } 1, \text{ tang. } \frac{2xy}{1-y^2} \\ & = \frac{c''y^3}{3} + \frac{c^{11}y^7}{7} + \frac{c^{21}y^{11}}{11} \&c. \end{aligned}$$

*Example 1.* Taking  $x$  equal to 1, we have

$$\begin{aligned} & \frac{1}{4} \text{ Log. } \frac{1+y}{1-y} + \frac{1}{2} \text{ Circ. Arc, rad. } 1, \text{ tang. } y \\ & = y + \frac{y^3}{5} + \frac{y^9}{9} \&c. = \text{fl. } \frac{y}{1-y^4}; \end{aligned}$$

$$\begin{aligned} \text{and } \frac{1}{4} \text{ Log. } \frac{1+y}{1-y} - \frac{1}{2} \text{ Circ. Arc, rad. } 1, \text{ tang. } y \\ & = \frac{y^3}{3} + \frac{y^7}{7} + \frac{y^{11}}{11} \&c. = \text{fl. } \frac{y^2}{1-y^4}; \end{aligned}$$

Circ. Arc, rad. 1, tang.  $\frac{2y}{1-y^2}$  being = 2 Circ. Arc, rad. 1, tang.  $y$ .

*Example 2.* If  $x$  be =  $\frac{1}{2}$ ;  $c'$ ,  $c''$ ,  $c'$ , &c. will be as in Ex. 2. art. 28. and

P

 $\frac{1}{4}$  Log.

$$\begin{aligned} \frac{1}{4} \text{Log. } \frac{1+y^2+y}{1+y^2-y} + \frac{1}{2} \text{Circ. Arc, rad. 1, tang. } \frac{y}{1-y^2} \\ = y + \frac{y^5}{5} - \frac{2y^9}{9} + \frac{y^{13}}{13} + \frac{y^{17}}{17} - \frac{2y^{21}}{21} \&c. \\ = \text{fl. } \frac{y}{1-y^{12}} + \text{fl. } \frac{y^5}{1-y^{12}} - 2 \text{ fl. } \frac{y^9}{1-y^{12}}; \end{aligned}$$

$$\begin{aligned} \text{and } \frac{1}{4} \text{Log. } \frac{1+y^2+y}{1+y^2-y} - \frac{1}{2} \text{Circ. Arc, rad. 1, tang. } \frac{y}{1-y^2} \\ = -\frac{2y^3}{3} + \frac{y^7}{7} + \frac{y^{11}}{11} - \frac{2y^{15}}{15} + \frac{y^{19}}{19} + \frac{y^{23}}{23} \&c. \\ = -2 \text{ fl. } \frac{y^3}{1-y^{12}} + \text{fl. } \frac{y^7}{1-y^{12}} + \text{fl. } \frac{y^{11}}{1-y^{12}}. \end{aligned}$$

But by Ex. 1. fl.  $\frac{y^3}{1-y^{12}}$  and fl.  $\frac{y^7}{1-y^{12}}$  are equal to

$$\frac{1}{12} \text{Log. } \frac{1+y^3}{1-y^3} + \frac{1}{6} \text{Circ. Arc, rad. 1, tang. } y^3,$$

and  $\frac{1}{12} \text{Log. } \frac{1+y^3}{1-y^3} - \frac{1}{6} \text{Circ. Arc, rad. 1, tang. } y^3$ , respectively.

Consequently we have

$$\begin{aligned} \frac{1}{4} \text{Log. } \frac{1+y^2+y}{1+y^2-y} + \frac{1}{2} \text{Log. } \frac{1+y^3}{1-y^3} + \frac{1}{2} \text{Circ. Arc, rad. 1, tang. } \frac{y}{1-y^2} \\ - \frac{1}{2} \text{Circ. Arc, rad. 1, tang. } y^3 = \text{fl. } \frac{y}{1-y^{12}} + \text{fl. } \frac{y^5}{1-y^{12}}; \end{aligned}$$

$$\begin{aligned} \text{and } \frac{1}{4} \text{Log. } \frac{1+y^2+y}{1+y^2-y} + \frac{1}{2} \text{Log. } \frac{1+y^3}{1-y^3} - \frac{1}{2} \text{Circ. Arc, rad. 1, tang. } \frac{y}{1-y^2} \\ + \frac{1}{2} \text{Circ. Arc, rad. 1, tang. } y^3 = \text{fl. } \frac{y^3}{1-y^{12}} + \text{fl. } \frac{y^7}{1-y^{12}}. \end{aligned}$$

33. Adding together the first two equations in Art. 30. and dividing by 2, we have

$$\frac{y^3 + c^{(n)}}{1+y^2+2yc} = c^{(n)} - y c^{(2n)} + y^2 c^{(3n)} - y^3 c^{(4n)} \&c.$$

Hence,

Hence, by taking  $x$  equal to 1, multiplying by  $\frac{-y^2}{\sqrt{1-x^2}}$  and taking the fluents, we find

$$\begin{aligned} & \frac{x}{2} - \text{Circ. Arc. rad. 1, tang. } \frac{1-y}{1+y} \times \left. \frac{1-x}{1+x} \right|^{\frac{1}{2}} \\ & = \text{Circ. Arc, rad. 1, tang. } \frac{y\sqrt{1-x^2}}{xy+1} = s'y - \frac{s''y^2}{2} + \frac{s'''y^3}{3} \&c. \end{aligned}$$

34. Writing  $-y$  instead of  $y$ , we have

$$\begin{aligned} & \frac{x}{2} - \text{Circ. Arc, rad. 1, tang. } \frac{1+y}{1-y} \times \left. \frac{1-x}{1+x} \right|^{\frac{1}{2}} \\ & = \text{Circ. Arc, rad. 1, tang. } \frac{y\sqrt{1-x^2}}{xy-1} = -s'y - \frac{s''y^2}{2} - \frac{s'''y^3}{3} \&c. \end{aligned}$$

and consequently (from the last two theorems)

$$\frac{1}{4} \text{Circ. Arc, rad. 1, tang. } \frac{2y}{1-y^2} \sqrt{1-x^2} = s'y + \frac{s''y^2}{3} + \frac{s'''y^3}{5} \&c.$$

and

$$\frac{1}{2} \text{Circ. Arc, rad. 1, tang. } \frac{2xy^2}{1+y^2-2x^2y^2} \sqrt{1-x^2} = \frac{s'y^2}{2} + \frac{s''y^4}{4} + \frac{s'''y^6}{6} \&c.$$

*Example.* If  $x$  be  $= \frac{1}{2}$ ;  $s'$ ,  $s''$ ,  $s'''$ , &c. will be equal to  $\frac{\sqrt{3}}{2}$ , 0,  $-\frac{\sqrt{3}}{2}$ , &c. respectively; and

$$\begin{aligned} & \frac{1}{3^{\frac{1}{2}}} \text{Circ. Arc, rad. 1, tang. } \frac{3^{\frac{1}{2}}y}{1-y^2} \\ & = y - \frac{y^3}{5} + \frac{y^7}{7} - \frac{y^{11}}{11} \&c. = \text{fl. } \frac{y}{1-y^5} - \text{fl. } \frac{y^4}{1-y^5}. \end{aligned}$$

Hence, and from art. 31. we have  $\text{fl. } \frac{y}{1-y^5} =$

$$\frac{1}{4} \text{Log. } \frac{1+y^2+y}{1+y^2-y} + \frac{1}{8} \text{Log. } \frac{1+y^2}{1-y^2} + \frac{1}{2\sqrt{3}} \text{Circ. Arc, rad. 1, tang. } \frac{3^{\frac{1}{2}}y}{1-y^2};$$



$$\text{and fl. } \frac{y^4 j}{1-y^6} =$$

$$\frac{1}{4} \text{Log. } \frac{1+y^2+y}{1+y^2-y} + \frac{1}{6} \text{Log. } \frac{1+y^2}{1-y^2} - \frac{1}{2\sqrt{3}} \text{Circ. Arc, rad. 1, tang. } \frac{3^{\frac{1}{2}}y}{1-y^2}$$

In the same manner, from the 3d theorem in art. 31. and the 3d theorem in this article, we find fl.  $\frac{y^2 j}{1-y^6} =$

$$\frac{1}{4} \text{Log. } 1+y^2+y^4 - \frac{1}{6} \text{Log. } 1-y^6 + \frac{1}{2\sqrt{3}} \text{Circ. Arc, rad. 1, tang. } \frac{3^{\frac{1}{2}}y^2}{2+y^2};$$

$$\text{and fl. } \frac{y^3 j}{1-y^6} =$$

$$\frac{1}{4} \text{Log. } 1+y^2+y^4 - \frac{1}{6} \text{Log. } 1-y^6 - \frac{1}{2\sqrt{3}} \text{Circ. Arc, rad. 1, tang. } \frac{3^{\frac{1}{2}}y^2}{2+y^2};$$

$$\text{or fl. } \frac{j}{1-y^3} =$$

$$\frac{1}{2} \text{Log. } 1+y+y^2 - \frac{1}{3} \text{Log. } 1-y^3 + \frac{1}{\sqrt{3}} \text{Circ. Arc, rad. 1, tang. } \frac{3^{\frac{1}{2}}y}{2+y};$$

$$\text{and fl. } \frac{y j}{1-y^3} =$$

$$\frac{1}{2} \text{Log. } 1+y+y^2 - \frac{1}{3} \text{Log. } 1-y^3 - \frac{1}{\sqrt{3}} \text{Circ. Arc, rad. 1, tang. } \frac{3^{\frac{1}{2}}y}{2+y}.$$

35. From the preceding article and art. 29. we have

$$\begin{aligned} \frac{1}{4} \text{Circ. Arc, rad. 1, tang. } \frac{2y}{1-y^2} \sqrt{1-x^2} + \frac{1}{8} \text{Log. } \frac{1+y^2+2y\sqrt{1-x^2}}{1+y^2-2y\sqrt{1-x^2}} \\ = s'y + \frac{s'y^3}{5} + \frac{s^3y^9}{9} \&c. \end{aligned}$$

$$\begin{aligned} \text{and } \frac{1}{4} \text{Circ. Arc, rad. 1, tang. } \frac{2y}{1-y^2} \sqrt{1-x^2} - \frac{1}{8} \text{Log. } \frac{1+y^2+2y\sqrt{1-x^2}}{1+y^2-2y\sqrt{1-x^2}} \\ = \frac{s''y^3}{3} + \frac{s^{11}y^7}{7} + \frac{s^{21}y^{11}}{11} \&c. \end{aligned}$$

*Example.*

*Example.* Taking  $x$  equal to  $\frac{1}{2}$ , we have  $s'$ ,  $s''$ ,  $s'$ , &c. as in the preceding example; and

$$\frac{1}{2\sqrt{3}} \text{Circ. Arc, rad. 1, tang. } \frac{3^{\frac{1}{2}}y}{1-y^2} + \frac{1}{4\sqrt{3}} \text{Log. } \frac{1+y^2+3^{\frac{1}{2}}y}{1+y^2-3^{\frac{1}{2}}y}$$

$$= y - \frac{y^5}{5} + \frac{y^{13}}{13} - \frac{y^{17}}{17} \text{ \&c.} = \text{fl. } \frac{y}{1-y^{12}} - \text{fl. } \frac{y^4j}{1-y^{12}};$$

and  $\frac{1}{2\sqrt{3}} \text{Circ. Arc, rad. 1, tang. } \frac{3^{\frac{1}{2}}y}{1-y^2} - \frac{1}{4\sqrt{3}} \text{Log. } \frac{1+y^2+3^{\frac{1}{2}}y}{1+y^2-3^{\frac{1}{2}}y}$

$$= \frac{y^7}{7} - \frac{y^{11}}{11} + \frac{y^{19}}{19} - \frac{y^{23}}{23} \text{ \&c.} = \text{fl. } \frac{y^6j}{1-y^{12}} - \text{fl. } \frac{y^{10}j}{1-y^{12}}.$$

Hence, and from Ex. 2. art. 32. the fluents of  $\frac{j}{1-y^{12}}$ ,  $\frac{y^4j}{1-y^{12}}$ ,  $\frac{y^6j}{1-y^{12}}$ , and  $\frac{y^{10}j}{1-y^{12}}$  may be readily obtained.

36. Multiplying the first equation in art. 28. by  $x + \sqrt{x^2 - 1} \Big|^{m-\frac{n}{2}}$  and the 2d. equat. by  $x - \sqrt{x^2 - 1} \Big|^{m-\frac{n}{2}}$ , bringing the products together according to our method, writing  $y^2$  instead of  $x$ , and afterwards multiplying by  $y^{p-1}y$  and taking the fluents, we have

$$\text{fl. } y^{p-1}y \times \frac{\binom{m}{s} \binom{n-m}{s} y^s}{1 + 2c y^q + y^{2q}} = \frac{\binom{m}{s} y^s}{p} - \frac{\binom{m+n}{s} y^{p+q}}{p+q} + \frac{\binom{m+2n}{s} y^{p+2q}}{p+2q} \text{ \&c.}$$

$$\text{and fl. } y^{p-1}y \times \frac{\binom{m}{c} \binom{n-m}{c} y^c}{1 + 2c y^q + y^{2q}} = \frac{\binom{m}{c} y^c}{p} - \frac{\binom{m+n}{c} y^{p+q}}{p+q} + \frac{\binom{m+2n}{c} y^{p+2q}}{p+2q} \text{ \&c.}$$

Whence, by considering  $y$  only as variable, we may obtain all the theorems which we, by considering  $x$  only as variable, have deduced in art. 28. 29. 30. and 33. and we may

may likewise obtain, from these two, some other theorems

not unworthy notice; the values of fl.  $y^{p-1}y \times \frac{\binom{m}{s} \binom{n-m}{s} y^s}{\binom{n}{s} (1 + 2c y^s + y^{2s})}$

and fl.  $y^{p-1}y \times \frac{\binom{m}{c} \binom{n-m}{c} y^c}{\binom{n}{c} (1 + 2c y^c + y^{2c})}$  being always assignable by means

of circular arcs and logarithms, as will be shewn in the Appendix.

Taking  $p$  and  $q$  equal to  $m$  and  $n$  respectively, we have

$$\begin{aligned} \text{fl. } y^{m-1}y \times \frac{\binom{m}{s} \binom{n-m}{s} y^s}{\binom{n}{s} (1 + 2c y^s + y^{2s})} &= \text{fl. } \frac{-y^{m-1}}{\sqrt{1-x^2}} \times \frac{\binom{m}{c} \binom{n-m}{c} y^c}{\binom{n}{c} (1 + 2c y^c + y^{2c})} \\ &= \frac{\binom{m}{s} y^s}{m} - \frac{\binom{m+n}{s} y^{m+n}}{m+n} + \frac{\binom{m+2n}{s} y^{m+2n}}{m+2n} \&c. \end{aligned}$$

$$\begin{aligned} \text{and fl. } y^{m-1}y \times \frac{\binom{m}{c} \binom{n-m}{c} y^c}{\binom{n}{c} (1 + 2c y^c + y^{2c})} &= \text{fl. } \frac{y^{m-1}}{\sqrt{1-x^2}} \times \frac{\binom{m}{s} \binom{n-m}{s} y^s}{\binom{n}{s} (1 + 2c y^s + y^{2s})} \\ &= \frac{\binom{m}{c} y^c}{m} - \frac{\binom{m+n}{c} y^{m+n}}{m+n} + \frac{\binom{m+2n}{c} y^{m+2n}}{m+2n} \&c. \end{aligned}$$

Hence the theorems in art. 33. and 30. may be obtained by taking  $m$  and  $n$  each equal to 1; and the theorems in art. 29. and 28. by taking  $m=1$  and  $n=2$ .

By expanding  $\left( x - \sqrt{x^2 - 1} \right)^{\frac{n}{2}} + y \cdot x + \sqrt{x^2 - 1} \left( x + \sqrt{x^2 - 1} \right)^{\frac{n}{2}}$  and  $\left( x + \sqrt{x^2 - 1} \right)^{\frac{n}{2}} + y \cdot x - \sqrt{x^2 - 1} \left( x - \sqrt{x^2 - 1} \right)^{\frac{n}{2}}$  by the binomial theorem, we may extend our analysis still farther.

37. It may be observed, that, as the value of the expression  $\frac{x^{p+1} - c^{p+1}}{p+1}$  becomes =  $\text{Log. } \frac{x}{c}$  when  $p$  is =  $-1$ , so, when  $p$  has that value ( $-1$ ), the value of the expression  $\frac{s^{(p+1)}}{p+1}$  becomes =  $x$ , and the value of the expression  $\frac{c^{(p+1)} - 1}{(p+1)^2}$  becomes =  $\frac{x^2}{2}$ : as appears by applying the rule derived from the doctrine of fluxions for finding the value of an algebraic fraction when its numerator and denominator both vanish together; the fluxion of  $s^{(p+1)}$  being =  $x^{(p+1)} c^p$ , and the fluxion of  $c^{(p+1)} = x^{(p+1)} s^p$ ,  $p$  only being considered as variable. Which being recollected, our theorems will be found useful in cases wherein they may without such recollection seem to fail.

I presume, I have now sufficiently explained this *new analytical Improvement*: by which the intelligent analyst may obtain the sums of many other series, and the fluents of many other fluxions, perhaps with greater facility than by any other method.

## P O S T S C R I P T.

The series considered in the following articles having relation to some of those in the preceding articles, it is thought not improper to take notice of them here, though their sums are obtained by a method different from our new method explained above.

38. It

38. It is well known that

$$\text{Log. } \frac{1}{1-x} \text{ is } = x + \frac{x^2}{2} + \frac{x^3}{3} \&c.$$

and, by substituting  $\frac{x}{x-1}$  for  $x$ , the same

$$\text{Log. } \frac{1}{1-x} \text{ is found } = -\frac{1}{1-x^{-1}} - \frac{1}{2} \frac{1}{1-x^{-1}} - \frac{1}{3} \frac{1}{1-x^{-1}} \&c.$$

From which equations, by substituting  $\frac{x-1}{x}$  (or  $1-x$ ) for  $x$ , we have

$$\begin{aligned} \text{Log. } x &= 1 - x^{-1} + \frac{1}{2} \frac{1}{1-x^{-1}} + \frac{1}{3} \frac{1}{1-x^{-1}} \&c. \\ \text{or } &= -\frac{1}{1-x} - \frac{1}{2} \frac{1}{1-x} - \frac{1}{3} \frac{1}{1-x} \&c. \end{aligned}$$

It follows therefore that

$$\text{fl. } \frac{x}{1-x} \text{ Log. } \frac{1}{1-x} \text{ is } = x + \frac{x^2}{2} + \frac{x^3}{3} \&c.,$$

$$\text{and fl. } \frac{x}{1-x} \text{ Log. } x = 1 - x + \frac{1-x}{2} + \frac{1-x}{3} \&c. - \frac{2a^2}{3},$$

the series  $1 + \frac{1}{2^2} + \frac{1}{3^2} \&c.$  being found  $= \frac{2a^2}{3}$ .

But the sum of the two fluents on one side is manifestly equal to  $\text{Log. } x \times \text{Log. } \frac{1}{1-x}$ ; which therefore is

$$= \begin{cases} x + \frac{x^2}{2} + \frac{x^3}{3} \&c. \\ + 1 - x + \frac{1-x}{2} + \frac{1-x}{3} \&c. - \frac{2a^2}{3}. \end{cases}$$

Hence, by taking  $x = \frac{1}{2}$ , it appears that, when  $x$  has that value

$\&c.$

$$x + \frac{x^2}{2^2} + \frac{x^3}{3^2} \&c. \text{ is } = \frac{x^2}{3} - \frac{1}{2} \text{sq. Log. } x.$$

39. Seeing that  $\frac{x}{x^2}$  is the fluxion of  $\frac{x^{-1}}{x}$ , we, from our two values of Log.  $x$  in the preceding article, get (by multiplying by  $\frac{x}{x} + \frac{x}{1-x}$ , or its equal  $\frac{-x}{x^2} \times \frac{x}{x-1}$ , and taking the fluents)

$$\left. \begin{aligned} 1 - x^{-1} + \frac{1-x^{-1}}{2^2} + \frac{1-x^{-1}}{3^2} \&c. \\ + 1 - x + \frac{1-x}{2^2} + \frac{1-x}{3^2} \&c. \end{aligned} \right\} = -\frac{1}{2} \text{sq. Log. } x:$$

and from hence the value of  $x + \frac{x^2}{2^2} + \frac{x^3}{3^2} \&c.$  when  $x$  is  $= \frac{1}{2}$ , may be found as above.

Moreover, taking  $1-x = \frac{x-1}{x}$ , we have

$$\left. \begin{aligned} \frac{1}{2} x x^2 + \frac{x^2}{2^2} + \frac{x^3}{3^2} \&c. \\ - x - \frac{x^2}{3^2} - \frac{x^3}{5^2} \&c. \end{aligned} \right\} = -\frac{1}{2} \text{sq. Log. } x;$$

$x^2$  being  $= 1-x = \frac{3-\sqrt{5}}{2}$ , and  $x = \frac{\sqrt{5}-1}{2}$

Therefore it appears that then

$$x^2 + \frac{x^2}{2^2} + \frac{x^2}{3^2} \&c. \text{ is } = \frac{1}{2} \times \frac{x^2}{x + \frac{x^2}{3^2} + \frac{x^2}{5^2} \&c.} - \frac{1}{2} \text{sq. Log. } x.$$

It also appears, by our value of Log.  $x \times \text{Log. } \frac{1}{1-x}$  in the preceding article, that, when  $x^2$  is  $= 1-x$ ,

Q

 $\frac{1}{2} \times$

$$\left. \begin{array}{l} \frac{1}{4} \times x^2 + \frac{x^4}{2^2} + \frac{x^6}{3^2} \&c. \\ x + \frac{x^3}{3} + \frac{x^5}{5} \&c. \end{array} \right\} \text{is} = \frac{2a^2}{3} - 2 \text{sq. Log. } x.$$

Consequently we have, in this case,

$$\begin{aligned} & x^2 + \frac{x^4}{2^2} + \frac{x^6}{3^2} \&c. \\ &= \frac{2}{15} a^2 - \frac{2}{5} \text{sq. Log. } x - \frac{1}{2} \times x + \frac{x^3}{3} + \frac{x^5}{5} \&c. \\ &= \frac{1}{2} \times x + \frac{x^3}{3^2} + \frac{x^5}{5^2} \&c. - \frac{2}{5} \text{sq. Log. } x, \end{aligned}$$

our former value of the same series.

Hence we find

$$x + \frac{x^3}{3} + \frac{x^5}{5} \&c. = \frac{a^2}{3} - \frac{1}{2} \text{sq. Log. } x,$$

$$x \text{ being} = \frac{\sqrt{5} - 1}{2}.$$

Now the value of this series being found, we, by means thereof, find

$$x^2 + \frac{x^4}{2^2} + \frac{x^6}{3^2} \&c. = \frac{a^2}{15} - \text{sq. Log. } x,$$

$$\text{or } \frac{1}{4} \times x^2 + \frac{x^4}{2^2} + \frac{x^6}{3^2} \&c. = \frac{a^2}{15} - \frac{1}{4} \text{sq. Log. } x.$$

$$\left. \begin{array}{l} \text{But } x + \frac{x^3}{3} + \frac{x^5}{5} \&c. \\ + \frac{1}{4} \times x^2 + \frac{x^4}{2^2} + \frac{x^6}{3^2} \&c. \end{array} \right\} \text{is} = x + \frac{x^3}{2^2} + \frac{x^5}{3^2} + \frac{x^7}{4^2} \&c.$$

$$\left. \begin{array}{l} \text{and } x + \frac{x^3}{3} + \frac{x^5}{5} \&c. \\ - \frac{1}{4} \times x^2 + \frac{x^4}{2^2} + \frac{x^6}{3^2} \&c. \end{array} \right\} = x - \frac{x^2}{2^2} + \frac{x^3}{3^2} - \frac{x^4}{4^2} \&c.$$

There-

## MEM. V.] OF THE SUMS OF SERIES.

Therefore it is evident, that

$$x + \frac{x^2}{2^2} + \frac{x^3}{3^2} + \frac{x^4}{4^2} \&c. \text{ is } = \frac{1}{2} a^2 - \text{sq. Log. } x,$$

$$\text{and } x - \frac{x^2}{2^2} + \frac{x^3}{3^2} - \frac{x^4}{4^2} \&c. = \frac{1}{15} a^2 - \frac{1}{2} \text{sq. Log. } x;$$

$$x \text{ being } = \frac{\sqrt{5}-1}{2}.$$

Mr. JOHN BERNOULLI, Mr. EULER, and some other authors have found the sums of the series  $y \pm \frac{y^2}{2^2} + \frac{y^3}{3^2} \&c.$  and  $y + \frac{y^2}{3^2} + \frac{y^3}{5^2} \&c.$  when  $y$  is  $= 1$ ; and the last mentioned gentleman, in his *Instit. Calc. Integ.* has also given the value of the series  $y + \frac{y^2}{2^2} + \frac{y^3}{3^2} \&c.$  when  $y$  is  $= \frac{1}{2}$ : which value I had before given in the *Philos. Transact.* for the Year 1760; together with the values of  $y \pm \frac{y^2}{2^2} + \frac{y^3}{3^2} \&c.$  and  $y + \frac{y^2}{3^2} + \frac{y^3}{5^2} \&c.$  in the first mentioned case.

In this Memoir, the value of the series  $y + \frac{y^2}{2^2} + \frac{y^3}{3^2} \&c.$  is assigned, not only in both those cases, but also when  $y$  is  $= \frac{\sqrt{5}-1}{2}$ , or  $= \frac{3-\sqrt{5}}{2}$ : its value being here found equal to  $\frac{2a^2}{3}$ ,  $\frac{a^2}{3} - \frac{1}{2} \text{sq. Log. } 2$ ,  $\frac{4a^2}{15} - \text{sq. Log. } \frac{\sqrt{5}-1}{2}$ , or  $\frac{1}{15} a^2 - \text{sq. Log. } \frac{\sqrt{5}-1}{2}$ , according as  $y$  is taken equal to  $1$ ,  $\frac{1}{2}$ ,  $\frac{3-\sqrt{5}}{2}$ , or  $\frac{\sqrt{5}-1}{2}$ .

Q 2

The



The value of  $y + \frac{y^3}{3} + \frac{y^5}{5} \&c.$  which before, I believe, had been assigned only when  $y$  is  $= 1$ , is found above, not only in that case, but likewise when  $y$  is  $= \frac{\sqrt{5}-1}{2}$ : and as the value of this last mentioned series (as also the value of  $y \pm \frac{y^3}{2} + \frac{y^5}{3} \&c.$ ) is useful, as well in assigning the sums of some other series of different forms, as in the calculation of certain fluents; I think it worth while to proceed to shew how the values of the said series  $y + \frac{y^3}{3} + \frac{y^5}{5} \&c.$  are obtained in two other cases.

40. It is obvious that

$$\text{fl. } \frac{x}{x} \cdot \text{fl. } \frac{x}{1-x^2} = \frac{1}{2} \text{Log. } x \times \text{Log. } \frac{1+x}{1-x} - \text{fl. } \frac{x}{1-x^2} \cdot \text{fl. } \frac{x}{x}$$

$$\text{is } = x + \frac{x^3}{3} + \frac{x^5}{5} \&c.$$

Let  $y$  be  $= \frac{1-x}{1+x}$ ; then will  $x$  be  $= \frac{1-y}{1+y}$ ,  $x = \frac{-2y}{1+y}$ ,

$\frac{x}{x} = \frac{-2y}{1-y}$ ,  $1-x^2 = \frac{4y}{1+y}$ , and  $\frac{x}{1-x^2} = \frac{-y}{2y}$ : and therefore it follows, that

$$\text{fl. } \frac{x}{1-x^2} \cdot \text{fl. } \frac{x}{x} \text{ will be } = \text{fl. } \frac{y}{y} \cdot \text{fl. } \frac{y}{1-y} = y + \frac{y^3}{3} + \frac{y^5}{5} \&c. - \frac{y^2}{2}$$

$$- g \text{Log. } \frac{1+x}{1-x}.$$

Consequently

$$x + \frac{x^3}{3} + \frac{x^5}{5} \&c. \text{ will be } = \frac{1}{2} \text{Log. } x \times \text{Log. } \frac{1+x}{1-x} + \frac{y^2}{2}$$

$$= y - \frac{y^3}{3} - \frac{y^5}{5} \&c. \quad \text{and}$$

$$\left. \begin{aligned} &\text{and } x + \frac{x^3}{3} + \frac{x^5}{5} \text{ \&c.} \\ &+ \frac{1-x}{1+x} + \frac{1-x^3}{3^2 \cdot 1+x} + \frac{1-x^5}{5^2 \cdot 1+x} \text{ \&c.} \end{aligned} \right\} = \frac{a^2}{2} + \frac{1}{2} \text{Log. } x \times \text{Log. } \frac{1+x}{1-x}.$$

Hence, by taking  $x = \sqrt{2} - 1$ , we find that, when  $y$  has that value,

$$y + \frac{y^3}{3} + \frac{y^5}{5} \text{ \&c. is } = \frac{a^2}{4} - \frac{1}{2} \text{sq. Log. } y:$$

and, by taking  $x = \frac{\sqrt{5}-1}{2}$ , (and referring to the preceding article for the value of  $x + \frac{x^3}{3} + \frac{x^5}{5} \text{ \&c.}$ ) we find that, when  $y$  is  $= \sqrt{5} - 2$ ,

$$y + \frac{y^3}{3} + \frac{y^5}{5} \text{ \&c. is } = \frac{a^2}{6} + \frac{1}{2} \text{sq. Log. } x - \frac{1}{2} \text{Log. } x \times \text{Log. } y.$$

By means of our conclusions respecting the values of the series  $y + \frac{y^3}{3} + \frac{y^5}{5} \text{ \&c.}$  and  $y \pm \frac{y^2}{2} + \frac{y^3}{3} \text{ \&c.}$  some useful theorems for the calculation of fluents may be obtained: for instance, it being known that the *whole* fluent of

$\frac{y}{\sqrt{r^2-y^2}} \cdot \text{fl. } \frac{y}{\sqrt{1-by^2}}$ , generated whilst  $y$  from 0 becomes

equal to  $r$ , is equal to  $r + \frac{br^3}{3} + \frac{b^2r^5}{5} \text{ \&c.}$  we, by referring to what is proved in this and the preceding article, easily infer that, the *whole* fluent of  $\frac{xy}{\sqrt{r^2-y^2}}$ , generated in the

time before mentioned, is equal to  $\frac{a^2}{3} - \frac{1}{2} \text{sq. Log. } r$ ,

$\frac{a^2}{4} - \frac{1}{2} \text{sq. Log. } r$ , or  $\frac{a^2}{6} + \frac{1}{2} \text{sq. Log. } \frac{1+r}{2} - \frac{1}{2} \text{Log. } r \times \text{Log.}$

$$\frac{1+r}{2},$$

$\frac{1+r}{2}$ , according as  $r$  is equal to  $\frac{\sqrt{5}-1}{2}$ ,  $\sqrt{2}-1$ , or  $\sqrt{5}-2$ , respectively;  $x$  denoting the circular arc whose radius is 1. and sine  $y$ , and  $a$  the quadrantal arc of the same circle.

Other theorems relating to fl.  $\frac{x}{x} \text{Log. } 1 \pm x$  and fl.  $\frac{x}{x} \text{Log. } \frac{1+x}{1-x}$  are more obvious.

41. By proceeding a step farther according to the method pursued in art. 38. and 39.

$$1 + \frac{1}{3^2} + \frac{1}{5^2} \&c. (= \frac{1}{7} \times 1 + \frac{1}{2^2} + \frac{1}{3^2} + \frac{1}{4^2} \&c.)$$

is found =  $\frac{1}{2}$  Cube Log.  $x - \frac{a^2}{3} \text{Log. } x + x + \frac{x^2}{2^2} + \frac{x^3}{3^2} \&c.$

and  $1 + \frac{1}{2^2} + \frac{1}{3^2} \&c. = \frac{1}{2}$  Cube Log.  $y - \frac{2a^2}{3} \text{Log. } y + \frac{1}{2} \times y^2 + \frac{y^3}{2^2} + \frac{y^3}{3^2} \&c.$

$x$  being =  $\frac{1}{2}$ ,  $y = \frac{\sqrt{5}-1}{2}$ ,  $y^2 = \frac{3-\sqrt{5}}{2}$ , and  $a$  as in the preceding article.

NOTE. All the *Logarithms* mentioned in this Memoir are of the *hyperbolic* kind.

---

---

## M E M O I R VI.

*A remarkable new Property of the Cycloid discovered, which suggests a new Method of regulating the Motion of a Clock.*

1. **L**ET  $A'B'P'E'Q'$ ,  $A''B''P''E''Q''$  be two similar, Fig. 38.  
curved, small tubes, situated exactly alike in a vertical plane; let a small ball be supposed to be put into each tube; and, both the balls  $P'$ ,  $P''$  being equal, let them be conceived to be connected by a perfectly flexible line, without weight, passing from  $P'$  up the tube wherein it is put to the top  $A'$ , and from thence to the top  $A''$  of the other tube, then down that other tube to  $P''$ : let that flexible line ( $P'A'A''P''$ ) be equal to  $E'B'A'A''B''E''$ ; and  $A'A''$ ,  $B'B''$ ,  $E'E''$ ,  $Q'Q''$  being horizontal lines, let  $B'E'$ ,  $E'Q'$ ,  $B''E''$ ,  $E''Q''$  be all equal, that, the balls being moved,  $P''$  may be at  $Q''$ ,  $E''$ , or  $B''$ , when  $P'$  shall be at  $B'$ ,  $E'$ , or  $Q'$  respectively. Then the ball  $P'$  being raised to  $B'$ , and left to descend from thence in the tube  $A'B'E'Q'$ ; and the ball  $P''$ , during the descent of  $P'$ , being drawn up the other tube from  $Q''$ , by means of the said connecting line; it is proposed to find the nature of the curve into which the tubes must be bent, that the time of descent of the ball  $P'$  (so connected with the ball  $P''$ ) from  $B'$  to  $Q'$  may always be the same, let the height  $B'E'$  be what it will.

Put

Put  $a$  for the length of the part  $B'E'$  of the tube into which the ball  $P'$  is supposed to be put;  $b$  for the vertical height of  $B'$  above  $E'$ ;  $x$  for the space passed over by  $P'$  in the tube in its descent;  $x$  for the vertical descent of  $P'$ ;  $y$  for the vertical ascent of  $P''$ ;  $v$  for the velocity of each of the balls;  $t$  for the time elapsed during the descent of  $P'$ ; and  $g$  for  $32\frac{1}{2}$  feet, the accelerative force of gravity: then will  $\frac{g \cdot x \cdot P'}{x}$  be the motive force by which the velocity  $v$  will be accelerated,  $\frac{g \cdot y \cdot P''}{x}$  the motive force by which  $v$  will be retarded, and  $\frac{1}{2}g \times \frac{x - y}{x}$  the actual accelerating force of each ball. Now, that  $P'$  may always arrive at  $E'$  in the same time, let the distance  $B'E'$  be what it will, its accelerating force must be always as the space to be passed over during such descent\*; that is,  $\frac{1}{2}g \times \frac{x - y}{x}$  must be  $= c \times \overline{a - x}$ ,  $c$  being some invariable quantity not yet known.

\* Let  $s$  be any space to be passed over,  $x$  a part of that space; and suppose that, in the time  $t$ , the moving body has passed over that part  $x$ , and has acquired a velocity  $v$  by the continued action of an accelerating force  $= c \times \overline{s - x}$ . Then will  $c \times \overline{s - x} \times \frac{x}{v}$  be  $= \dot{v}$ , and consequently  $c \cdot s \cdot x - \frac{1}{2}c \cdot x^2 = \frac{1}{2}v^2$ ,  $v$  being  $= 0$  when  $x$  is  $= 0$ . Moreover  $t$  ( $= \frac{x}{v}$ ) will be  $= \frac{x}{c \cdot \frac{1}{2} \sqrt{2s \cdot x - x^2}}$ : and hence  $t$  is found  $= \frac{1}{c \cdot \frac{1}{2}} \times$  Circ. Arc, radius  $r$ , versed sine  $\frac{x}{r}$ . Consequently, taking  $x$  equal to  $s$ , we find the *whole* time of passing over the space  $s = \frac{1}{c \cdot \frac{1}{2}} \times$  the quadrantal arc of the circle whose radius is  $r$ ; which,  $c$  being given, is always the same, let  $s$  be what it will.

Whence

Whence we have  $\frac{g}{c} \times \overline{x-y} = 2az - 2xz$ ; and, by taking the fluents, we find  $\frac{g}{c} \times \overline{x-y} = 2az - z^2$ .

2. Let ABPERQN be a semi-cycloid inverted, the diameter MHpIrKN of whose generating circle is  $d$ : let AB be  $= c$ , BE = EQ =  $a$ , HI =  $b$ , Hp =  $x$ , Kr =  $y$ , and BP = QR =  $z$ ; BH, Pp, EI, Rr, and QK being each parallel to the horizontal line AM. We shall then,

Fig. 39.

by the nature of the curve, have AN =  $2d$ , HN =  $\frac{2d-c}{4d}$ ,  
 $Np = \frac{2d-c-x}{4d}$ , NK =  $\frac{2d-c-2a}{4d}$ , Nr =  $\frac{2d-c-2a+z}{4d}$ ,  
 $\frac{2d-c-x-2z}{4d} = \text{HN} - Np = x$ , and  $\frac{2d-c-2a+z-2d-c-2a}{4d}$

= Nr - NK =  $y$ . Hence, it appearing by subtraction that  $x - y$  is =  $\frac{4az - 2z^2}{4d}$ , we have  $\frac{g}{c} \times \overline{x-y} = \frac{g}{c} \times \frac{4az - 2z^2}{4d}$ ;

which, if  $c$  be =  $\frac{g}{2d}$ , will be =  $2az - z^2$ , and the equation the same as that which we have deduced in the preceding article. It appears therefore, that our cycloid is the curve required: and, the accelerating force of the ball P' being =  $\frac{g}{2d} \times \overline{a-z}$ , the time of its descent from B' to Q'

(= twice the time of descent from B' to E') =  $\frac{2d}{g} \frac{1}{2}$  × semi-circle, rad. 1, which,  $d$  being given, will be the same, let B'E' be what it will; and will be equal to twice the time of free descent, from B to N, in the same cycloid; or the limit of the time of vibration (in a circular arc) of a pendulum whose length is  $2d$ .

R

It

It is obvious that the consequence will be the same, if  $P'$ ,  $P''$  be similar, slender chains perfectly flexible.

When  $P'$  shall have descended from  $B'$  to  $Q'$ ;  $P''$ , having been drawn up from  $Q''$  to  $B''$ , will begin to descend from the last mentioned point and draw  $P'$  upwards, so that a vibratory motion will ensue, which will be such, that, abstracting from friction, the time of vibration will be the same, from what point soever  $P'$  may begin to move, and whatever may be the length of the line connecting the balls or chains. By means of which line, a rod applied to a clock may be made to vibrate in any plane whatever: and, only small parts of the cycloidal tubes being requisite, the mechanism may, in a little room, be so adapted, by taking the diameter  $d$  of a proper length, (agreeable to what is proved above,) that any given number of vibrations shall be performed in a given time.

The evolute of the cycloid being a similar cycloid, the balls ( $P'$ ,  $P''$ ) may be easily made to describe any cycloidal arcs by evolution: and, by substituting evolutes instead of tubes, the friction of the movement may be diminished; but it will then take up more room.

---



---

## M E M O I R VII.

*Of the Motion of a Body, keeping always in the same given Plane, whilst acted on by any Force, or Forces, urging it continually to change its Direction in that Plane.*

1. **L**ET a body (B) be supposed to describe the tra- Fig. 40.  
 jectory ABb about the center C: let (CB) its  
 distance from C be denoted by  $y$ : let  $B \times f$  denote a mo-  
 tive force urging it continually towards the said center:  
 let  $B \times g$  denote another such force always acting on it at  
 right angles to (CB) the radius vector, or ray drawn from  
 C to the body: let  $v$  denote the velocity of the body  
 from the center C, in the variable direction CB; (which  
 is sometimes called its *paracentric velocity*);  $u$  its angular  
 velocity about C, measured at the distance  $r$  therefrom;  
 and  $w$  its velocity at B in the direction Bb, at right angles  
 to CB; (which is sometimes called its *circulatory velocity*.)  
 Then, the right line BP being a tangent to the trajectory  
 at B; if the right line CP, drawn in any manner from C  
 to that tangent, be denoted by  $p'$ ; and the sine of the  
 angle CPB (to the radius 1) by  $s$ ; the sine of the angle  
 CBP will be  $= \frac{p's}{y}$ , and the absolute velocity of the body,  
 in the direction of the tangent PBb, (or particle Bb of  
 the curve,) will be  $= \frac{wy}{p's}$ : which would continue in-  
R 2
variable



variable if no force acted on the body. In which case it would continue to move in the right line  $PBb$ , and  $p'$  and  $s$  would remain invariable. Therefore  $\frac{\dot{w}y + w\dot{y}}{p's}$ , the fluxion of its velocity, would, in that case, be = 0, and  $\dot{w} = -\frac{w\dot{y}}{y} = -\frac{wy}{y}$ ,  $w$  being =  $\frac{wy}{y}$ . But the motive forces  $B \times f$  and  $B \times g$  (answering to the accelerating or retarding forces  $f$  and  $g$ )\* acting on the body,  $\dot{w}$  will be =  $\frac{wy + wy}{y}$ . Consequently, the directions  $BC$ ,  $Bd$  being at right angles to each other,  $\frac{wy + 2wy}{y}$  (the excess of  $\frac{wy + wy}{y}$ , the fluxion of  $w$  when the forces act, above  $-\frac{wy}{y}$ , which would be the fluxion of  $w$  if the forces ceased to act) will be the fluxion of the velocity generated or destroyed by the action of the force  $B \times g$  only: and, the motive force into  $\left(\frac{y}{v}\right)$  the fluxion of the time being equal to the fluxion

\* The quantity here called an accelerating or retarding force is the measure of the force causing acceleration or retardation; and is denoted or expressed by the *velocity* that would be generated or destroyed in the same or an equal body, in a given time, by the action of the same or an equal force uniformly continued in the direction in which the body, during such uniform action, might move.

The measure of such a force may also, in comparing the effects of forces, be denoted or expressed by the *space* which a body would be made to pass over in a given time by the action of the same or an equal force uniformly continued; such space being always the *half* of the space which a body would pass over with the velocity acquired in that time by the action of such force, and the halves of quantities being to each other as the whole quantities are to each other respectively: but due care must be taken, that these two ways of expressing the measure of such force be not confounded and both used in the same process.

of

of the quantity of motion generated or destroyed by that force,  $B \times g \times \frac{j}{v}$  will be  $= B \times \frac{\dot{u}y + 2uj}{r}$ , and  $g = \frac{v}{r} \times \frac{\dot{u}y + 2uj}{j}$ .

Moreover  $\sqrt{v^2 + w^2}$ , the velocity of the body in the direction  $PBb$ , would be invariable if the forces ceased acting: therefore  $\frac{v\dot{v} + w\dot{w}}{\sqrt{v^2 + w^2}}$  (the fluxion of that velocity)

would then be  $= 0$ , and  $\dot{v} = -\frac{w\dot{w}}{v} = \frac{u^2yj}{r^2v}$ ,  $w$  (in that case)

being  $= \frac{uy}{r}$  and  $\dot{w} = -\frac{uj}{r}$ , by what is said above. But,

the motive forces  $B \times f$  and  $B \times g$  continuing to act on the body, the fluxion of  $v$  will in general be expressed by  $\dot{v}$ . Consequently (the directions  $BC$ ,  $Bd$  being at right

angles to each other)  $\frac{u^2yj}{r^2v} - \dot{v}$  will be the fluxion of the

velocity destroyed or generated by the motive force  $B \times f$ ; and therefore ( $f$ ) the retarding or accelerating force produced by the action of the said motive force (being equal to the last mentioned fluxion divided by the fluxion of the

time) will be equal to  $\frac{u^2yj - r^2v\dot{v}}{r^2j}$ .

The general theorems

$$f = \frac{u^2yj - r^2v\dot{v}}{r^2j}, \text{ and } g = \frac{v}{r} \times \frac{\dot{u}y + 2uj}{j},$$

which we have investigated (with others that will be suggested by the particular nature or circumstances of the proposition to be considered) will be sufficient for the purpose of determining every thing that may be required respecting

respecting the trajectory and the motion of the body therein; as every force which can act on the body, in any direction (in our given plane) different from  $BC$ ,  $Bd$ , (the two directions in which we have supposed the forces  $B \times f$ ,  $B \times g$  to act,) may be resolved into two forces acting in those two particular directions.

2. It is obvious that our two forces  $f$  and  $g$  may be resolved into two others  $h$  and  $k$ ; the former in the direction of the tangent to the trajectory, retarding or accelerating the absolute velocity of the body; and the latter in a direction at right angles to the said tangent, neither retarding nor accelerating the absolute velocity of the body, but only changing its direction. Which forces found by so resolving the forces above investigated are

$$h = \frac{-v}{\sqrt{r^2 v^2 + u^2 y^2}} \times \frac{r^2 v \dot{v} + u \dot{u} y^2 + u^2 y \dot{y}}{r y},$$

$$k = \frac{1}{\sqrt{r^2 v^2 + u^2 y^2}} \times \frac{r^2 v^2 \dot{u} y - r^2 v \dot{v} u y + 2 r^2 v^2 u \dot{y} + u^2 y^2 \dot{y}}{r^2 y}.$$

3. It is likewise obvious that one single force ( $F$ ) compounded of the forces  $f$  and  $g$ , or  $h$  and  $k$ , is sufficient to cause a projectile to describe any trajectory whatever, with a velocity varying in any possible manner; the direction in which  $F$  must act being properly varied at every instant: which single force is  $= \sqrt{f^2 + g^2}$ ; the direction in which it acts at every instant making angles with the radius vector and tangent to the trajectory, such that their sines shall be to radius as  $g$  to  $F$ , and  $\frac{fuy + grv}{\sqrt{r^2 v^2 + u^2 y^2}}$

to F, and their cosines to radius as  $f$  to F, and  $\frac{frv - gvy}{\sqrt{r^2v^2 + u^2y^2}}$  to F respectively.

4. If  $\dot{x}$  be the fluxion of the circular arc described by a point in the radius vector, at the distance  $r$  from the center C,  $v$  will be to  $\frac{uy}{r}$  as  $\dot{y}$  to  $\frac{y\dot{x}}{r}$ , and  $uy = v\dot{x}$ .

5. If the trajectory, instead of being referred to the center C, be referred to a base; and,  $y$  being the ordinate corresponding (at right angles) to the abscissa  $x$ ,  $v$  be the velocity of the body from the base, (in the direction of the said ordinate,) and  $u$  its velocity in a direction parallel to the base; the force ( $f'$ ) urging it towards the base will be  $= -\frac{v\dot{v}}{y}$ , and the force ( $g'$ ) urging it in a direction parallel to the base will be  $= \frac{v\dot{u}}{y}$ . Moreover  $uy$  will be  $= v\dot{x}$ .

6. Which two last mentioned forces may be resolved into two others,  $h'$  and  $k'$ : whereof the former, in the direction of the tangent to the trajectory, (retarding or accelerating the velocity of the body,) will be  $= -\frac{v^2\dot{v} + v\dot{u}u}{\sqrt{v^2 + u^2} \times y}$ ; and the latter, in a direction at right angles to the said tangent, (changing the direction of the body,)  $= \frac{v^2\dot{u} - v\dot{v}u}{\sqrt{v^2 + u^2} \times y}$ .

The use of the theorems, obtained above with such facility, is far more extensive than the common doctrine  
of

of centripetal forces; as will in some measure appear by the following articles.

7. If only the centripetal force  $B \times f$  act on the body;  $g$  being = 0, we have  $u\dot{y} + 2u\dot{y} = 0$ ; and, by taking the fluents, after multiplying by  $y$ , we get  $uy^2 = a^2b$ ,  $b$  being the value of  $u$  when  $y$  is =  $a$ . Therefore,  $u$  being =  $\frac{a^2b}{y^2}$ ,  $w$  ( $= \frac{u\dot{y}}{r}$ ) will be =  $\frac{a^2b}{ry}$ ; and (U) the absolute velocity of the body in its trajectory =  $\frac{\sqrt{a^4b^2 + r^2v^2y^2}}{ry}$ , the value of  $\sqrt{v^2 + w^2}$ : which (being to  $w$  as  $y$  to  $p$ ) is also =  $\frac{a^2b}{rp}$ ,  $p$  denoting the perpendicular from C, (the center of force) to the tangent to the trajectory at the point where the body shall then be. Moreover the force

$$\begin{aligned} f \left( = \frac{u^2y}{r^2} - \frac{v\dot{v}}{y} \right) & \text{ will be } = \frac{a^4b^2}{r^2y^3} - \frac{v\dot{v}}{y} \\ & = \frac{a^4b^2}{r^2y^3} - \frac{a^4b^2}{2y} \times \text{the fluxion of } \frac{y^2}{y^2x^2} = \frac{a^4b^2p}{r^2p^3y}, \\ v \text{ being } & = \frac{u\dot{y}}{r} = \frac{a^2b\dot{y}}{y^2x} = \frac{a^2b\sqrt{y^2 - p^2}}{rpy}. \end{aligned}$$

Hence, by substitution, the requisite force may be found, when the nature or equation of the curve is given: or, if the force be given in terms of  $p$ ,  $v$ ,  $u$ ,  $y$ , or  $x$ ; the nature of the curve described by the projectile, and its motion therein, may be determined.

*Example 1.* If,  $r$  being taken =  $a$ , the centripetal force  $f$  be =  $\frac{a^2b^2}{y^3}$  ( $= \frac{u^2y}{r^2}$ ) = the centrifugal force arising from the  
circulatory

circulatory velocity of the projectile;  $\dot{v}$  will be = 0, and  $\dot{z}$  ( $= \frac{uj}{v} = \frac{uj}{c}$ ) =  $\frac{a^2 bj}{cy^2}$ : whence  $z = \frac{ab}{c} \times \frac{y-a}{y}$ .

It is remarkable, that, when  $y$  is infinite,  $\rho$  and  $z$  will each be =  $\frac{ab}{c}$ .

*Example 2.* Let  $f$  be supposed =  $\frac{Av\dot{v}}{Bj}$ ; as it will be in the case of a body B revolving on a horizontal plane, about a given point (C), and drawing another body A, on the same plane, directly towards that point, by means of a string connecting the bodies\*.

In this instance,  $f$ , which we above found equal to  $\frac{a^2 b^2}{r^2 j^2} - \frac{v\dot{v}}{j}$ , being =  $\frac{Av\dot{v}}{Bj}$ , we have, from that equation,  $\frac{a^2 b^2 j}{r^2 j^2} = \frac{A+B}{B} v\dot{v}$ : whence, by taking the fluents, we get  $\frac{a^2 b^2}{r^2} \times \frac{y^2 - a^2}{j^2} = \frac{A+B}{B} v^2$ ,  $v$  being supposed = 0 when  $y$  is =  $a$  and  $z = b$ . Therefore  $v$  will be =  $\frac{bP^{\frac{1}{2}}}{j} \times \sqrt{y^2 - a^2} = \frac{a^2 bj}{y^2 z}$ , the value of  $v$  found above; and  $\dot{z}$ , in consequence, =  $\frac{a^2}{P^{\frac{1}{2}}} \times \frac{y^{-1} j}{\sqrt{y^2 - a^2}}$ ;  $r$  being taken equal to  $a$ , and P being put for  $\frac{B}{A+B}$ . By taking the fluents again,  $z$  is found

\* In this case  $f$  is =  $\frac{T}{B}$ , T being the tension of the string: and it is obvious that  $(Tj =) \frac{Tj}{v}$  is =  $Av\dot{v}$ : therefore  $f$  is =  $\frac{Av\dot{v}}{Bj}$ .

$= \frac{1}{p^{\frac{1}{2}}} \times \text{circ. arc, rad. } a, \text{ sec. } y:$  and  $t$  the time elapsed (computed from the time of  $y$  being  $= a$ ) is, from the equation  $\dot{r} = \frac{\dot{x}}{u} = \frac{\dot{y}}{v} = \frac{y\dot{y}}{bP^{\frac{1}{2}}\sqrt{y^2-a^2}}$ , found  $= \frac{\sqrt{y^2-a^2}^{\frac{1}{2}}}{bP^{\frac{1}{2}}}$ .

It is observable that,  $z$  being  $(= \frac{da}{P^{\frac{1}{2}}}) = \frac{1}{P^{\frac{1}{2}}} \times$  the quadrantal arc of the circle whose radius is  $a$ , when  $y$  is infinite; the body B, after having made a number of revolutions  $(= \frac{1}{4P^{\frac{1}{2}}} = \frac{1}{4} \times \frac{A+B}{B})$  about the center C, will fly off to an infinite distance from that point; approaching continually to a rectilinear asymptote, whose distance from C is  $= \frac{A+B}{B} \times a$ , parallel to the ultimate direction of the radius vector: which direction will be known, the ultimate value of  $z$  being (manifestly  $= \frac{A+B}{B} \times da$ ) as above mentioned.

*Example 3.* Let  $f = \frac{a^2 b^2}{y^2} - \frac{a^2 b^2}{2y} \times$  the fluxion of  $\frac{y^2}{y^2 z^2}$  be supposed  $= \frac{M}{y^m} + \frac{N}{y^n}$ ,  $r$  being taken  $= a$ . Then, by multiplying by  $2y$  and taking the fluents, we have  $b^2 + c^2 - \frac{a^2 b^2}{y^2} - \frac{a^2 b^2 y^2}{y^4 z^2} = \frac{2Ma^{1-m}}{m-1} + \frac{2Na^{1-n}}{n-1} - \frac{2My^{1-m}}{m-1} - \frac{2Ny^{1-n}}{n-1}$ ,  $c$  being the value of  $v$  when  $y$  is  $= a$ .

Whence

$$\text{Whence } \dot{z} \text{ is found} = \frac{a^2 b y^{-1} j}{\sqrt{K y^2 + \frac{2M y^{3-m}}{m-1} + \frac{2N y^{3-n}}{n-1} - a^2 b^2}}$$

$$K \text{ being} = b^2 + c^2 - \frac{2M a^{1-m}}{m-1} - \frac{2N a^{1-n}}{n-1}.$$

And it follows from what is done above, that

$$p \text{ will be} = \frac{ab}{\sqrt{K + \frac{2M y^{1-m}}{m-1} + \frac{2N y^{1-n}}{n-1}}}$$

$$u = \frac{a^2 b}{y^2}, \quad v = \sqrt{K + \frac{2M y^{1-m}}{m-1} + \frac{2N y^{1-n}}{n-1} - a^2 b^2 y^{-2}};$$

$\sqrt{v^2 + w^2}$  (the absolute velocity of the body in its trajectory)

$$= \sqrt{v^2 + \frac{u^2 y^2}{a^2}} = \sqrt{v^2 + \frac{a^2 b^2}{y^2}} = \sqrt{K + \frac{2M y^{1-m}}{m-1} + \frac{2N y^{1-n}}{n-1}};$$

$$\text{and } t \text{ (the fluxion of the time)} = \frac{y \dot{y}}{\sqrt{K y^2 + \frac{2M y^{3-m}}{m-1} + \frac{2N y^{3-n}}{n-1} - a^2 b^2}}$$

Now, though we cannot in general determine the trajectory, and the motion of the projectile therein, from these equations, without sometimes having recourse to infinite series; yet there are certain cases wherein the curve, and the motion of the body describing it, may be determined by means of logarithms, or circular or elliptic arcs: concerning some of which cases, I purpose to subjoin a few remarks.

*Remark 1.* If  $m$  be  $= 2$  and  $n = 3$ ,  $\dot{z}$  will be equal to  $\frac{a^2 b y^{-1} j}{\sqrt{K y^2 + 2M y + N - a^2 b^2}}$ ,  $K$  being  $= b^2 + c^2 - 2M a^{-1} - N a^{-2}$ :

S 2

and



and it appears by our Appendix, that, in all possible cases, the fluent  $z$  will be assigned by logarithms or circular arcs.

$$\text{If } K \text{ be } = \frac{M^2}{N - a^2 b^2}, \dot{z} \text{ will be } = \frac{a^2 b \sqrt{N - a^2 b^2}}{M} \times \frac{y^{-1} j}{y + \frac{N - a^2 b^2}{M}}$$

$$\text{and } z = \frac{a^2 b}{\sqrt{N - a^2 b^2}} \times \text{Log.} \frac{aM}{aM + N - a^2 b^2} \times \frac{y - \frac{N - a^2 b^2}{M}}{y};$$

$$\text{in which case } c^2 \text{ is } = \frac{M^2 a^{-2}}{N - a^2 b^2} \times a + \frac{N - a^2 b^2}{M}.$$

Consequently, if, whilst the body is continually urged *towards* the center C by a force  $\frac{N}{y^2}$ , it be continually urged *from* that center by a force  $\frac{M}{y^2}$ , N being greater than  $a^2 b^2$ ; the projectile, supposing  $y$  to increase after being =  $a$ , (that is, supposing  $c$  to be positive,) will revolve in a spiral about C, and for ever recede therefrom; yet never will arrive at the periphery of the circle whose center is C, and radius =  $\frac{N - a^2 b^2}{M}$ , a given quantity greater than  $a$ : or, supposing  $y$  to decrease after being =  $a$ , (that is, supposing  $c$  negative,) the body will revolve about C, and continually approach nearer and nearer to that point; yet never will get within the circle whose center is C and radius =  $\frac{N - a^2 b^2}{M}$ , a given quantity less than  $a$ : M being accordingly less than  $\frac{N - a^2 b^2}{a}$  when the body recedes from the center, or greater than  $\frac{N - a^2 b^2}{a}$  when it approaches towards the center.

*Remark*

*Remark 2.* If  $m$  be  $= 2$  and  $n = 4$ ,  $\dot{z}$  will be equal to  $\frac{a^2 b y^{-\frac{1}{2}}}{\sqrt{K y^3 + 2 M y^2 - a^2 b^2 y + \frac{1}{3} N}}$ ,  $K$  being  $= b^2 + c^2 - 2 M a^{-1} - \frac{1}{3} N a^{-3}$ ; and it appears by our Tables, that the fluent  $z$  will always be assigned by elliptic or circular arcs, or by logarithms, or by algebraic quantities.

If  $M$  be  $= 0$ ,  $\dot{z}$  will be  $= \frac{a^2 b y^{-\frac{1}{2}}}{\sqrt{K y^3 - a^2 b^2 y + \frac{1}{3} N}}$ ; which will be  $= \frac{3^{\frac{1}{2}} a A y^{-\frac{1}{2}}}{\sqrt{y^3 - 3 A^2 y + 2 A^3}}$  ( $= \frac{3^{\frac{1}{2}} a A y^{-\frac{1}{2}}}{y - A \sqrt{y + 2 A}}$ ) when  $K$  is  $= \frac{N}{3 A^3}$ ,  $A$  being  $= \frac{N}{a^2 b^2}$ . Hence, by taking the fluents, we find

$$z = a \times \text{Log.} \frac{y - A}{a - A} \times \frac{a + 2A + \sqrt{6aA + 3a^2}}{y + 2A + \sqrt{6Ay + 3y^2}}$$

Consequently, if  $\sqrt{\frac{b^2}{3aA^3} \times (a - A)^2 \times a + 2A}$  (the value of  $c$ ) be positive, the projectile will describe a spiral about  $C$ , and for ever recede therefrom; yet never will arrive at the periphery of the circle whose center is  $C$  and radius  $\frac{N}{a^2 b^2}$ : or, if the value of  $c$  be negative, the body will revolve about  $C$ , and continually approach nearer and nearer to that point; yet never will get within the circle whose center is  $C$  and radius  $\frac{N}{a^2 b^2}$ :  $N$  being greater or less than  $a^3 b^2$ , according as the body recedes from, or approaches towards the center.

$$\text{If } K \text{ be } = -\frac{4M^2}{3a^2 b^2} \text{ and } N = \frac{1}{4} \frac{a^2 b^2}{M},$$

$\dot{z}$  will

$z$  will be  $= \frac{3^{\frac{1}{2}} A a y^{-\frac{1}{2}} j}{A - y)^{\frac{1}{2}}}$ ; and, by Theorem III. Table L.

$$z = 2 \cdot 3^{\frac{1}{2}} a \times \frac{j^{\frac{1}{2}}}{A - y)^{\frac{1}{2}}} - \frac{a^{\frac{1}{2}}}{A - a)^{\frac{1}{2}}}, \text{ A being } = \frac{1}{2} \frac{a^2 b^2}{M}.$$

Consequently, if,  $M$  being a positive quantity less than  $\frac{1}{2} a b^2$ ,  $\left( \frac{b \times \overline{A - a}^{\frac{1}{2}}}{3^{\frac{1}{2}} a^{\frac{1}{2}} A} \right)$  the value of  $c$  be taken positive, the projectile will describe a spiral having a circular asymptote whose center is  $C$  and radius  $\frac{1}{2} \frac{a^2 b^2}{M}$ .

*Remark 3.* If,  $n$  being  $= 4$ ,  $m$  be  $= 3$ ; or  $n$  being  $= 5$ ,  $m$  be equal to 2, 3, or 4; the value of  $z$  will, by our Tables, be always assigned by elliptic or circular arcs, or by logarithms.

If,  $n$  being  $= 5$ ,  $M$  be  $= 0$ ,

$$z \text{ will be } = \frac{a^2 b j}{\sqrt{K y^4 - a^2 b^2 y^2 + \frac{1}{2} N}}$$

which, when  $K$  is  $= \frac{1}{2} \frac{N}{A^2}$ , will be  $= \frac{2^{\frac{1}{2}} A a j}{y^2 - A^2}$ ,  $A$  being  $= \frac{N^{\frac{1}{2}}}{a b}$ . Hence  $z$  is found  $= \frac{a}{2^{\frac{1}{2}}} \times \text{Log. } \frac{A + a}{A - a} \times \frac{A - y}{A + y}$ .

Consequently, if  $\sqrt{\frac{b^2}{2 a^2 A^2} \times \overline{a^2 - A^2}}$  (the value of  $c$ ) be taken either positive or negative, the projectile will describe a spiral having a circular asymptote whose center is  $C$  and radius  $\frac{N^{\frac{1}{2}}}{a b}$ ;  $N$  being greater or less than  $a^4 b^2$ , according

ording as the body recedes from, or approaches towards, the center.

8. The remarkable circumstance of the trajectory continually approaching to a circular asymptote, I believe, was first taken notice of by MR. MACLAURIN, in the case wherein ( $g$  being  $= 0$ )  $f$  is  $= \frac{N}{y^3}$ : afterwards Mr. SIMPSON observed, that the same thing will happen in an infinity of other cases,  $g$  being  $= 0$ ,  $n$  greater than 3, and  $f = \frac{N}{y^n}$ .

It is farther observable, that such circumstance always takes place when,  $g$  being  $= 0$ , the value of  $f$  is expressed in any manner whatever: provided the values of  $b$ ,  $c$ , and  $y$ , determined from the equations  $\overline{b^2 + c^2} \cdot y^2 - a^2 b^2 - 2y^2 \times$  the fluent of  $f\dot{y} = 0$  and  $\dot{y} = \dot{p} \left( = \frac{abf\dot{y}}{b^2 + c^2 - 2 \text{ fluent of } f\dot{y}} \right)^{\frac{1}{2}}$ ,

be real, the fluent mentioned in these expressions being generated whilst  $y$  from being  $= a$  becomes  $= A =$  the value of  $y$  in the equation ( $f y^3 = a^2 b^2$ ) resulting from those two equations. Which equations are obtained by considering  $C$  as the center of curvature of the trajectory when the value of  $v$  is  $= 0$ . From whence it follows, that ( $A$ ) the value of  $y$  in the equation  $f y^3 = a^2 b^2$  will be the radius of the asymptotic circle; the value of  $f$  when  $y$  is  $= a$  being so adapted, that  $A$  shall be greater or less than  $a$ , according as the body is considered as ascending or descending towards such asymptote, that is, according as  $c$  is taken positive or negative.

9. If only the force  $g$  act on the projectile;  $f$  being  $= 0$ , we have  $\frac{v^2}{r} = \frac{v\dot{v}}{y\dot{y}}$ : by means of which equation and that

that of the proposed curve, the requisite value of  $\frac{v}{r}$   $\times \frac{uy + 2uj}{j}$  ( $=g$ ) may be readily computed; or,  $g$  being given, the trajectory, and the motion of the body therein, may be determined.

*Example 1.* To find the force requisite (at every instant) to cause a projectile to describe a logarithmic, or equiangular spiral, by acting thereon always in a direction at right angles to the radius vector?

Let  $p$  be  $=my$ ,  $m^2 + n^2$  being  $=1$ ; then will  $\dot{z}$  be  $=\frac{mry}{ny}$   $=\frac{uj}{v}$ : whence  $u = \frac{mrv}{ny}$ . Therefore  $\frac{v\dot{v}}{yy}$  ( $=\frac{u^2}{r^2}$ ) will be  $=\frac{m^2v^2}{n^2y^2}$ , and  $\frac{\dot{v}}{v} = \frac{m^2}{n^2} \cdot \frac{j}{y}$ : hence  $v$  is found  $=c \cdot \left(\frac{y}{a}\right)^{\frac{n^2}{m^2}}$ . Consequently  $u$  will be  $=b \cdot \left(\frac{y}{a}\right)^{\frac{n^2-n^2}{n^2}}$ , and  $g = \frac{b^2}{mna} \cdot \left(\frac{y}{a}\right)^{\frac{2n^2-n^2}{n^2}}$ ;  $r$  being taken  $=a$ , and  $nb$  being then  $=mc$ .

If  $m$  be  $=\frac{1}{\sqrt{3}}$  and  $n = \sqrt{\frac{2}{3}}$ ,  $g$  will be equal to the invariable quantity  $\frac{3b^2}{2^{\frac{1}{2}}a}$ .

*Example 2.* Let the trajectory be a conic section, whereof  $C$  is a focus, and whose semi-axes are  $\frac{d}{1 \mp e}$  and  $\frac{d}{\sqrt{1 \mp e^2}}$ ; the curve being a circle when  $e$  is  $=0$ ; and an ellipsis, a parabola, or a hyperbola, according as  $e$  is less, equal to, or greater than  $1$ .

Then

Then will  $p$  be  $= \frac{dy^{\frac{1}{2}}}{\sqrt{2d + e^2 - 1.y}} = \frac{xy^2}{\sqrt{r^2v^2 + u^2y^2}}$ : and, ex-terminating  $v\dot{v}$ , we find

$$\frac{\dot{u}}{u} = \frac{dy}{e^2 - 1.y^2 + 2dy - d^2} - \frac{2\dot{y}}{y};$$

whence, by taking the fluents,  $u$  is found  $= \frac{a^2b}{y^2} \times$

$\frac{\left[ \frac{e+1.y-d \times e-1.a+d}{e-1.y+d \times e+1.a-d} \right]^{\frac{1}{2e}}}{\left[ \frac{e-1.a+d}{e+1.a-d} \right]^{\frac{1}{2e}}}$ . Consequently  $v$  will be equal to

$$\frac{a^2b}{dry} \times \frac{\left[ \frac{e-1.a+d}{e+1.a-d} \right]^{\frac{1}{2e}}}{\left[ \frac{e-1.y+d}{e-1.y+d} \right]^{\frac{1+e}{2e}}}$$

Which being known, we are thereby enabled to obtain, by substitution, the requisite value of the force

$$g = \frac{a^2b^2}{r^2y^2} \times \frac{\left[ \frac{e-1.a+d}{e+1.a-d} \right]^{\frac{1}{e}}}{\left[ \frac{e+1.y-d}{e-1.y+d} \right]^{\frac{2+e}{2e}}}$$

By which it appears, that the body cannot begin to move from the vertex of the curve unless  $e$  be  $= 2$ . In which case the curve is an equilateral hyperbola, whose semi-axes are each  $= d$ ; and  $g$  must then be  $= \frac{a^2b^2}{r^2y^2} \times$

$\frac{\left[ \frac{a+d}{3a-d} \right]^{\frac{1}{2}} \times \frac{1}{y+d}}$ . Therefore, taking  $b^2 = \overline{3a-d}^{\frac{1}{2}}$  and  $r = a = \frac{d}{3}$ , it follows, that the body resting at the vertex of the curve may be moved from thence by the force  $g$ , (at first  $= \frac{3^{\frac{1}{2}}}{2d^{\frac{1}{2}}}$ ) and be afterwards made to describe such

T

hyper-

hyperbola by the continued action of the force  $g =$

$$\frac{2d^{\frac{1}{2}}}{3^{\frac{1}{2}}y^2y + d}$$

*Example 3.* Let  $u$  be always  $= \frac{Mb}{N+y^2}$ ; as it will be in the case of a ball moving within a tube (or a ring upon a rod) revolving in a horizontal plane about one of its ends; the ball and tube, (or ring and rod) after the first percussive impulse, being disturbed by no force but their mutual pressure against each other\*.

$$\text{Then will } \frac{v\ddot{v}}{y\dot{y}} (= \frac{u^2}{r^2}) \text{ be } = \frac{M^2b^2}{r^2.N+y^2}, \text{ and } v\dot{v} = \frac{M^2b^2y\dot{y}}{r^2.N+y^2} =$$

$$\text{hence } v \text{ is found } = \sqrt{c^2 + \frac{M^2b^2}{r^2.N+a^2} - \frac{M^2b^2}{r^2.N+y^2}}. \text{ Confe-}$$

$$\text{quently } g \text{ will be } = \frac{2MNb}{r^2.N+y^2} \times \sqrt{r^2c^2 + \frac{M^2b^2}{N+a^2} - \frac{M^2b^2}{N+y^2}}$$

$$\text{and } \dot{z} (= \frac{u\dot{y}}{v}) = \frac{Mrb.N+y^2)^{-\frac{1}{2}}\dot{y}}{\sqrt{r^2c^2 + Mb^2 \times N+y^2 - M^2b^2}} \cdot \frac{M}{N+a^2} \text{ being } = 1.$$

$$\text{Which value of } \dot{z} \text{ becomes } = \frac{\frac{1}{2}Mrbx^{-\frac{1}{2}}(x-N)^{-\frac{1}{2}} \times \dot{x}}{\sqrt{r^2c^2 + Mb^2 \times x - M^2b^2}}, \text{ upon}$$

substituting  $x$  for  $N+y^2$ ; and therefore it appears by our Appendix, that the fluent  $x$  will always be assigned by elliptic arcs, or by logarithms.

\* It will appear by a subsequent Memoir, that, in this case, when the tube (or rod) is very slender,  $N (= M - a^2)$  is  $= \frac{r^2T}{3B}$ ;  $l$  being the length of the tube (or rod),  $T$  its weight, and  $B$  the weight of the ball (or ring).

If,

If, taking  $r = a$ ,  $N$  be  $= \frac{a^2 b^2}{c^2 - b^2}$ ;  $M$  will be  $= \frac{a^2 c^2}{c^2 - b^2}$ ,

$$\dot{z} = \frac{aN^{\frac{1}{2}}y^{-1}\dot{y}}{\sqrt{N+y^2}}, \quad z = \frac{1}{2}a \times \text{Log.} \frac{b+c}{b-c} \times \frac{N^{\frac{1}{2}} - \sqrt{N+y^2}}{N^{\frac{1}{2}} + \sqrt{N+y^2}}, \quad \dot{t} (= \frac{\dot{z}}{v})$$

$$= \frac{aN^{\frac{1}{2}}}{bM} y^{-1} \dot{y} \times \sqrt{N+y^2}, \quad \text{and } t = \frac{bz}{c^2} + \frac{\sqrt{c^2 - b^2}}{c^2} \times \sqrt{N+y^2} - \frac{a}{c}.$$

Whence it is evident, that, their first velocities being adapted accordingly, the tube and ball within it will so revolve about the center  $C$ , to which one end of the tube is fixed, that the ball will continually approach nearer and nearer to that center; yet never will arrive at it.

In other cases, if the ball first moves towards that center, it will, in a finite time, make its nearest approach thereto; and afterwards it will continually recede therefrom till it comes to the revolving end of the tube: or,  $c$  being greater than  $\frac{M}{N}^{\frac{1}{2}} \times b$ , it will, in a finite time, arrive at the center, with a velocity  $= \sqrt{c^2 - \frac{Mb^2}{N}}$ ,  $r$  being taken  $= a$ .

10. In general, we may, by means of our theorems, either find the requisite force, or forces, from the nature or equation of the curve and some circumstance respecting the motion of the projectile or the action of the force, or forces, thereon; or, from some such circumstance and a knowledge of the force, or forces, acting on the body, its trajectory may be found, and every thing else that may be required relative to its motion.

T 2

*Example*



*Example 1.* To find the force and its direction requisite (at every instant) to cause a body to revolve in a trajectory, so that it shall, in equal times, recede equal spaces from, and describe equal angles about, a given point; that is, that it shall describe the *spiral of Archimedes* with an invariable angular velocity?

In this case  $v$  and  $u$  being invariable, we have, by art. 1.  $f = \frac{b^2 y}{r^2}$  and  $g = \frac{2bc}{r}$ : therefore  $\sqrt{f^2 + g^2}$ , the force sought, must be  $= \frac{b}{r^2} \sqrt{4r^2 c^2 + b^2 y^2}$ ; and,  $f$  being to  $g$  as  $\frac{by}{2r}$  to  $c$ , if a perpendicular to the radius vector, drawn from the center  $C$  to the tangent to the trajectory at the point (B); where the body shall at any time be, be bisected by a right line drawn from that same point of contact, the line from  $B$  at right angles to that bisecting line will be the direction in which that force must act.

*Example 2.* To find the force, or forces, requisite (at every instant) to cause a projectile to describe an ellipsis with an invariable angular velocity about one of its foci?

In any elliptic orbit,  $v$  is  $= \frac{uy}{dr} \times \sqrt{e^2 - 1.y^2 + 2dy - d^2}$ ; the semi-axes being  $\frac{d}{1-e^2}$  and  $\frac{d}{\sqrt{1-e^2}}$ , and  $p = \frac{dy^{\frac{3}{2}}}{\sqrt{2d+e^2-1.y}}$ : therefore, first supposing  $u$ , the angular velocity about the focus  $C$ , to be invariable; we have  $v^2 = \frac{b^2 y^2}{d^2 r^2} \times \sqrt{e^2 - 1.y^2 + 2dy - d^2}$ ; and, by substitution,

$f =$

$$f = \frac{b^2 y}{r^2} - \frac{v\dot{v}}{j} = \frac{b^2 y^2}{dr^2} - \frac{2b^2 y}{d^2 r^2} \times \sqrt{e^2 - 1} \cdot y^2 + 2dy - d^2,$$

$$g = \frac{2b^2 v}{r} = \frac{2b^2 y}{dr^2} \times \sqrt{e^2 - 1} \cdot y^2 + 2dy - d^2,$$

$$F = \sqrt{f^2 + g^2} = \frac{b^2 y^2}{dr^2} \times \sqrt{1 + 2 \cdot \frac{1-e^2}{d^2} \cdot \frac{y'-y}{y} \times \sqrt{1-e^2} \cdot yy'' - d^2};$$

$y$  being the distance of the body from the focus C, and  $y''$  ( $= \frac{2d + e^2 - 1 \cdot y}{1 - e^2}$ ) its distance from the other focus.

Secondly: supposing C to be one focus of the orbit, and the angular velocity about the other focus to be invariable;  $u$  will be  $= \frac{b'' y''}{y}$  and  $v = \frac{b'' y''}{dr} \times \sqrt{1 - e^2} \cdot yy'' - d^2$ ,  $b''$  denoting the said angular velocity measured at the distance  $r$  from the respective focus: consequently, by substitution, we find in this case

$$f = \frac{u^2 y}{r^2} - \frac{v\dot{v}}{j} = \frac{b''^2 y''^2}{dr^2} - \frac{b''^2 y''}{d^2 r^2} \cdot \frac{y'' - y}{y} \times \sqrt{1 - e^2} \cdot yy'' - d^2,$$

$$g = \frac{v}{r} \cdot \frac{uy + 2uy}{j} = \frac{b''^2 y''}{dr^2} \cdot \frac{y'' - y}{y} \times \sqrt{1 - e^2} \cdot yy'' - d^2,$$

$$F = \sqrt{f^2 + g^2} = \frac{b''^2 y''^2}{dr^2} \times \sqrt{1 + 2 \cdot \frac{e^2 - 1}{d^2} \cdot \frac{y'' - y}{y''} \times \sqrt{1 - e^2} \cdot yy'' - d^2}.$$

*Remark.* It is observable that  $\sqrt{1 - e^2} \cdot yy'' - d^2$  is  $= 0$  when the body is at either end of the transverse axis, and  $y'' - y = 0$  when it is at either end of the conjugate axis, let the excentricity of the orbit be what it will: and

when the excentricity is small,  $\frac{y'' - y}{y} \times \sqrt{1 - e^2} \cdot yy'' - d^2$  is always a very small quantity; therefore it may then be rejected:

rejected as inconsiderable; and,  $g$  being at the same time inconsiderable, the only force to be considered, as affecting the motion of the body, will be  $f = \frac{b''^2 y''^2}{dr^2}$ , in our second case.

Now, if a body be made to revolve in the same ellipsis by the action of a centripetal force at  $C$ , that force will be  $= \frac{b^2 d^2}{(1+e)^2 \cdot r^2 y^2}$ ,  $b$  denoting the angular velocity of the body at the apsis nearest to the focus  $C$ ; and the angular velocity of the body about the other focus will, at the same time, be to  $b$  as  $1-e$  to  $1+e$ . Moreover, if  $b''$  be assumed in that proportion to  $b$ , our force  $f = \frac{b''^2 y''^2}{dr^2}$

will become  $= \frac{(1-e)^2 \cdot b^2 y''^2}{(1+e)^2 \cdot dr^2}$ ; which is to the centripetal

force  $\frac{b^2 d^2}{(1+e)^2 \cdot r^2 y^2}$ , just now mentioned, as  $y^2 y''^2$  to  $\frac{d^2}{(1-e^2)^2}$ .

But  $y^2 y''^2$  is  $= \frac{d^2}{(1-e^2)^2}$  when the body is at either apsis; and,

if the excentricity be small,  $e^2$  being then very small, and the focal distances  $y$  and  $y''$  differing very little from the semi-parameter  $d$ ,  $y^2 y''^2$  will always be nearly  $= \frac{d^2}{(1-e^2)^2}$ .

Therefore, in such an orbit, the motion of the body made to describe it by the action of the centripetal force

$\frac{b^2 d^2}{(1+e)^2 \cdot r^2 y^2}$ , which is reciprocally as the square of the distance of the body from the focus  $C$ , will be nearly the

same as the motion of a body revolving in the same orbit  
with

with the invariable angular velocity  $\frac{1-e}{1+e} \times b$  about the other focus.

Thus, without any regard to the area of the trajectory described, we see the physical ground of the celebrated *Bishop WARD's* method of determining the place of a Planet in its orbit: which method, it appears by what is here shewn, is not merely an hypothesis, as it is commonly called; nor is it indeed, in any elliptic orbit, strictly true; but it may properly be considered as an useful approximation in orbits nearly circular.

*Example 3.* To find how a ball will move within a straight slender tube, revolving in a vertical plane, with an invariable angular velocity, about an axis at the point C of the tube?

The tube being continued both ways from C; let the Fig. 41. circular arc  $z$ , whose radius is  $r$  and sine  $x$ , measure the angle of elevation BCH, (above the horizontal line CH,) of the ascending branch CB of the tube, in which the ball is supposed at first to move: Then, putting  $m = 16\frac{1}{2}$  feet, the force of gravity urging the ball towards the axis C

will be  $= \frac{2mx}{r}$ ; which will be  $= f = \frac{b^2 y}{r^2} - \frac{v\dot{v}}{j}$ ,  $b$  denoting the angular velocity of the tube about the said axis, measured at the distance  $r$  therefrom. Moreover  $i$  will be  $= \frac{z}{b} = \frac{j}{v}$ : therefore  $v$  will be  $= \frac{bj}{z}$ , and  $\dot{v} = \frac{b\dot{j}}{z}$ ,  $z$  being considered as invariable; and, by substituting accordingly, we have  $\frac{2mx}{r} = \frac{b^2 y}{r^2} - \frac{b^2 \dot{j}}{z^2}$ , and consequently  $\frac{2mrx\dot{z}^2}{b^2} = yz^2 - r^2 \ddot{y}$ .

Hence,

Hence, by multiplying by  $N^{\frac{x}{r}}$  and taking the fluents according to what is shewn in the Appendix, we get.

$$\frac{mr}{b^2} \times x N^{\frac{x}{r}} \dot{z} - r N^{\frac{x}{r}} \dot{x} = y N^{\frac{x}{r}} \dot{z} - r N^{\frac{x}{r}} \dot{y} - \frac{a-d}{b} \dot{z};$$

$d$  being  $= \frac{cr}{b} - \frac{mr^2}{b^2}$ ,  $N$  = the number whose hyperbolic logarithm is 1; and  $a$  and  $c$  being the values of  $y$  and  $v$  respectively, when  $x$  and  $z$  are each = 0,  $\dot{x} = \dot{z}$ , and  $y = \frac{cr}{b}$ ;  $c$  being considered as positive when the ball at first moves in the ascending branch of the tube so that  $y$  (then =  $a$ ) increases.

From which equation of the fluents, we have

$$r y N^{-\frac{x}{r}} - y N^{-\frac{x}{r}} \dot{z} = \frac{mr}{b^2} \times r x N^{-\frac{x}{r}} - x N^{-\frac{x}{r}} \dot{z} - \frac{a-d}{b} N^{-\frac{2x}{r}} \dot{z};$$

from whence, by again taking the fluents, we get

$$y N^{-\frac{x}{r}} = \frac{mr x N^{-\frac{x}{r}}}{b^2} + \frac{1}{2} \frac{a-d}{b} N^{-\frac{2x}{r}} + \frac{1}{2} \frac{a+d}{b}.$$

Consequently the general equation of the curve described by the ball is

$$y = \frac{mr x}{b^2} + \frac{1}{2} \frac{a-d}{b} N^{-\frac{2x}{r}} + \frac{1}{2} \frac{a+d}{b} N^{\frac{x}{r}};$$

and  $v$  will be  $= \frac{m\sqrt{r^2-x^2}}{b} - \frac{b}{2r} \frac{a-d}{b} N^{-\frac{x}{r}} + \frac{b}{2r} \frac{a+d}{b} N^{\frac{x}{r}}.$

*Remark 1.* If  $a$  be = 0, and  $c = \frac{mr}{b}$ , that is, if the ball begins to move from C, with the horizontal velocity

city  $\frac{mr}{b}$ , in the ascending branch of the tube; the equation of the curve described by the ball becomes  $y = \frac{cr}{b}$ . Which, answering to a circle, whose diameter is  $\frac{cr}{b}$ , touching the horizontal line passing through C, suggests this *remarkable inference*: the ball being put in motion at C, with the velocity  $\frac{mr}{b}$  as just now mentioned, it will revolve uniformly in a circle (whose diameter is  $\frac{cr}{b}$ ) standing upon the horizontal line with which the tube at first is supposed to coincide! and it will continue so to revolve (moving up and down alternately in the different branches of the tube) so long as the uniform motion of the tube is continued, making two complete revolutions whilst the tube makes one revolution! and the uniform velocity ( $\frac{mr}{b}$ ) wherewith the ball so revolves in such circle will be to its velocity along the tube every where as  $r$  to  $\sqrt{r^2 - x^2}$ .

That the ball, in the case adverted to in this Remark, will describe a circle, will, without finding the general equation of the trajectory, readily appear upon enquiring what angular velocity the tube must have, that the ball, when  $a$  is = 0, shall describe a given circle, or part of a given circle, touching the horizontal line at C. Thus:  $y$  being =  $2x$ , and  $v$  to  $u$  as  $\dot{y}$  to  $\frac{r\dot{y}}{\sqrt{4r^2 - y^2}}$ , by the nature of the circle, when  $r$  is the radius of the given one as well as of that upon which  $x$  is measured; we, by substituting properly in the equation  $\frac{2mx}{r} = \frac{u^2 y}{r^2} + \frac{v\dot{y}}{y}$  (=  $f$ ), have

U

 $\frac{my}{r}$

$$\frac{my}{r} = \frac{2u^2y}{r^2} - \frac{4u\dot{u}}{j} + \frac{u\dot{u}y^2}{r^2j};$$

whence it plainly appears, that  $u$  may be equal to the invariable quantity  $\left(\frac{mr}{2}\right)^{\frac{1}{2}}$  ( $= b$ ); which,  $c$  being  $= \sqrt{2mr}$  ( $= 2b$ ) is agreeable to what is said above.

Or, by taking the fluents, after multiplying by  $y$  and bringing the fluxionary equation into a convenient form, we find

$$u \text{ must be } = \frac{\sqrt{mry^4 - 8mr^2y^2 + 32b^2r^4}}{2^{\frac{1}{2}} \times 4r^2 - y^2};$$

$b$  being the initial value of  $u$ , and  $c = 2b$ .

If the velocity  $u$  were regulated according to this equation, and  $b$  were less than  $\left(\frac{mr}{2}\right)^{\frac{1}{2}}$ , the ball could only describe a part of the given circle; of which part, the chord would be  $= 2r \times \sqrt{1 - \frac{mr - 2b^2}{mr}}$ .

In a similar manner, we may find how the angular velocity of the tube must be regulated, that the ball shall describe a given ellipsis, or part of a given ellipsis, or other figure, touching the horizontal line at C.

*Remark 2.* If  $a$  be  $= \frac{cr}{b} - \frac{mr^2}{b^2}$ , our equation of the trajectory becomes

$$y = aN^{\frac{2}{r}} - \frac{a}{r} - \frac{c}{b} \cdot x;$$

and  $v$  will then be  $= aN^{\frac{2}{r}} - \frac{a}{r} - \frac{c}{b} \cdot \sqrt{r^2 - x^2} \times \frac{b}{r}$ .

*Remark*

*Remark 3.* If  $a$  be  $= \frac{nr^2}{b^2} - \frac{cr}{b}$ , the equation of the spiral described by the ball will be

$$y = aN^{-\frac{n}{r}} + \frac{a}{r} + \frac{c}{b} \cdot x;$$

$$\text{and } v \text{ will be } = \frac{a}{r} + \frac{c}{b} \cdot \sqrt{r^2 - x^2} - aN^{-\frac{n}{r}} \times \frac{b}{r}.$$

In this case, supposing both  $a$  and  $c$  positive, the ball will revolve in a spiral B'B''B''' &c. having a circular asymptote CD; and will make two revolutions in such spiral whilst the tube makes one revolution: which revolutions in the spiral will be alternately without and within the circle whose diameter is  $a + \frac{cr}{b}$ , standing upon, and touching at the point C, the horizontal line passing through that point; to the periphery of which circle, the ball, going alternately into and out of it at the said point C, will continually approach nearer and nearer, but never can absolutely revolve in it!

In each of the spirals which, in a preceding article, we found to have circular asymptotes, the describing projectile either keeps always without or always within the asymptotic circle. The spiral here described is perhaps still of a more extraordinary nature, being in its revolutions alternately without and within its asymptotic circle; and, at its ingress and egress, always intersecting the periphery of that circle in the same point C. And this spiral is a curve, that (abstracting from resistance within the tube) a ball would actually be made to describe by the force of gravity, whilst carried about C by the uniform



motion of the tube revolving in a vertical plane as above explained.

The spiral described by the ball will be of the same kind, when  $a$  and  $c$  are not both positive, provided  $a + \frac{cr}{b}$  be equal to the positive quantity  $\frac{mr^2}{b^2}$ ; but it may sometimes make more than one revolution before it enters the asymptotic circle.

*Remark 4.* The angular velocity about C being invariable, the force  $B \times g$  wherewith the tube, by the nature of the motion, would necessarily urge the ball in a direction at right angles thereto, if gravity did not act,

$$\text{is} = \frac{2Bbv}{r} = 2Bm \cdot \frac{\sqrt{r^2 - x^2}}{r} + \frac{Bb^2}{r^2} \cdot a + d \cdot N^{\frac{x}{r}} - \frac{Bb^2}{r^2} \cdot a - d \cdot N^{-\frac{x}{r}} :$$

to which adding  $2Bm \cdot \frac{\sqrt{r^2 - x^2}}{r}$ , the pressure against the tube arising from the gravity of the ball B, we have

$$4Bm \cdot \frac{\sqrt{r^2 - x^2}}{r} + \frac{Bb^2}{r^2} \cdot a + d \cdot N^{\frac{x}{r}} - \frac{Bb^2}{r^2} \cdot a - d \cdot N^{-\frac{x}{r}},$$

the quantity expressing the whole pressure arising from the gravity of the ball and the mutual action or re-action of the ball and tube against each other.

Which pressure, it may be observed, is  $= 2Bm + \frac{2Bbc}{r}$ , the instant the tube is moved from a horizontal position; and it is the same let  $a$  be what it will: whereas the pressure, if the tube rested in that position, would be only  $2Bm$ ; which likewise will be the pressure when it is first moved, if  $c$  be  $= 0$ .

*Remark*

*Remark 5.* If the ball be supposed to descend from the point C along a revolving plane, (instead of being included in a tube,) whilst the plane itself moves uniformly about that point, from a horizontal position; the pressure against the plane will, by taking  $a$  and  $c$  each = 0, be found by

our theorem =  $Bm \times \frac{4\sqrt{r^2 - x^2}}{r} - N^{\frac{x}{r}} - N^{-\frac{x}{r}}$ . Which

being = 0 when the angle measured by the arc  $x$  is =  $47^\circ 11' 54''$ ; it follows that the ball so descending will quit the plane when, by revolving as just now mentioned, it comes to make that angle ( $47^\circ 11' 54''$ ) with the horizon, let the invariable angular velocity of the plane be what it will.

*Remark 6.* Retaining the plane (continued above the axis) instead of the tube, if  $a$  be =  $\frac{mr^2}{b^2}$ , and  $c = 0$ ; that is, if the ball be laid at a distance =  $\frac{mr^2}{b^2}$  from the axis of motion, upon the ascending part of the plane, afterwards made to revolve (from a horizontal position) as above; the

pressure will be =  $2Bm \times \frac{2\sqrt{r^2 - x^2}}{r} - N^{-\frac{x}{r}}$ . Which will

be = 0 when  $2\sqrt{r^2 - x^2}$  is =  $rN^{-\frac{x}{r}}$ ; that is, when the angle measured by the arc  $x$  becomes =  $83^\circ 17' 21''$ ; and consequently the ball, at that instant, will quit the plane,  $y$  being at the same time =  $1.2268 \times a$ : the ball at first ascending up the plane, and continuing to do so till  $y$  be-

comes =  $1.2361 \times a$ ;  $\sqrt{r^2 - x^2}$  being then =  $rN^{-\frac{x}{r}}$ , and the inclination of the plane to the horizon  $74^\circ 3' 58''$ .

*Remark*

Fig. 43. *Remark 7.* If the straight slender tube DBPE be fastened to the lever CP, and the tube with a ball in it be made to revolve about the axis C, in a vertical plane, so that the angular velocity of CP shall be invariable; we may, in such case, find how the ball will move, by observing that its relative velocity from the point P within the tube will be the same, let the length of the lever CP (which we will suppose at right angles to the tube) be what it will: provided the angular velocity of CP, the first position of the tube, the first distance of the ball from P, and its first velocity towards D, be respectively the same.

Accordingly, having respect to a tube so moved, and supposing  $b$  now to denote the velocity of a point in CP at the distance  $r$  from C;  $z$  the circular arc described by such point;  $x$  the sine of that arc; Y the variable distance of the ball from P; A the value of Y when  $x$  and  $z$  are each = 0; V the variable relative velocity of the ball within the tube towards the end D thereof; and C its relative velocity, towards the same end, at the commencement of the motion, when the tube is supposed to be parallel to the horizontal line CH:

$$Y = \sqrt{y^2 - P^2} \text{ will be } = \frac{mr^x}{b^x} + \frac{1}{2} \cdot A + D \cdot N^{\frac{x}{r}} + \frac{1}{2} \cdot A - D \cdot N^{-\frac{x}{r}};$$

$$V = \frac{b\dot{Y}}{\dot{x}} = \frac{m\sqrt{r^2 - x^2}}{b} + \frac{b}{2r} \cdot A + D \cdot N^{\frac{x}{r}} - \frac{b}{2r} \cdot A - D \cdot N^{-\frac{x}{r}};$$

$$v, \text{ the velocity of the ball from the center C, } = \frac{\sqrt{y^2 - P^2}}{y} \times V;$$

U, the absolute velocity of the ball in the curve it will describe,

$$= \frac{\sqrt{b^2 y^2 + r^2 V^2 - 2brPV}}{r};$$

and

and the pressure =  $\frac{2m\sqrt{r^2-x^2}}{r} + \frac{2bV}{r} - \frac{b^2P}{r^2} \times B$ ;

D being =  $\frac{rC}{b} - \frac{mr^2}{b^2}$ , and P = the perpendicular CP.

*Remark 8.* If A be = 0 and  $C = \frac{mr}{b}$ ,  $y^2$  will be =  $P^2 + \frac{C^2x^2}{b^2}$ : and the ball will describe a geometrical oval, which will be a line of the fourth order, whose equation is

$$W^2 = \frac{1}{2}P^2 + RX - X^2 \pm \frac{1}{2}P\sqrt{P^2 + 4RX - 4X^2};$$

X being the abscissa, measured upwards from C on a line at right angles to CH, and W the correspondent ordinate parallel to CH; and R being put for  $\frac{rC}{b} = \frac{mr^2}{b^2}$ .

The oval described by the ball will be as in Figure 44, 45, or 46; according as R is greater, equal to, or less than P; R, when less than P, being greater than  $\frac{1}{2}P$ . Fig. 44.  
Fig. 45.  
Fig. 46.

If R be less than  $\frac{1}{2}P$ , the oval will be every where concave towards C.

*Remark 9.* If A + D be = 0, the ball will describe a spiral whose equation is  $y^2 = P^2 + \frac{C^2x^2}{b^2} + \frac{2ACxN^{-\frac{x}{r}}}{b} + A^2N^{-\frac{2x}{r}}$ ; of which spiral, one or other of the ovals mentioned in the preceding Remark will be an asymptote.

*Remark 10.* It is obvious that, by making the tube revolve uniformly in some certain plane between the horizontal and vertical, the sine of whose inclination to the horizon (to the radius  $r$ ) shall be  $s'$ , we may make the ball describe a curve whose equation is

$$Y = \frac{ms'x}{b^2} + \frac{1}{2} \cdot \overline{A + D} \cdot N^{\frac{x}{r}} + \frac{1}{2} \cdot \overline{A - D} \cdot N^{-\frac{x}{r}};$$

where

where  $\frac{ms'}{b^2}$ , the coefficient of  $x$ , may be of any value between 0 and  $\frac{mr}{b^2}$ ; and where  $D$  is  $= \frac{rC}{b} - \frac{ms'}{b^2}$ .

*Remark 11.* If  $s'$  be taken = 0, the equation of the curve which the ball will describe becomes

$$Y = \sqrt{y^2 - P^2} = \frac{1}{2}.A + \frac{rC}{b}.N^{\frac{x}{r}} + \frac{1}{2}.A - \frac{rC}{b}.N^{-\frac{x}{r}};$$

$$\text{and } V \text{ will be } = \frac{bA + rC}{2r}.N^{\frac{x}{r}} - \frac{bA - rC}{2r}.N^{-\frac{x}{r}}.$$

Which equations relate to the curve described by a ball within a straight tube at the lever  $CP$ , revolving in a horizontal plane about the center  $C$ .

And if  $Ab - rC$  be = 0, the equation of the curve becomes

$$y^2 = P^2 + A^2 N^{\frac{2x}{r}}.$$

*Remark 12.* If,  $s'$  being = 0,  $C$  be  $= -\frac{Ab}{r}$ ; that is, if the ball at first moves towards the point  $P$  of the tube, with the relative velocity  $\frac{Ab}{r}$ ; our equation of the curve described by the ball becomes

$$y^2 = P^2 + A^2 N^{-\frac{2x}{r}}; \text{ and } V \text{ will be } = -\frac{AN^{-\frac{x}{r}}}{r}.$$

By which it appears that the ball, in that case, will revolve about the center  $C$  in a spiral, continually approaching nearer and nearer to the circumference of the circle whose center is  $C$  and radius =  $P$ , yet never will arrive at it.

In

In other cases, (the tube revolving in a horizontal plane,) if the ball at first moves towards the point P of the tube with a relative velocity less than  $b$ ,  $r$  being taken  $= A$ , it will, in a finite time, make its nearest approach to that point; (when  $Y$  will be  $= \frac{A\sqrt{b^2 - C^2}}{b}$ ;) and afterwards it will continually recede therefrom: or,  $C$  being less than  $-b$ , the ball will, in a finite time, arrive at P with the relative velocity  $\sqrt{C^2 - b^2}$ ; and, after touching the circumference of the circle just now mentioned, continually recede from it.

*Remark 13.* The relation between ( $y$ ) the radius vector Fig. 47. and ( $V$ ) the relative velocity of the ball moving in any straight tube, revolving uniformly in a horizontal plane about  $C$ , being (by what is said above) expressed by the equation  $b^2 y^2 = r^2 V^2$ , let the perpendicular  $P$  and the angle made by the tube and radius vector be what they will; it follows, that the same equation will express the relation between the radius vector and the relative velocity of a ball moving in any curved tube, carried about the center  $C$  in a horizontal plane so that the angular velocity of the lever  $CQ$  (to which we suppose the curved tube to be fastened) shall be expressed by the invariable quantity  $b$ . From which equation, by taking the fluents, we get

$$b^2 \cdot y^2 - a^2 = r^2 \cdot V^2 - C^2;$$

and consequently  $V = \sqrt{C^2 + \frac{b^2}{r^2} \cdot y^2 - a^2}$ ;

$C$  being the value of  $V$  when  $y$  is  $= a$ .

X

There-

$$\text{Therefore } \dot{z} = \frac{b\dot{Z}}{V} \text{ will be } = \frac{b\dot{Z}}{\sqrt{C^2 + \frac{b^2}{r^2}y^2 - a^2}}$$

$\dot{Z}$  denoting the fluxion of the curve  $QB$ , or (which is the same thing) of the relative space passed over by the ball within such curved tube.

Whence the trajectory of the ball may be found, and its motion therein;  $\dot{Z}$  being given in terms of  $y$  and  $\dot{y}$ .

**Fig. 48.** *Remark 14.* If the form of the tube be the circumference of a circle whose radius is  $r$ , and  $C$  be a point therein,  $\dot{Z}$  will be  $= \frac{2r\dot{y}}{\sqrt{4r^2 - y^2}}$ ; and consequently

$$\dot{x} = \frac{2b\dot{r}\dot{y}}{\sqrt{4r^2 - y^2} \times \sqrt{C^2 + \frac{b^2}{r^2}y^2 - a^2}}$$

Therefore, by what is shewn in Table XII. of the Appendix, the motion of a ball in a circular tube, revolving uniformly in a horizontal plane about a point in such tube, will be determined by means of elliptic arcs.

When, supposing  $a=0$ ,  $C$  (the velocity of the ball at the point  $C$ ) is  $=2b$ ; the time in which the tube will make one revolution, about the axis  $C$ , will be to the time in which the ball will make one revolution in the tube, as  $4q$  to  $e + \sqrt{e^2 - 2q}$ ; that is, as 3.14159265 to 1.3110287771;  $q$  denoting the quadrantal arc of a circle whose radius is 1, and  $e$  the quadrantal arc of an ellipse whose semi-axes are  $2\frac{1}{2}$  and 1.

*Example 4.* Let the body  $B'$  describe any trajectory whatever about the center  $C$ , by means of known forces ( $f$  and  $g$ ) acting thereon; and let another body  $B''$  describe

scribe such a trajectory about the same center, that the bodies shall always be equidistant therefrom, and the angular velocity of the former to that of the latter (about that center) as  $u$  to  $e + mu$ : to find the force, or forces, requisite to retain the body  $B''$  in such a trajectory,  $e$  and  $m$  being invariable quantities?

The forces acting on  $B'$  being

$$f = \frac{u^2 y}{r^2} - \frac{v\dot{v}}{j}, \text{ and } g = \frac{v}{r} \times 2u + \frac{\dot{u}y}{j};$$

those acting on  $B''$  will be

$$f'' = \frac{(e + mu)^2 y}{r^2} - \frac{v\dot{v}}{j}, \text{ and } g'' = \frac{v}{r} \times 2.e + mu + \frac{m\dot{u}y}{j};$$

the additional forces therefore must be

$$f'' - f = \frac{e^2 + 2em + m^2 - 1.u^2}{r^2} \times y,$$

$$\text{and } g'' - g = \frac{v}{r} \times 2.e + m - 1.u + \frac{m - 1.\dot{u}y}{j}.$$

*Remark 1.* If  $e$  and  $g$  be each = 0, the only additional force requisite will be  $f'' - f = \frac{m^2 - 1.u^2 y}{r^2} = \frac{m^2 - 1.a^2 b^2}{r^2 y^3};$

$u$  being then =  $\frac{a^2 b}{y^2}$ , by art. 7.

which agrees with the conclusion deduced by Sir ISAAC NEWTON and others, relative to the requisite additional force in this case.

*Remark 2.* If  $m$  be = 1 and  $g = 0$ ,

$$f'' - f \text{ will be } = \frac{e^2 + 2eu}{r^2} \times y = \frac{e^2 y}{r^2} + \frac{2a^2 b e}{r^2 y}, \text{ and } g'' - g = \frac{2ev}{r};$$

$u$  being as in the preceding Remark.

*Example 5.* Let the projectile be supposed to move near the earth's surface, in a medium, the resistance whereof



is  $MDU^n$ ;  $D$  denoting the density of the medium,  $M$  some invariable quantity, and  $U$  the absolute velocity of the body in the path it may describe.

Fig. 49. Then the resistance in a direction parallel to the horizontal base  $ACD$  will be  $= \frac{MDU^n \dot{x}}{z}$ ; and in a direction perpendicular to the same base  $= \frac{MDU^n \dot{y}}{z}$ ;  $x$ ,  $y$ , and  $z$ , respectively denoting the abscissa  $AC$  (parallel to the horizon), the ordinate  $CB$  (at right angles thereto,) and the length  $AB$  of the trajectory. Therefore, by art. 5.

$$f' = -\frac{v\dot{v}}{j} \text{ will be } = 2m + \frac{MDU^n \dot{y}}{z} = 2m + \frac{MDv^n z^{n-1}}{j^{n-1}}$$

$$g' = \frac{v\dot{u}}{j} \dots\dots = -\frac{MDU^n \dot{x}}{z} = -\frac{MDvu^{n-1} z^{n-1}}{z^{n-2} j}$$

$$\frac{\dot{x}}{U} = \frac{\dot{y}}{u} = \frac{j}{v} = i;$$

$v$  being the velocity of the body from the base  $ACD$ ,  $u$  its velocity in a direction parallel to the said base, and  $2m$  the accelerative force of gravity in the medium.

Whence, by exterminating  $v$  and  $\dot{v}$ , we have  $2m\dot{x}^2 = -u^2\dot{y}$ ;  $\dot{x}$  being considered as invariable.

From which equations every thing relative to the trajectory, and the motion of the projectile therein, may be determined.

*Remark 1.* If  $n$  be  $= 1$ , and  $MD$  equal to the invariable quantity  $d$ ; we have, from the values of  $g'$ ,  $\dot{u} = -d\dot{x}$ : whence, by taking the fluents, we find  $u = b - dx$ ,  $b$  being the value of  $u$  when  $x$  is  $= 0$ . Consequently we have, by sub-

substitution,  $2m\dot{x}^2 = -\overline{b-dx}^2 \ddot{y}$ , or  $\ddot{y} = -\frac{2m\dot{x}^2}{\overline{b-dx}^2}$ .

Hence, by taking the fluents, we have  $\dot{y} = e\dot{x} - \frac{2m\dot{x}}{d\overline{b-dx}}$ ;

and, by taking the fluents again,  $y$  is found  $= e\dot{x} + \frac{2m}{d^2} \times$

Log.  $\frac{b-dx}{b}$ ;  $e$  being  $= \frac{c}{b} + \frac{2m}{bd}$ , where  $c$  denotes the initial vertical velocity of the projectile.

*Remark 2.* If  $n$  be  $= 2$ , and MD as in the preceding remark; we have, from the values of  $g'$ ,  $\frac{\dot{u}}{u} = -d\dot{x}$ : whence, by taking the fluents, we get Log.  $\frac{u}{b} = -dx$ ; and consequently  $u = bN^{-dx}$ , N being the number whose hyperbolic logarithm is 1. Therefore, by substitution, we have  $2m\dot{x}^2 = -b^2\ddot{y}N^{-2dx}$ . Let  $s\dot{x}$  be  $= \dot{z} = \sqrt{\dot{x}^2 + \dot{y}^2}$ : then will  $\dot{y}$  be  $= \dot{x}\sqrt{s^2-1}$ ,  $\ddot{y} = \frac{ss\dot{x}}{\sqrt{s^2-1}}$ ,  $\frac{2m\dot{x}}{s} = -\frac{b^2ss}{\sqrt{s^2-1}} \times N^{-2dx}$ , and  $\frac{b^2s^2s}{\sqrt{s^2-1}} = -2m\dot{x}N^{-2dx}$ .

Hence  $\frac{b^2}{2} \times s\sqrt{s^2-1} + \text{Log.} \frac{s + \sqrt{s^2-1}}{q + \sqrt{q^2-1}} = \frac{m}{d}N^{-2dx} - \frac{m}{d}$ ;

$q$  being  $= \frac{C}{b}$ , where C denotes the initial absolute velocity of the projectile.

Now, from the equation so found, expressing the relation of  $s$  and  $z$ , the value of this last mentioned quantity may be found in terms of  $s$ ; and consequently the values of  $\dot{x}$  ( $= \frac{\dot{z}}{s}$ ) and  $\dot{y}$  ( $= \frac{\dot{z}\sqrt{s^2-1}}{s}$ ) will be had in terms of  $s$  and  $z$ : from whence the values of  $x$  and  $y$  may be found in terms of  $z$ .

---



---

M E M O I R E VHL.

*Of the Motion of a Body in (or upon) a Spherical Surface; in (or upon) which it is retained by some Force urging it towards the Center of the Sphere, whilst it is continually impelled by some other Force, or Forces, to change its Direction in (or upon) that Surface.*

Fig. 50. 1. **L**ET the body B be supposed to describe the curve Bb in (or upon) the spherical surface CPBbd, wherein C is a given point: let the sine of the arc CB (part of a great circle passing through C and B) be denoted by  $y$ : let  $B \times f$  denote a motive force continually urging the body towards C, in the direction of the tangent to the said arc CB at the point B, where the body is supposed to be: let  $B \times g$  denote another such force always acting on it, in a direction Bd at right angles to such tangent, and to the plane of the great circle CB: let  $v$  denote the velocity of the body from C, in the variable direction of the said tangent;  $u$  the angular velocity of a plane considered as revolving with the point B about the diameter from C, measuring such velocity at the distance  $r$  (= the radius of the sphere) from the sphere's center; and  $w$  the velocity of the body at B, in the direction Bd. Then, BP being a great circle touching the track (Bb) of the body in B; if the sine of the arc CP (of another great circle) be denoted by  $\rho$ , and the sine of the angle CPB (to the radius

radius 1) by  $s$ ; the sine of the angle CBP will be  $= \frac{p's}{y}$ , and the absolute velocity of the body in the direction (Bb) of the tangent to its track at B will be  $= \frac{wy}{p's}$ : which would continue invariable, and the body would describe the great circle PBb, if no force acted thereon but that which, by urging it towards the center of the sphere, would be requisite to retain it in (or upon) the surface thereof. In which case, the great circles BP and CP keeping their positions,  $p'$  and  $s$  would remain invariable; and  $\frac{wy + wj}{p's}$ , the fluxion of the velocity of the body, would be  $= 0$ : therefore  $w$  would then be  $= -\frac{wj}{y} = -\frac{uj}{r}$ ,  $w$  being  $= \frac{uj}{r}$ . But the forces  $B \times f$  and  $B \times g$  acting on the body,  $\dot{w}$  will be  $= \frac{\dot{u}y + uj}{r}$ . Consequently, the directions in which those forces act being at right angles to each other,  $\frac{\dot{u}y + 2uj}{r}$  (the excess of  $\frac{\dot{u}y + uj}{r}$ , the fluxion of  $w$  when those forces act, above  $\frac{-uj}{r}$ , which would be the fluxion of  $w$  if those forces ceased to act) will be the fluxion of the velocity generated or destroyed by the action of the force  $B \times g$  only: and, the motive force into  $\left(\frac{rj}{v\sqrt{r^2 - j^2}}\right)$  the fluxion of the time being equal to the fluxion of the quantity of motion generated or destroyed by that force,  $\frac{Bgrj}{v\sqrt{r^2 - j^2}}$  will be  $= B \times \frac{\dot{u}y + 2uj}{r}$ , and  $g = \frac{v\sqrt{r^2 - j^2}}{r^2} \times \frac{\dot{u}y + 2uj}{j}$ .

Moreover  $\sqrt{v^2 + w^2}$ , the velocity of the body in its track Bb, would be invariable if the forces  $B \times f$  and  $B \times g$  ceased!

ceased acting; and it would then (as observed above) describe a great circle: therefore  $\frac{v\dot{v} + w\dot{w}}{\sqrt{v^2 + w^2}}$  (the fluxion of that velocity) would, in that case, be = 0; and  $\dot{v} = -\frac{w\dot{w}}{v} = \frac{u^2 y \dot{y}}{r^2 v}$ ;  $w$  being then =  $\frac{uy}{r}$  and  $\dot{w} = -\frac{u\dot{y}}{r}$ , by what is said above. But the said forces continuing to act on the body, the fluxion of  $v$  will in general be expressed by  $\dot{v}$ . Consequently (the forces acting in directions at right angles to each other)  $\frac{u^2 y \dot{y}}{r^2 v} - \dot{v}$  will express the fluxion of the velocity destroyed or generated by the motive force  $B \times f$ ; and therefore ( $f$ ) the retarding or accelerating force produced by the action of the said motive force (being equal to the last mentioned fluxion divided by the fluxion of the time) will be =  $\frac{\sqrt{r^2 - y^2}}{r} \times \frac{u^2 y \dot{y} - r^2 v \dot{v}}{r^2 j}$ ,  $\frac{r \dot{y}}{v \sqrt{r^2 - y^2}}$  expressing the fluxion of the time, as before observed.

The forces

$$f \left( = \frac{\sqrt{r^2 - y^2}}{r} \times \frac{u^2 y \dot{y} - r^2 v \dot{v}}{r^2 j} \right) \text{ and } g \left( = \frac{v \sqrt{r^2 - y^2}}{r^2} \times 2u + \frac{u \dot{y}}{j} \right),$$

according to the method pursued in art. 2. of the preceding Memoir, may be resolved into two others  $h$  and  $k$ ; the former in the direction of the tangent to the track of the body, retarding or accelerating its absolute velocity; and the latter in a direction at right angles to such tangent, changing the direction of the body. Which forces  $h$  and  $k$  being to the forces denoted by the same symbols in the article just now mentioned manifestly as  $\sqrt{r^2 - y^2}$  to  $r$  respectively, I omit specifying their values particularly,

particularly, and proceed to explain the use of the theorems already investigated.

2. If only the force  $B \times f$  act on the body:  $g$  being  $= 0$ , we have  $uy + 2uy = 0$ ; whence we get  $uy^2 = a^2b$ ,  $b$  being the value of  $u$  when  $y$  is  $= a$ . Therefore,  $u$  being  $= \frac{a^2b}{y^2}$ ,  $w (= \frac{uy}{r})$  will be  $= \frac{a^2b}{ry}$ ; and (U) the absolute velocity of the body  $= \frac{\sqrt{4b^2 + r^2v^2y^2}}{ry}$ . Moreover the force

$$f \left( = \frac{\sqrt{r^2 - y^2}}{r} \times \frac{u^2y}{r^2} - \frac{v\dot{v}}{j} \right) \text{ will be } = \frac{\sqrt{r^2 - y^2}}{r} \times \frac{a^4b^2}{r^2y^4} - \frac{v\dot{v}}{j}$$

$$= \frac{\sqrt{r^2 - y^2}}{r} \times \frac{a^4b^2}{r^2y^4} - \frac{a^4b^2r^2}{2j} \times \text{the fluxion of } \frac{y^{-4}j^2}{z^2, r^2 - y^2},$$

$$v \text{ being } = \frac{\pi u j}{z \sqrt{r^2 - y^2}} = \frac{a^2 b r y^{-2} j}{z \sqrt{r^2 - y^2}};$$

where  $z (=ur)$  denotes the fluxion of the spherical angle BCP, the arc CP being supposed to keep its position.

*Example 1.* If the body describes a circle whose radius is  $a$ , in a plane at right angles to the diameter of the sphere drawn from C,  $v$  will be  $= 0$ , and  $y$  always  $= a$ : consequently the force  $f$  must, in that case, be  $= \frac{a^4b^2\sqrt{r^2 - a^2}}{r^3} = \frac{a^2\sqrt{r^2 - a^2}}{ar}$ ,  $d$  denoting the invariable velocity of the body in the circle it describes.

Which theorem is of considerable use in enquiries relative to the motion of a point, or body, in (or upon) the  
Y
surface

surface of a sphere, as will appear in a subsequent Memoir.

*Example 2.* Let the body be supposed to describe a *loxodromic*, making an angle, whose sine (to the radius 1) is  $s$ , with the great circles passing through the pole C.

In which case  $v$  will be  $= \frac{\sqrt{1-s^2} \cdot uy}{rs} = \frac{a^2 b \sqrt{1-s^2}}{rsy}$ , and

$\dot{v} = -\frac{a^2 b \sqrt{1-s^2} \times j}{r s y^2}$ : consequently the force  $f$  must then

$$\text{be} = \frac{\sqrt{r^2-y^2}}{r} \times \frac{a^4 b^2}{r^2 y^3} + \frac{a^4 b^2 \cdot 1-s^2}{r^2 s^2 y^3} = \frac{a^4 b^2 \sqrt{r^2-y^2}}{r^2 s^2 y^3}.$$

This case coincides with that in the preceding example, if  $s$  be taken equal to 1;  $y$  being then always equal to the invariable quantity  $a$ .

*Example 3.* One end of a string (whose length is  $r$ ) being fastened to an immoveable point, let a small ball be fastened to the other end thereof; and, after stretching the string straight out in any direction, and putting the ball in motion in a direction making a given angle with the horizon, let it be left to move in such a track as the action of its gravity and the tension of the string shall cause it to describe.

Then, supposing C to be the lowest point of the spherical surface in which the ball will move, and putting  $2m$  ( $= 32\frac{1}{2}$  feet) for the accelerative force of gravity,

$$f \text{ will be } = \frac{2my}{r} = \frac{\sqrt{r^2-y^2}}{r} \times \frac{a^4 b^2}{r^2 y^3} - \frac{v\dot{v}}{j}:$$

from which equation we have  $v\dot{v} = \frac{a^4 b^2 j}{r^2 y^3} - \frac{2myj}{\sqrt{r^2-y^2}}$ ;

and hence, by taking the fluents,

$$v =$$

$$v = \sqrt{4m \times \sqrt{r^2 - y^2} - \sqrt{r^2 - a^2} - \frac{a^2 b^2}{r^2 y^2} + c^2 + \frac{a^2 b^2}{r^2}};$$

$a$ ,  $b$ , and  $c$  being the respective values of  $y$ ,  $u$ , and  $v$ , at the commencement of the motion.

Therefore

$$\frac{-rj}{v \sqrt{r^2 - y^2}} = t \text{ will be } = \frac{-r \sqrt{r^2 - y^2}^{-\frac{1}{2}} \times j}{\sqrt{4m \times \sqrt{r^2 - y^2} - \sqrt{r^2 - a^2} - \frac{a^2 b^2}{r^2 y^2} + n^2}}$$

$n^2$  being put for  $c^2 + \frac{a^2 b^2}{r^2}$ ;

or, substituting  $x$  for  $\sqrt{r^2 - y^2} - \sqrt{r^2 - a^2}$ ,

$$t \text{ will be } = \frac{r^2 x}{\sqrt{4mr^2 x + n^2 r^2 \times a^2 - 2\sqrt{r^2 - a^2} \times x - x^2 - a^2 b^2}}$$

and, by our Appendix,  $t$  will always be assigned by means of the arcs of the conic sections.

*Remark 1.* If  $c$  be = 0,  $t$  will be

$$= \frac{\frac{1}{2} r m^{-\frac{1}{2}} x^{-\frac{1}{2}}}{\sqrt{a^2 - \frac{a^2 b^2 \sqrt{r^2 - a^2}}{2mr^2} - \frac{a^2 b^2 + 8mr^2 \sqrt{r^2 - a^2}}{4mr^2} \cdot x - x^2}}$$

and the value of  $t$  may always be assigned by means of elliptic arcs; or, by what is done in Mem. III. the time  $t$  may be compared with the time of descent of a common pendulum in a circular arc.

Let  $R$  denote the length of such a pendulum,  $D$  the vertical height from which its bob descends,  $x$  its vertical descent, and  $T$  the time of its descent:

$$\text{then will } t \text{ be } = \frac{rT}{R};$$

Y 2

R being



R being  $= \sqrt{P^2 + Q^2}$ , and  $D = \sqrt{P^2 + Q^2} - Q$ ,

when  $a^2 = \frac{a^2 b^2 r - b}{2mr^2}$  the value of  $P^2$  is positive;

or  $R = \frac{Q + \sqrt{P^2 + Q^2}}{2}$ , and  $D = Q - \sqrt{P^2 + Q^2}$ ,

when,  $\left(\frac{ab}{r}\right)$  the initial velocity of the ball being greater than  $a \cdot \frac{2m}{r-b}$ ,  $P^2$  is negative;

Q being put for  $\frac{a^2 b^2}{8mr^2} + r - h = \frac{4r^2 - 3a^2 - P^2}{4r - b}$ ,

and  $h$  for  $r - \sqrt{r^2 - a^2}$ , the height of the ball at the commencement of the motion above the horizontal plane touching the spherical surface at its lowest point C.

The ascent and descent of the ball will, it is obvious, be limited by two horizontal circles; in which the places of the apsides will continually vary: but

I have not found that  $x$ , the fluent of  $\frac{a^2 b r y^{-2} y}{v \sqrt{r^2 - y^2}} = u t$ , (which measures, on a horizontal plane, the angle described by the ball about the vertical diameter of the surface in which it moves,) can be assigned without having recourse to an infinite series, or to some curve of a higher kind than the conic sections.

*Remark 2.* If,  $c$  being  $= 0$ , the initial velocity  $\left(\frac{ab}{r}\right)$  of the ball, in a horizontal direction at right angles to the string, be  $= \sqrt{8m \cdot h - r}$ ,

$$t \text{ will be } = \frac{\frac{1}{2} m^{-\frac{1}{2}} r x^{-\frac{1}{2}}}{\sqrt{4r^2 - 3a^2 - x^2}}$$

And

And it follows, that the time of revolving from apsis to apsis (that is, from the highest point of its track to the lowest, or from the lowest to the highest) will be to the time of descent (of another body) in the quadrantal arc, whose lowest point is C and radius  $r$ , as  $r^{\frac{1}{2}}$  to  $\sqrt{4r^2 - 3a^2}^{\frac{1}{2}}$ : which time of descent will be  $= \frac{r^{\frac{1}{2}}}{m^{\frac{1}{2}}} \times \frac{e + \sqrt{e^2 - 2g}}{2} = 1.31102877 \times \frac{r^{\frac{1}{2}}}{m^{\frac{1}{2}}}$ , by Mem. III.  $q$  denoting the quadrantal arc of a circle whose radius is 1, and  $e$  the quadrantal arc of an ellipsis whose semi-axes are  $2^{\frac{1}{2}}$  and 1.

*Remark 3.* If  $b$  be  $= 0$ , the ball will move in a vertical plane, and

$$t \text{ will be } = \frac{\frac{1}{2} m^{-\frac{1}{2}} r x}{\sqrt{\frac{e^2}{4m} + x \times 2r - b + x \times b - x}}$$

And it appears by our Tables, that  $t$  will then be assigned by means of elliptic arcs: except  $c$  be  $= \sqrt{4m \cdot 2r - h}$ ; when,

$$t \text{ being } = \frac{\frac{1}{2} m^{-\frac{1}{2}} r x}{2r - b + x \cdot \sqrt{b - x}},$$

$$t \text{ will be } = \frac{r^{\frac{1}{2}}}{2^{\frac{1}{2}} m^{\frac{1}{2}}} \times \text{Log. } \frac{\sqrt{2r - \sqrt{b}}}{\sqrt{2r + \sqrt{b}}} \times \frac{\sqrt{2r + \sqrt{b - x}}}{\sqrt{2r - \sqrt{b - x}}}$$

*Remark 4.* If,  $b$  being  $= 0$ ,  $c$  be  $= \sqrt{8mr}$  and  $h = 2r$ ;

$$t \text{ will be } = \frac{\frac{1}{2} m^{-\frac{1}{2}} r x^{-\frac{1}{2}}}{\sqrt{4r^2 - x^2}};$$

and the ball will revolve in a vertical circle, so that the time of revolution will be to the time of descent (of another

other body) in the quadrantal arc, whose lowest point is C and radius  $r$ , as  $2^{\frac{1}{2}}$  to 1; which time of descent is specified in the preceding Remark.

*Remark 5.* Taking  $v$  and  $\dot{v}$  each = 0, in the value of  $f$  above written, we get

$$\frac{2ma}{r} = \frac{ab^2\sqrt{r^2-a^2}}{r^3}, \text{ or } \frac{2mr^2}{\sqrt{r^2-a^2}} = b^2 = \frac{d^2r^2}{a^2},$$

$d$  denoting the initial velocity of the ball.

Whence we have  $d = a \times \frac{\sqrt{2m}}{\sqrt{r^2-a^2}}$ : and therefore it follows, that the ball, having a velocity given it equal to that particular value of  $d$ , in a horizontal direction at right angles to the string, will revolve uniformly in a circle in a horizontal plane; as has been demonstrated in a different manner by Mr. HUYGENS and others.

*Remark 6.* When the initial velocity of the ball is nearly  $= a \times \frac{\sqrt{2m}}{r-b}$  its track (projected on a horizontal plane) will be nearly a circle, and the time of its revolution therein nearly  $= \frac{4q\sqrt{r-b}}{2m^{\frac{1}{2}}}$ ,  $q$  being as in Remark 2.

Moreover the *limit* of the time of moving from apsis to apsis, upon taking the initial velocity nearer and nearer to  $a \times \frac{\sqrt{2m}}{r-b}$ , appears, by such comparison as is mentioned

in Remark 1, to be  $= \frac{2qr\sqrt{r-b}}{2m^{\frac{1}{2}} \times \sqrt{4r^2-3a^2}}$ . Consequently, when the track differs but little from a circle, the angle between apsis and apsis will be nearly  $= \frac{r \cdot 180^\circ}{\sqrt{4r^2-3a^2}}$ ; as  
found

found by Mr. EULER, (in his *Mechanics*,) by a very different method.

3. If only the force  $g$  act on the body;  $f$  being = 0,  $\frac{u^2}{r^2}$  will be =  $\frac{v\dot{v}}{y\dot{y}}$ : and, by means of that equation, the value of the force  $g$  ( $= \frac{v\sqrt{r^2-y^2}}{r^2} \times \frac{\dot{u}y + 2v\dot{y}}{y}$ ) may be readily computed, when the equation of the curve described by the body is given.

*Example.* If the proposed curve be the loxodromic specified in Example 2. of the preceding article;  $u$  being =  $\frac{rsv}{\sqrt{1-s^2}y}$ ,  $\frac{v\dot{v}}{y\dot{y}}$  ( $= \frac{u^2}{r^2}$ ) will be =  $\frac{s^2v^2}{1-s^2y^2}$ ,  $\frac{\dot{v}}{v} = \frac{s^2}{1-s^2} \cdot \frac{\dot{y}}{y}$ ,

and  $v = c \cdot \frac{y^{\frac{1}{1-s^2}}}{a}$ . Consequently  $u$  will be =  $b \cdot \frac{y^{\frac{2s^2-1}{1-s^2}}}{a}$ ,

and  $g = \frac{ab^2\sqrt{r^2-y^2}}{r^2s\sqrt{1-s^2}} \cdot \frac{y^{\frac{3s^2-1}{1-s^2}}}{a}$ ,  $crs$  being =  $ab\sqrt{1-s^2}$ .

If  $s$  be =  $\frac{1}{\sqrt{3}}$ ,  $g$  must be =  $\frac{3ab^2\sqrt{r^2-y^2}}{2^{\frac{1}{2}}r^2}$ .

4. In general, from proper data, we may, by means of our theorems investigated in art. 1. find not only the requisite force or forces, but also every thing else that may be required relative to the track of the body and its motion therein.

*Example 1.* Suppose the body to describe with an invariable velocity, the loxodromic specified in Ex. 2. art. 2. and referred to in the last example.

Then,

Then,  $d$  being the given velocity of the body,  $v$  will be  $= d\sqrt{1-s^2}$ , and  $\frac{uy}{r} = ds$ . Consequently, by substitution, we have  $f = \frac{d^2 s^2 \sqrt{r^2 - y^2}}{ry}$ , and  $g = \frac{d^2 s \sqrt{1-s^2} \sqrt{r^2 - y^2}}{ry}$ . The single force requisite to cause the body so to move in the proposed curve is  $\sqrt{f^2 + g^2} = \frac{d^2 s \sqrt{r^2 - y^2}}{ry}$ ; the variable direction in which it must act being always at right angles to the track of the body.

*Example 2.* If the body describe the loxodromic referred to in the preceding example, so that the velocity  $u$  shall always be equal to the invariable quantity  $b$ ;  $s$  will be to  $\frac{by}{r}$  as  $\sqrt{1-s^2}$  to  $v = \frac{b\sqrt{1-s^2}y}{rs}$ . Consequently, by substitution, we have in this case,

$$f = \frac{b^2 \cdot 2s^2 - 1 \cdot y \sqrt{r^2 - y^2}}{r^2 s^2}, \text{ and } g = \frac{2b^2 \sqrt{1-s^2} y \sqrt{r^2 - y^2}}{r^2 s}$$

The requisite single force is  $\sqrt{f^2 + g^2} = \frac{b^2 y \sqrt{r^2 - y^2}}{r^2 s^2}$ ; and its variable direction such, that the angle it must make with the great circle passing through the place of the body and the pole C shall always be bisected by the track of the body.

*Example 3.* Let  $u$  and the absolute velocity of the body be supposed invariable.

Then, those velocities being denoted by  $b$  and  $d$  respectively,  $v^2 + \frac{b^2 y^2}{r^2}$  will be  $= d^2$ ,  $v = \sqrt{d^2 - \frac{b^2 y^2}{r^2}}$ , and  $t = \frac{\pi}{b}$

=

$$= \frac{rj}{v\sqrt{r^2-y^2}} = \frac{r^2j}{\sqrt{r^2-y^2} \times \sqrt{d^2r^2-b^2y^2}}$$
 And it appears, by our Appendix, that  $t$  and  $z$  will always be assigned by means of elliptic arcs: except when  $d$  is  $= b$ ; and then,  $\dot{t}$  ( $= \frac{\dot{z}}{b} = \frac{rj}{v\sqrt{r^2-y^2}}$ ) being  $= \frac{r^2j}{b.r^2-y^2}$ ,  $t$  ( $= \frac{z}{b}$ ) will be  $= \frac{r}{2b} \times \text{Log.} \frac{r-a}{r+a} \cdot \frac{r+y}{r-y}$ ; and it follows, that, in this particular case, the body (describing a spiral) will continually approach nearer and nearer to the great circle of which C is a pole, yet never can arrive thereat.

The single force requisite to cause the body to move according to our supposition is  $\sqrt{f^2+g^2} = \frac{2bd\sqrt{r^2-y^2}}{r^2}$ ; and the variable direction in which it must act will always be at right angles to the track of the body.

*Example 4.* Supposing the velocities  $u$  and  $v$  to be invariable, and equal to  $b$  and  $c$  respectively: it appears by our theorems, that

$$f \text{ must be } = \frac{b^2y\sqrt{r^2-y^2}}{r^3}, \text{ and } g = \frac{2bc\sqrt{r^2-y^2}}{r^2};$$

and the requisite single force  $\sqrt{f^2+g^2}$ , compounded of those two forces, must be  $= \frac{b\sqrt{r^2-y^2} \times \sqrt{4c^2r^2+b^2y^2}}{r^3}$ ; the direction in which it must act being inclined to the plane of the great circle BC (passing through the place of the body and the pole C) in an angle whose sine shall be to radius as  $2c$  to  $\frac{\sqrt{4c^2r^2+b^2y^2}}{r}$ , and in a plane at right angles to the plane of the said great circle.

Z

Remark.

*Remark.* Mr. SIMPSON, in Lem. 2. pag. 3. of his *Miscell. Tracts*, computes (by a different method) the value of the force  $g$  (making it as above) upon the supposition that, if  $u$  be invariable,  $v$  will likewise be so; without taking any notice that  $v$  will not be invariable when  $u$  is so, unless the body be acted on by a force  $f = \frac{b^2 y \sqrt{r^2 - y^2}}{r^3}$  as well as by the force  $g = \frac{2bc \sqrt{r^2 - y^2}}{r^2}$ .\*

Indeed he has considered the velocity  $b$  as very small: and then the force  $f$  will be very small, but not absolutely  $= 0$ , nor yet indefinitely small; for  $b$  (though small) being finite, the value of  $f$  will be also finite.

Mr. DE LA LANDE, in his *Astronomy* (art. 3547.), proposing to explain Mr. SIMPSON'S computation, has observed, that, only the force  $g$  acting on the body,  $v$  will not be invariable when  $u$  is so. But, without computing the requisite force  $f$ , or the exact value of the force  $g$  when  $f$  is  $= 0$ , he neglects a part of the force  $g$ , and entirely neglects the force  $f$ , as being what he calls *infiniments petits du troisieme ordre*; whereas they are not generally *infiniments petits* of any order whatever, being assignable finite quantities which may in some cases be considerable, and therefore should not be neglected without first computing their values, and shewing that, in the case in question, they are really inconsiderable.

*Example 5.* If,  $u$  being invariable, the force  $f$  be  $= 0$ ;  $u^2 v \dot{v}$  will be  $= b^2 y \dot{y}$ , and  $v = \frac{\sqrt{c^2 r^2 + b^2 y^2}}{r}$ ,  $c$  being the

\* If only this force  $g$  act on the body, neither  $v$  nor  $u$  will be invariable;  $u$  being, in that case,  $= \frac{2bcy}{uy + 2ux}$ , and  $u^2 = \frac{r^2 v \dot{v}}{y \dot{y}}$ , as appears by the theorems in art. 1.

value

value of  $v$  when  $y$  is  $= 0$ . Therefore it follows, that  $g$  will, in that case, be  $= \frac{2b\sqrt{r^2 - y^2} \times \sqrt{c^2 r^2 + b^2 y^2}}{r^2}$ ,

$$\text{and } t = \frac{z}{b} = \frac{rj}{v\sqrt{r^2 - y^2}} = \frac{r^2 j}{\sqrt{r^2 - y^2} \times \sqrt{c^2 r^2 + b^2 y^2}}$$

and, by our Appendix, it appears that  $t$  and  $z$  will always be assigned by means of elliptic arcs.

*Remark 1.* If  $c$  be  $= b$ , the time in which the revolving great circle (CB) will make one revolution, about the diameter drawn from the pole C, will be to the time in which the body will make one revolution in that circle, as  $2q$  to  $e + \sqrt{e^2 - 2q}$ ; that is, as 1.57079632 to 1.3110287771;  $e$  and  $q$  being as specified in Rem. 2. Ex. 3. Art. 2.

*Remark 2.* The consequence is obvious when a number of bodies, kept from flying from the spherical surface by any force whatever tending to its center, follow one another in the revolving circle CB from the pole C, each having the same velocity ( $c$ ) when at that pole. If only the force  $g = \frac{2b\sqrt{r^2 - y^2} \times \sqrt{c^2 r^2 + b^2 y^2}}{r^2}$  act on each, they will always be found (forming a kind of ring) in one and the same great circle CB, supposed to revolve uniformly about the diameter drawn from C; but they will not revolve uniformly in such circle, unless each be continually acted on by the two forces  $f$  and  $g$ , whose values are computed in the 4th Example.

*Example 6.* If a small ball move within a slender circular tube, or a small ring upon a slender circular rod,  
Z 2
whilst



whilst such tube or rod revolves uniformly about a diameter thereof perpendicular to the horizon;  $u$  being invariable, and  $f = \frac{2my}{r}$ ,  $r^2 v \dot{v}$  will be  $= b^2 y \dot{y} - \frac{2mr^2 y \dot{y}}{\sqrt{r^2 - y^2}}$ ,

$$\text{and } v = \frac{\sqrt{c^2 r^2 - 4mr^2 + b^2 y^2 + 4mr^2 \sqrt{r^2 - y^2}}}{r};$$

$c$  being the value of  $v$  when  $y$  is  $= 0$ .

Therefore

$$\dot{t} = \frac{\dot{x}}{b} = \frac{r \dot{y}}{v \sqrt{r^2 - y^2}} \text{ will be } = \frac{r^2 \sqrt{r^2 - y^2}^{-\frac{1}{2}} \dot{y}}{\sqrt{c^2 r^2 - 4mr^2 + b^2 y^2 + 4mr^2 \sqrt{r^2 - y^2}}}$$

which, by substituting  $x$  for  $\sqrt{r^2 - y^2}$ , will be adapted to our Tables, and we have another instance of their use in assigning the fluent by means of elliptic or circular arcs, or logarithms, as the case may require.

This case coincides with the preceding, if, abstracting from gravity,  $m$  be considered as  $= 0$ .

---



---

## M E M O I R IX.

*Of the Motion of a Body in any variable Plane.*

1. **L**ET the plane BCD, in which the body (B) is always to be found, be supposed to revolve about an immoveable axis GD, with the angular velocity  $u$  measured on the circular arc  $z$ , described about the center C by a point in that plane at the distance  $r$  from the said axis: let  $v$  denote the velocity, and  $y$  the distance of the body from that axis;  $v$  its velocity, and  $x - k$  the distance it has moved in the direction  $B\beta$  parallel to the same axis: and let  $B \times f$  denote a motive force continually urging the body towards the said axis CD;  $B \times g$  another such force urging it in a direction at right angles to the said plane BCD; and  $B \times h$  a third force urging it in the said direction  $B\beta$ . Fig. 51.

Then, by the method pursued in Memoir VII.

$$f \text{ is found} = \frac{u^2 y j - r^2 v \dot{v}}{r^2 j}, \quad g = \frac{v}{r} \times \frac{u j + 2 u \dot{j}}{j},$$

$$h = \frac{v \dot{v}}{j} = \frac{v \dot{v}}{x};$$

and  $\dot{t}$ , the fluxion of the time, will be  $= \frac{\dot{x}}{v} = \frac{j}{v} = \frac{\dot{x}}{u}$ .

From these theorems others may be readily derived, by the resolution or composition of forces or motion, to suit

suit particular propositions : and the theorems so derived, with others that the particular circumstances of the proposition to be considered may suggest, will enable the intelligent analyst to proceed in determining the path of the body and its motion therein in any possible case whatever; as every force that can act on the body in any direction different from the particular directions in which we have supposed the forces  $B \times f$ ,  $B \times g$ ,  $B \times h$  to act may be resolved into two or three others acting in two or three of those particular directions.

2. If the force  $g$  be  $= 0$ ;  $\dot{u}y + 2uy$  being  $= 0$ , we have  $xy^2 = a^2b$ ,  $b$  being the value of  $x$  when  $y$  is  $= a$ : and  $f$  will be  $= \frac{a^2b^2}{r^2y^3} - \frac{v\dot{v}}{j}$ .

Let  $B$ , the place of the body, be projected in  $B''$  by a perpendicular on a plane  $CB''$  at right angles to the axis  $CD$ ; and let  $Z$  be the fluxion of the line in which the point  $B''$  shall be found,  $U$  the velocity of that point in that line, and  $p$  the perpendicular from  $C$  on the tangent to the same line at the point  $B''$  corresponding to the place of the body: then will  $U$  be  $= \frac{\sqrt{a^2b^2 + r^2v^2y^2}}{ry}$ ,

$$f = \frac{a^2b^2}{r^2y^3} - \frac{v\dot{v}}{j} = \frac{a^2b^2}{r^2y^3} - \frac{a^2b^2}{2j} \times \text{the fluxion of } \frac{j^2}{y^2z^2} = \frac{a^2b^2\dot{p}}{r^2p^3j},$$

$$v \text{ being } = \frac{Uj}{Z} = \frac{uj}{x} = \frac{a^2bj}{y^2z} = \frac{a^2b\sqrt{y^2 - p^2}}{rpy}.$$

Fig. 52. 3. Supposing the body to move near the earth's surface in the concave superficies of a solid of revolution, generated by the line  $DB$  revolving about the axis  $DC$ ; let the abscissa  $Db$ , and the correspondent ordinate  $bB$  at right

right angles thereto, be denoted by  $x$  and  $y$  respectively; and let  $\xi$  denote the fluxion of the said line DB.

Then, DbC being perpendicular to the horizon,  $g$  will be  $= 0$ ; and  $\frac{bx}{\xi}$  will denote the force upwards, in the direction of the tangent to the line DB, arising from the action of the force  $B \times h$ ;  $\frac{fy}{\xi}$  the force downwards, in the direction of the said tangent, arising from the action of the force  $B \times f$ ; and  $\frac{fy - bx}{\xi}$  ( $= \frac{a^2 b^2 j}{r^2 y^3 \xi} - \frac{v \dot{v}}{\xi} - \frac{v \dot{v}}{\xi} = \frac{a^2 b^2 j}{r^2 y^3 \xi} - \frac{V \dot{V}}{\xi} = -\frac{W \dot{W}}{\xi}$ ) the actual force downwards in the said direction;  $V$  denoting the velocity of the body upwards in that same direction, and  $W$  its absolute velocity in the path it describes.

Which last mentioned force will be  $= \frac{2m\dot{x}}{\xi}$ , the force arising from the gravity of the body in the direction of the said tangent,  $2m$  denoting ( $32\frac{1}{2}$  feet) the accelerating force of gravity directly downwards. Therefore

$$\frac{a^2 b^2 j}{r^2 y^3} - V \dot{V} \text{ will be } = 2m\dot{x};$$

and consequently  $V = \sqrt{d^2 + \frac{a^2 b^2}{r^2} + 4mk - 4mx - \frac{a^2 b^2}{r^2 y^2}}$

$d$  being the value of  $V$  when  $x$  is  $= k$  and  $y = a$ .

Whence we have

$$\dot{x} = \frac{\xi}{V} = \frac{\xi}{\sqrt{d^2 + \frac{a^2 b^2}{r^2} + 4mk - 4mx - \frac{a^2 b^2}{r^2 y^2}}}$$

and

$$\text{and } \dot{z} = \frac{a^2 b \dot{t}}{y^2} = \frac{a^2 b y^{-2} \dot{t}}{\sqrt{d^2 + \frac{a^2 b^2}{r^2} + 4mk - 4mx - \frac{a^4 b^2}{r^2 y^2}}}$$

and hence, the relation of  $x$  and  $y$  being given, the values of  $t$  and  $z$  may be found; by which means the place of the body at any time will be determined.

$\dot{V}$  being = 0 when,  $y$  being =  $a$ ,  $\frac{a^4 b^2 j}{r^2 y^3}$  is =  $2mx$ ; it follows, that, if the initial velocity of the body be =  $\frac{2amx}{j}$  and in a horizontal direction, the body will revolve uniformly in a circle parallel to the horizon.

Fig. 53. *Example 1.* Let the superficies in which the body moves be that of an inverted cone, the radius of whose base is to its perpendicular height as  $r$  to  $q$ .

Then,  $x$  being supposed to begin at the vertex D,  $\frac{rx}{q}$  will be =  $y$ ,  $\frac{aq}{r} = k$ ,

$$V = \sqrt{d^2 + \frac{a^2 b^2}{r^2} + \frac{4amq}{r} - 4mx - \frac{a^4 b^2 q^2}{r^2 x^2}}$$

$$\dot{t} = \frac{r^2 s x \dot{x}}{q \sqrt{d^2 r^4 + a^2 b^2 r^2 + 4amqr^3 x^2 - 4mr^4 x^3 - a^4 b^2 q^2}}$$

$$\text{and } \dot{z} = \frac{a^2 b q t x^{-1} \dot{x}}{\sqrt{d^2 r^4 + a^2 b^2 r^2 + 4amqr^3 x^2 - 4mr^4 x^3 - a^4 b^2 q^2}}$$

$s$  being put for  $\sqrt{q^2 + r^2}$ :

and, by the help of our Tables, the value of  $t$  may be assigned by means of elliptic arcs: the value of  $z$  also may sometimes be so assigned; for instance, when,  $d$  being = 0 and  $r = a$ , the initial velocity of the body is =  $2\sqrt{2^{\frac{1}{2}} - 1} \times \sqrt{mq}$ .

*Remark*

*Remark 1.* If  $d$  be = 0,

$$V \text{ will be } = x^{-1} \sqrt{4m \times x - \frac{aq}{r} \times R' - x \times R'' + x};$$

$$R' \text{ being } = \frac{\sqrt{a^4 b^4 + 16 a^3 b^2 m q r + a^2 b^2}}{8 m r^2},$$

$$R'' \dots = \frac{\sqrt{a^4 b^4 + 16 a^3 b^2 m q r - a^2 b^2}}{8 m r^2}.$$

Hence it appears, that the ascent and descent of the body will be limited by two horizontal circles; whose heights above D, in this case, will be  $\frac{aq}{r}$  and  $R'$ : and the places of the apfides (in those circles) will vary every revolution.

*Remark 2.* If the initial velocity of the body be =  $\left(\frac{2amq}{r}\right)^{\frac{1}{2}}$ , and in a direction parallel to the horizon, the two limiting circles (mentioned in the preceding Remark) will coincide, and the body will revolve in the circle of coincidence, always keeping the same distance from the vertex of the cone; the value of  $\dot{V}$  being then = 0, as well as  $V = 0$ .

When the initial velocity is nearly =  $\left(\frac{2amq}{r}\right)^{\frac{1}{2}}$ , and in a direction parallel to the horizon, the path of the body will be nearly a circle: and the limit of the angle described between apfis and apfis may be found by the method commonly pursued in such cases; or by means of the limit of the time of moving from apfis to apfis, as in the like case in the preceding Memoir.

*Remark 3.* If, abstracting from gravity,  $m$  be considered as = 0,

A a

$t$  will

$$\dot{s} \text{ will be } = \frac{r^2 s \dot{x}}{q \sqrt{d^2 r^2 + a^2 b^2 r^2 \times x^2 - a^4 b^2 q^2}}$$

$$\dot{x} \dots \dots = \frac{a^2 b q s x^{-1} \dot{x}}{\sqrt{d^2 r^2 + a^2 b^2 r^2 \times x^2 - a^4 b^2 q^2}}$$

$$t = \frac{s}{d^2 r^2 + a^2 b^2} \times \sqrt{\frac{d^2 r^2 + a^2 b^2 r^2}{q^2} \times x^2 - a^4 b^2} - a d r,$$

$$\text{and } z = \begin{cases} \text{Circ. arc, rad. } s, \text{ sec. } \frac{r s \sqrt{d^2 r^2 + a^2 b^2}}{a^2 b q} \times x. \\ -\text{Circ. arc, rad. } s, \text{ sec. } \frac{s \sqrt{d^2 r^2 + a^2 b^2}}{a b} \end{cases}$$

or, supposing  $d = 0$ ,

$$t \text{ will be } = \frac{s}{a b} \times \sqrt{\frac{r^2 x^2}{q} - a^2}$$

$$\text{and } z = \text{circ. arc, rad. } s, \text{ sec. } \frac{r s x}{a q}$$

**Fig. 54.** *Example 2.* If the superficies wherein the body moves be that of a parabolic conoid, generated by the revolution of the parabola whose equation is  $q x = y^2$  about its own axis;  $x$  being supposed to begin at D,  $t$  will be  $= \frac{a^2}{q}$ ,

$$V = \sqrt{d^2 + \frac{a^2 b^2}{r^2} + \frac{4 a^2 m}{q} - 4 m x - \frac{a^4 b^2}{q r^2 x}}$$

$$\dot{s} = \frac{\frac{1}{2} q^{\frac{1}{2}} r \dot{x} \sqrt{q+x}}{\sqrt{d^2 q r^2 + a^2 b^2 q + 4 a^2 m r^2 \times x - 4 m q r^2 x^2 - a^4 b^2}}$$

$$\text{and } \dot{x} = \frac{\frac{1}{2} a^2 b q^{-\frac{1}{2}} r x^{-1} \dot{x} \sqrt{q+x}}{\sqrt{d^2 q r^2 + a^2 b^2 q + 4 a^2 m r^2 \times x - 4 m q r^2 x^2 - a^4 b^2}}$$

and in this Example, as well as in the preceding one, the value of  $t$  may, by having recourse to our Tables, be assigned.

assigned by means of elliptic arcs : the value of  $z$  also may *sometimes* be so assigned ; for instance, when,  $d$  being = 0,

the initial velocity of the body is =  $2a \times \frac{mq}{q^2 \pm a^2}^{\frac{1}{2}}$ .

Moreover it is obvious, that remarks respecting the motion of the body, similar to those in the preceding example, may be here made : except that, if  $m$  were here considered as = 0, the values of  $t$  and  $z$  would not be assigned by the like means as in that example ; but the value of  $t$  would be assigned by means of a logarithm, and the value of  $z$  by means of a circular arc and a logarithm.

4. If the body be acted on by the force  $F$ , urging it towards the center  $C$  in the axis  $CD$  ; and by the force  $G$ , urging it towards the plane  $CAB''$  passing through  $C$  at right angles to the said axis, and always in a direction perpendicular to that plane ; Fig. 55.

$$g \text{ being } = 0, f = \frac{a^2 b^2}{r^2 y^2} - \frac{v \dot{v}}{y} \text{ will be } = \frac{y}{\sqrt{x^2 + y^2}} \times F,$$

$$\text{and } h = \frac{v \dot{v}}{x} = -G - \frac{x}{\sqrt{x^2 + y^2}} \times F.$$

Hence, when the values of the forces  $F$  and  $G$  are given in terms of  $x$  and  $y$ , the path of the body and its motion therein may be determined.

*Example.* Let the forces  $F$  and  $G$  be supposed to be as the distances of the body from the center  $C$  and from the plane  $CAB''$  respectively ; that is, let  $F$  be supposed =  $m\sqrt{x^2 + y^2}$  and  $G = nx$ ,  $m$  and  $n$  being invariable quantities.

A a 2

Then



Then will  $\frac{a^2 b^2 \dot{y}}{r^2 y^3} - v \dot{v}$  be  $= m y \dot{y}$ , and  $v \dot{v} = -\overline{m+n.x\dot{x}}$ :

$$\text{whence we have } v = \sqrt{m a^2 + \frac{a^2 b^2}{r^2} - m y^2 - \frac{a^2 b^2}{r^2 y^2}}$$

$$v = \sqrt{e^2 + m + n.k^2 - m + n.x^2}$$

$$\dot{t} = \frac{\dot{x}}{\sqrt{e^2 + m + n \times k^2 - x^2}} = \frac{r y \dot{y}}{\sqrt{m a^2 r^2 + a^2 b^2 \times y^2 - m r^2 y^4 - a^2 b^2}}$$

$$\text{and } \dot{z} = \frac{a^2 b r y^{-1} \dot{y}}{\sqrt{m a^2 r^2 + a^2 b^2 \times y^2 - m r^2 y^4 - a^2 b^2}}$$

$v$  being supposed  $= 0$  and  $v = e$ , when  $x$  is  $= k$  and  $y = a$ .

Therefore  $p$  will be  $= \frac{a^2 b}{\sqrt{m a^2 r^2 + a^2 b^2 - m r^2 y^2}}$ : and it follows, that the curve of projection will be an ellipsis whose center will be C, and semi-axes  $\frac{ab}{m^{\frac{1}{2}} r}$  and  $a$ ; which curve becomes a circle whose radius is  $a$ , when,  $v$  being  $= 0$ , the initial velocity of the body is  $= \sqrt{m a^2 + e^2}$ .

From those fluxional equations we find, by taking the fluents,

$$\begin{aligned} t &= \frac{1}{m+n} \times \text{circ. arc, rad. } a, \text{ cof. } a \times \frac{m+n.kx + e \sqrt{e^2 + m + n.k^2 - x^2}}{e^2 + m + n.k^2} \\ &= \frac{1}{m^{\frac{1}{2}} a} \times \text{circ. arc, rad. } a, \text{ cofine } \frac{a^2 b^2 - m r^2 y^2}{b^2 - m r^2}^{\frac{1}{2}}, \end{aligned}$$

$$\text{and } z = r \times \text{circ. arc, rad. } 1, \text{ cofine } \frac{1}{y} \times \frac{a^2 b^2 - m r^2 y^2}{b^2 - m r^2}^{\frac{1}{2}}.$$

It is evident that, the ordinate  $bB''$  being at right angles to the shorter semi-axis  $CbA$  (denoted by  $a$ ), the abscissa

$Cb$

Cb will be  $= \frac{a^2 b^2 - m r^2 y^2}{b^2 - m r^2} \Big|^{1/2}$ ; and, if that abscissa be represented by  $w$ ,

$$t \text{ will be } = \frac{r}{m+n} \times \text{circ. arc, rad. } a, \text{ cosine } w :$$

therefore, from comparing this value of  $t$  with that above found in terms of  $x$ , it follows, that

$$\frac{m+n.kx + e \sqrt{e^2 + m+n.k^2 - x^2}}{e^2 + m+n.k^2} \text{ will be}$$

$$= \frac{r}{a} \times \text{cosine of } \frac{m+n}{m} \Big|^{1/2} \text{ times the arc, rad. } a, \text{ cosine } w ;$$

which, by art. 2. Mem. V. is

$$= \frac{w + \sqrt{w^2 - a^2}^q + w - \sqrt{w^2 - a^2}^q}{2a^q}, \text{ } q \text{ being put for } \frac{m+n}{m} \Big|^{1/2}.$$

Consequently the value of  $x$  will from hence be obtained in terms of the abscissa  $w$ .

The points in which the path of the body intersects the plane CAB" being distinguished by the name of the *nodes*; let the body, when  $y$  is  $= a$ , be supposed to be *below* the said plane and approaching towards an ascending node. Then,  $x$  being  $= 0$  at the time the body comes to the node,  $w$  will at the same time be the cosine of  $\frac{A}{q}$ , A de-

noting the arc whose radius is  $a$  and cosine  $\frac{ae}{\sqrt{e^2 + m+n.k^2}}$ ;

and, when the body comes to the next descending node,  $w$  will be the cosine of  $\frac{A''}{q}$ , A'' denoting the arc (of the

same

same circle) whose cosine is  $\frac{-ae}{\sqrt{e^2 + m + nk^2}}$ , the quantity

expressed by  $\sqrt{e^2 + m + nk^2} - x^2$  becoming negative after being equal to 0, as it will be when the body shall be at the greatest distance from the said plane CAB": moreover, when the body comes again to an ascending node,  $w$  will be the cosine of  $\frac{A'''}{q}$ ,  $A'''$  denoting the arc whose radius

is  $a$  and cosine  $\frac{ae}{\sqrt{e^2 + m + nk^2}}$ , which is the same as the

cosine of the arc  $A$ . Therefore, denoting the quadrantal arc of the circle whose radius is  $a$  by  $Q$ , it evidently follows, that the successive places of the nodes (varying in antecedentia) will be determined by taking (in the ellipsis of projection) the abscissa  $w$  equal to the cosine of  $\frac{A}{q}$ ,

$\frac{A+2Q}{q}$ ,  $\frac{A+4Q}{q}$ ,  $\frac{A+6Q}{q}$ , &c. successively.

The position of the ray CB" when the body shall beat the greatest distance from the plane CAB" is found by tak-

ing the value of  $v$  equal to 0, and considering  $\frac{\sqrt{e^2 + m + nk^2}}{\sqrt{m+n}}$

the value of  $x$  at such time, as alternately positive and negative. By which means it appears, that the successive positions of the said ray CB", when the body shall be at the greatest distance from the said plane, will be determined by taking (in the said ellipsis of projection) the ab-

scissa  $w$  equal to the cosine of  $\frac{A+Q}{q}$ ,  $\frac{A+3Q}{q}$ ,  $\frac{A+5Q}{q}$ , &c.

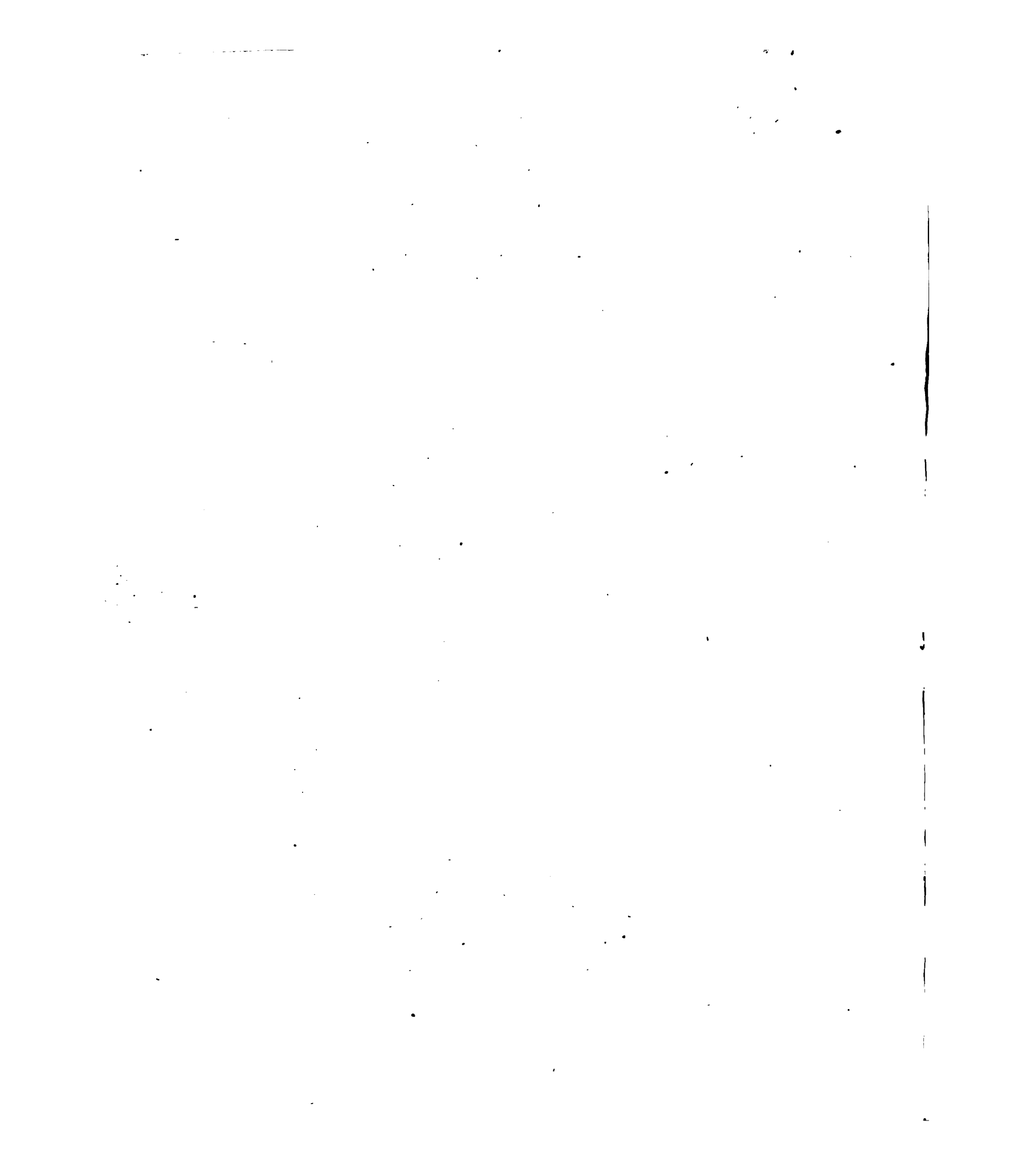
successively;  $\frac{(m+n)^{\frac{1}{2}} \times ak}{\sqrt{e^2 + m + nk^2}}$  (=  $\mp$  the sine of the arc  $A$ )

being

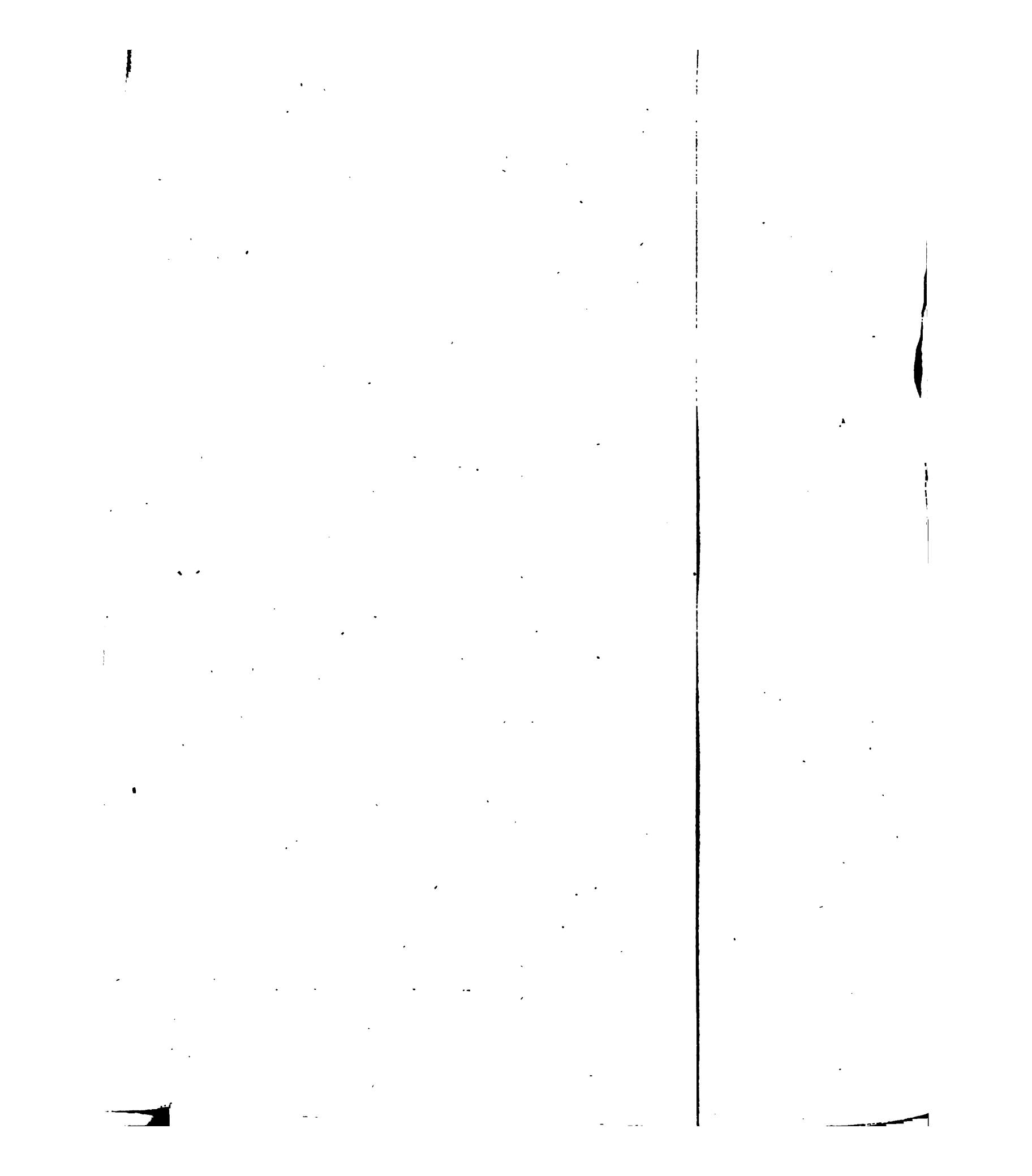
being the cosine of  $A + Q$ , considering  $k$  as negative, agreeable to the supposition above.

It is manifest that, if  $q (= \frac{m+n}{m})^{\frac{1}{2}}$  be a rational number  $= \frac{N}{D}$  in its least terms, the nodes, after  $\frac{1}{2}D$  or  $D$  revolutions of the body, (according as  $D$  is even or odd,) will successively fall in the same points as before.

The theorems in the preceding Memoir, it is obvious, might be deduced from those in art. 1. of this Memoir: and many other instances might be given of the use of the principal theorems in these last three Memoirs; some other instances of their use may probably appear hereafter in the course of this Work, when we come to consider some particular subjects in distinct Memoirs.

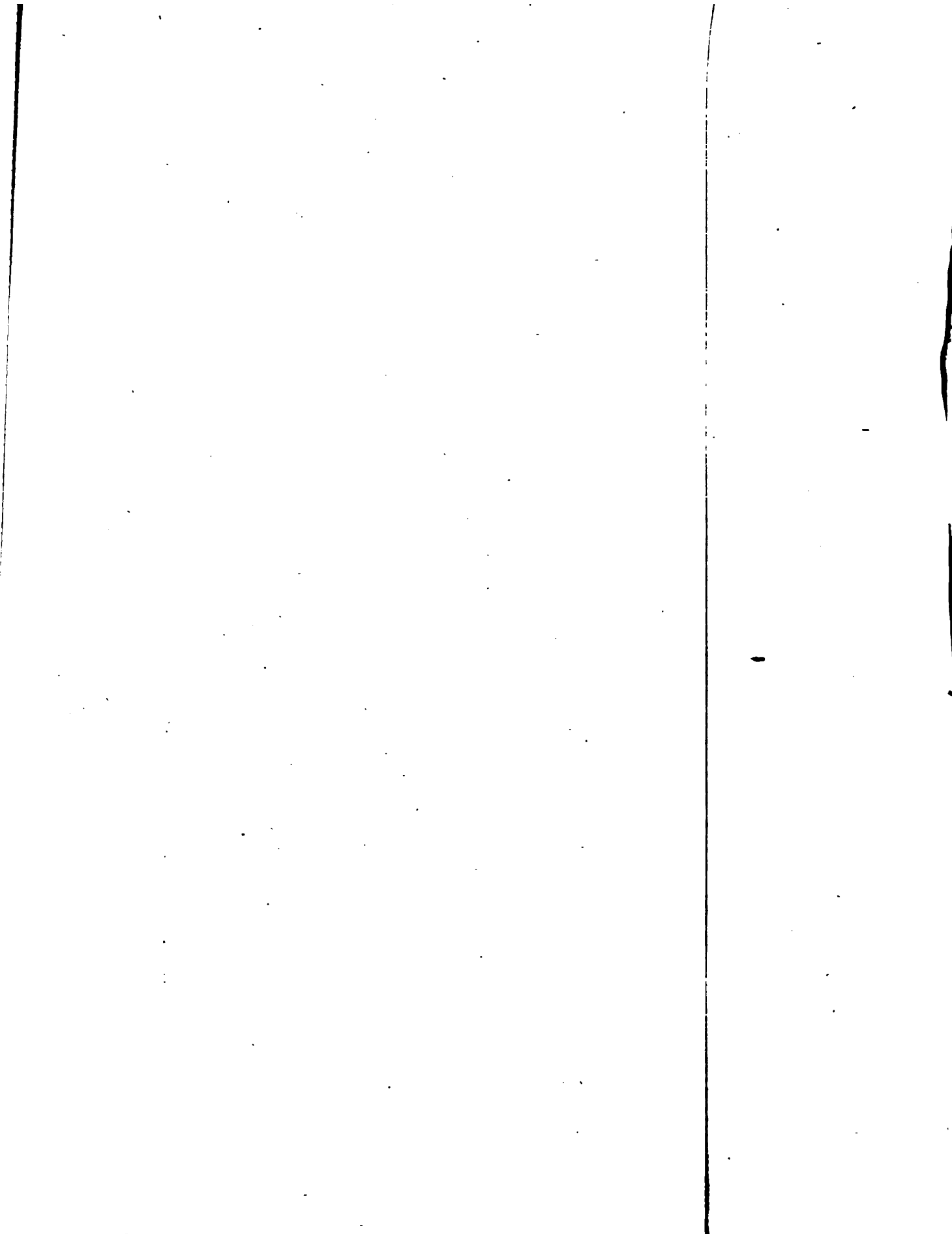




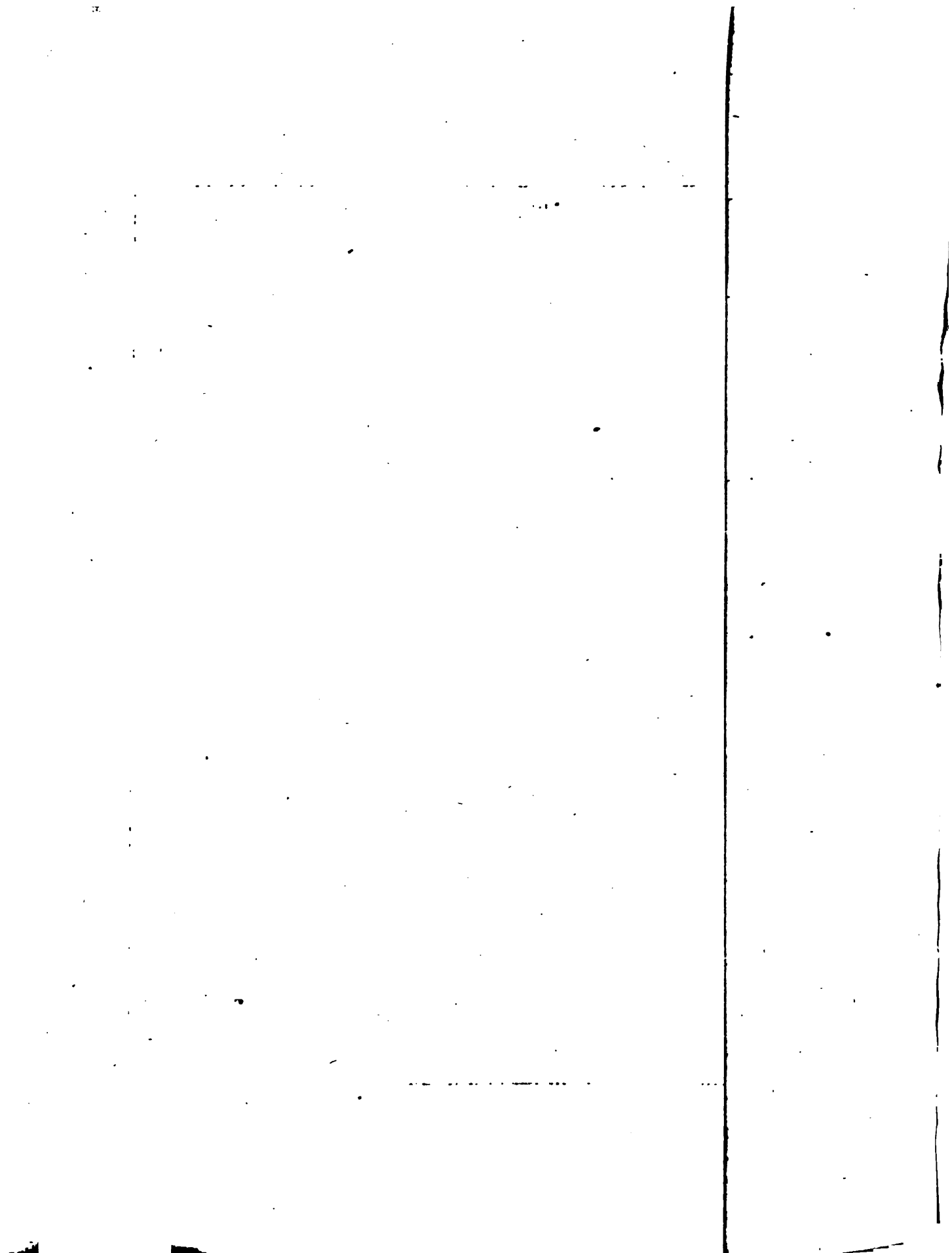


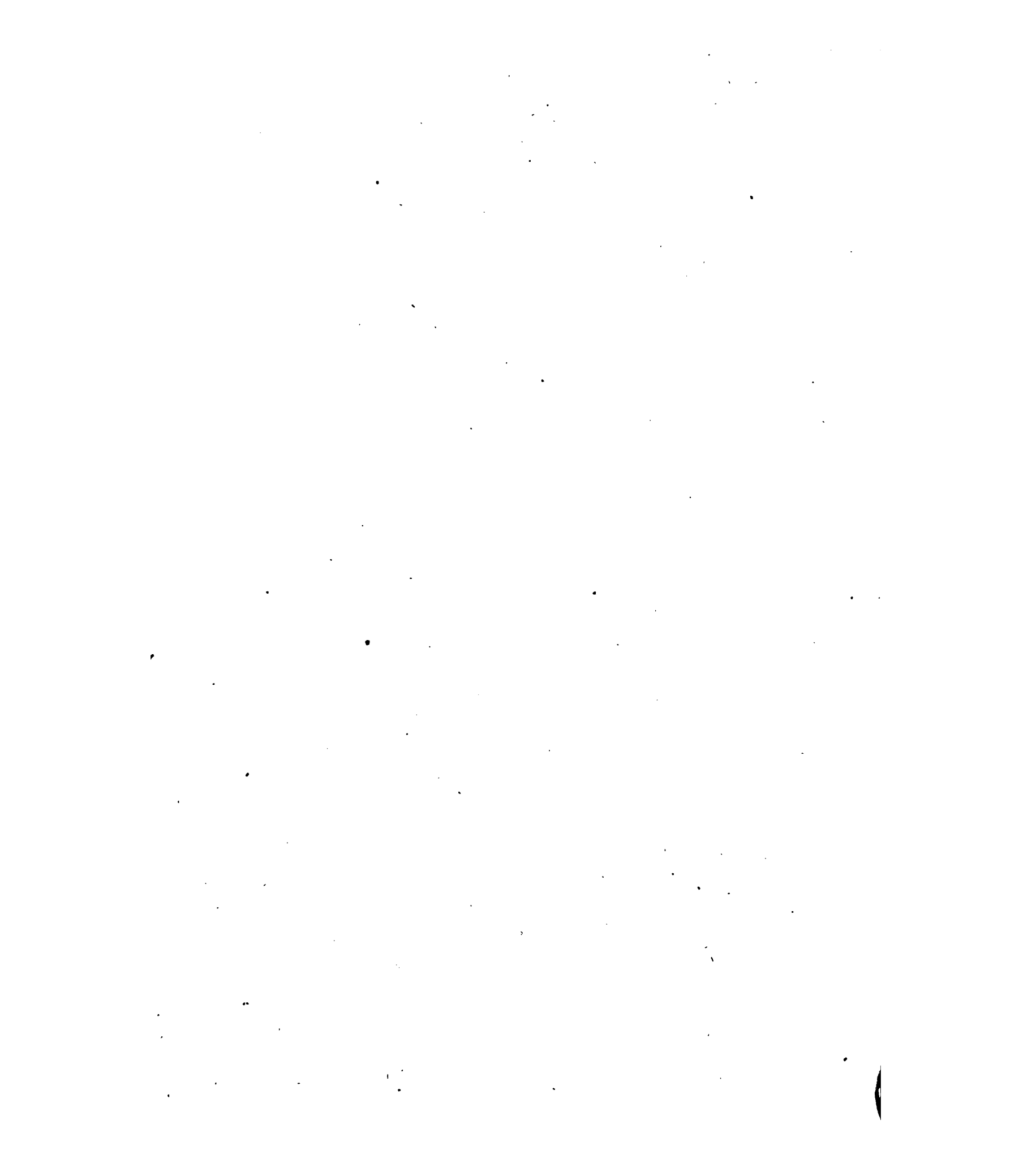


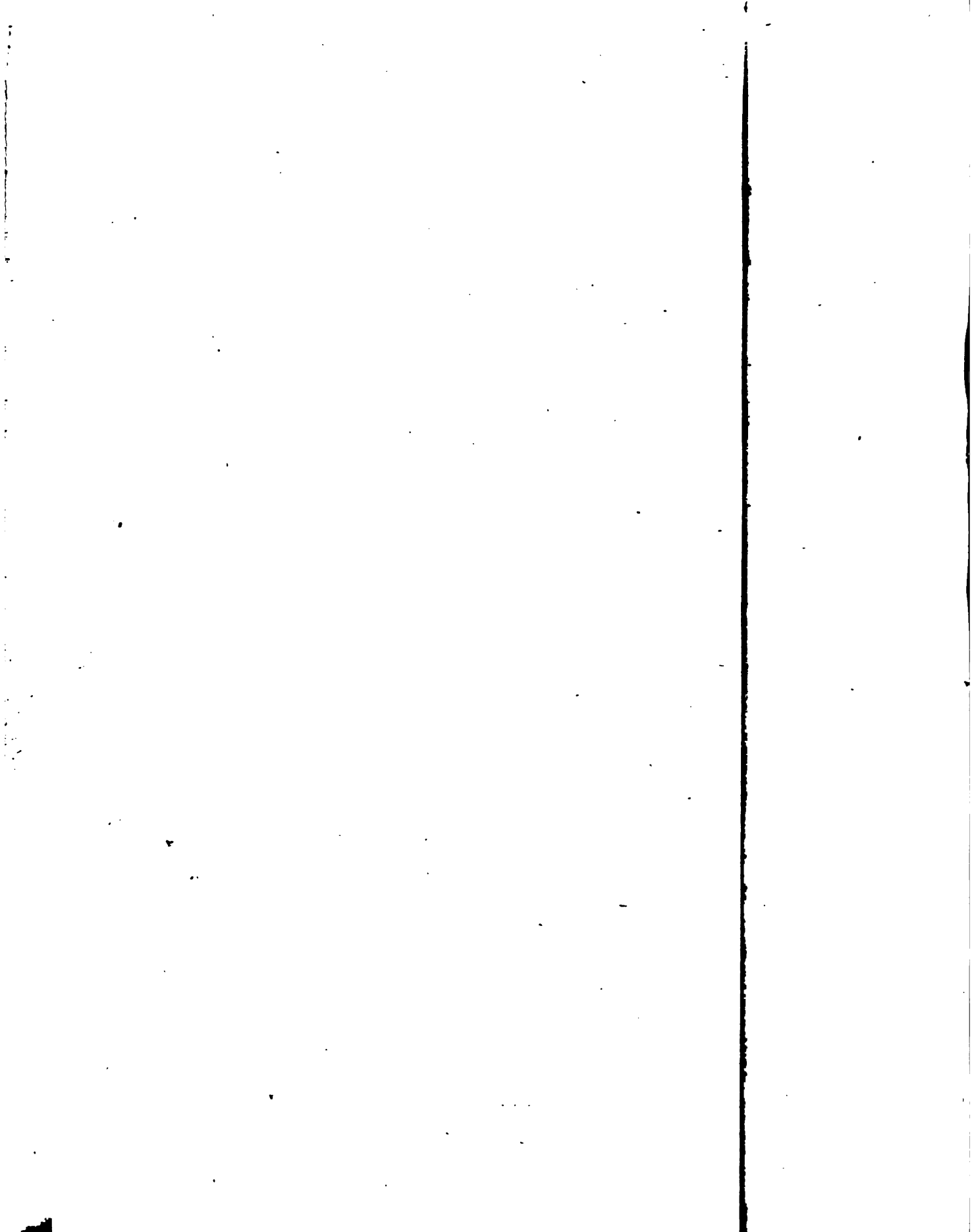














1870

1871

1872

1873

1874

1875

1876

1877

1878

1879

1880

1881

1882

1883

1884

1885

1886

1887

1888

1889

1890

1891

1892

1893

1894

1895

1896

1897

1898

1899

1900

1870

1871

1872

1873

1874

1875

1876

1877

1878

1879

1880

1881

1882

1883

1884

1885

1886

1887

1888

1889

1890

1891

1892

1893

1894

1895

1896

1897

1898

1899

1900

1870

1871

1872

1873

1874

1875

1876

1877

1878

1879

1880

1881

1882

1883

1884

1885

1886

1887

1888

1889

1890

1891

1892

1893

1894

1895

1896

1897

1898

1899

1900

A P P E N D I X:

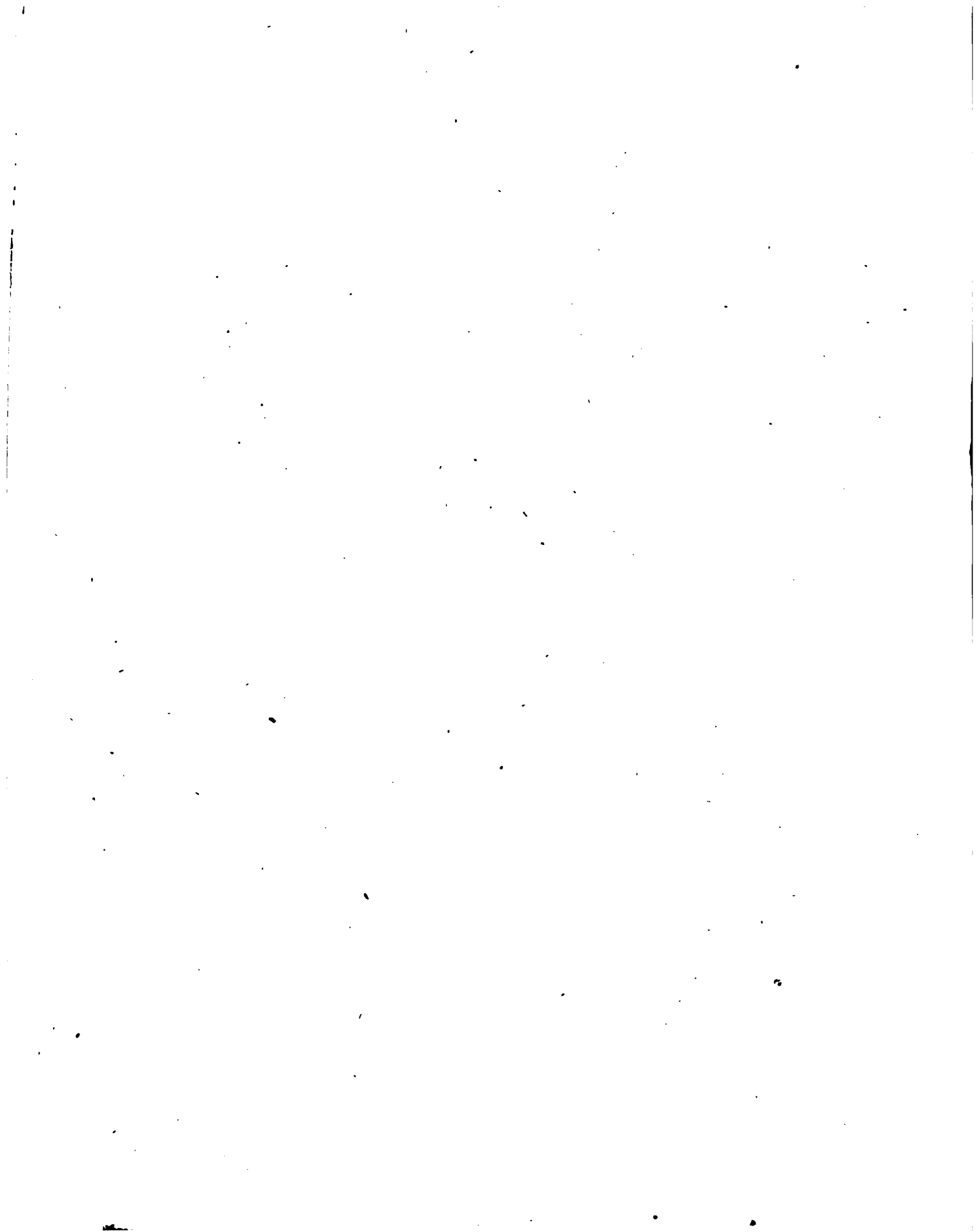
CONTAINING

TABLES of THEOREMS

FOR THE

CALCULATION of FLUENTS.





---

---

# TABLES of THEOREMS

FOR THE

CALCULATION of FLUENTS.

---

## T A B L E I.

### THEOREM I.

$$\dot{F} = x^n \dot{x}.$$

$$F = \frac{x^{n+1} - c^{n+1}}{n+1}.$$

$c$  the quantity to which  $x$  is equal when  $F = 0$ .

NOTE. When  $n$  is  $= -1$ , the expression for the value of

$F$  becomes  $= \text{Log. } \frac{x}{c}$ .

## T A B L E I.

## T H E O R E M II.

$$\dot{F} = \overline{a^n + bx^n}^p \times x^{n-1} \dot{x}.$$

$$F = \frac{\overline{a^n + bx^n}^{p+1} - \overline{a^n + bc^n}^{p+1}}{bn \cdot p + 1}.$$

$c$  the value of  $x$  when  $F = 0$ .

NOTE. When  $p$  is  $= -1$ , the expression for the value of

$$F \text{ becomes } = \frac{1}{bn} \times \text{Log.} \frac{a^n + bx^n}{a^n + bc^n}.$$

## T H E O R E M III.

$$\dot{F} = \frac{x^{np-1} \dot{x}}{a^n + bx^n}^{p+1}.$$

$$F = \frac{1}{npa^n} \times \frac{x^{np}}{a^n + bx^n}^p - \frac{c^{np}}{a^n + bc^n}^p.$$

$c$  the value of  $x$  when  $F = 0$ ,

NOTE. When  $p$  is  $= 0$ , the expression for the value of

$$F \text{ becomes } = \frac{1}{na^n} \times \text{Log.} \frac{x^n}{c^n} \times \frac{a^n + bc^n}{a^n + bx^n}.$$

T H E O-

## T H E O R E M I V.

$$\dot{F} = \frac{x^{n-1} \dot{x}}{\sqrt{2ax^n + x^{2n}}}$$

$$F = K + \frac{1}{n} \times \text{Log. } a + x^n + \sqrt{2ax^n + x^{2n}}.$$

NOTE. K here, and in the following theorems, denotes some invariable quantity; which will be determined in any equation whereto it may appertain, by properly substituting therein any known contemporary values of the variable quantities. And when K is so determined the equation may be said to be fitly adjusted, or corrected.

## T H E O R E M V.

$$\dot{F} = \frac{x^{n-1} \dot{x}}{\sqrt{ba^{2n} + x^{2n}}}$$

$$F = K + \frac{1}{n} \times \text{Log. } x^n + \sqrt{ba^{2n} + x^{2n}}.$$

## T H E O R E M VI.

$$\dot{F} = \frac{x^{n-1} \dot{x}}{a^{2n} - x^{2n}}$$

$$F = K + \frac{1}{2na^n} \times \text{Log. } \frac{a^n + x^n}{a^n - x^n}.$$

T H E O.

## T A B L E I.

## T H E O R E M VII.

$$\dot{F} = \frac{x^{-1} \dot{x}}{\sqrt{a^{2n} + bx^{2n}}}.$$

$$F = K + \frac{1}{2na^n} \times \text{Log.} \frac{a^n - \sqrt{a^{2n} + bx^{2n}}}{a^n + \sqrt{a^{2n} + bx^{2n}}}.$$


---

## T H E O R E M VIII.

$$\dot{F} = \frac{x^{n-1} \dot{x}}{\sqrt{2a^n x^n - x^{2n}}}.$$

$$F = K + \frac{1}{na^n} \times \text{Circ. Arc, rad. } a^n, \text{ versed sine } x^n.$$


---

## T H E O R E M IX.

$$\dot{F} = \frac{x^{n-1} \dot{x}}{\sqrt{a^{2n} - x^{2n}}}.$$

$$F = K + \frac{1}{na^n} \times \text{Circ. Arc, rad. } a^n, \text{ sine } x^n.$$


---

## T H E O R E M X.

$$\dot{F} = \frac{x^{n-1} \dot{x}}{a^{2n} + x^{2n}}.$$

$$F = K + \frac{1}{na^{2n}} \times \text{Circ. Arc, rad. } a^n, \text{ tang. } x^n.$$

T H E O-

T A B L E I.

7

T. H E O R E M XI.

$$\dot{F} = \frac{x^{-1} \dot{x}}{\sqrt{x^{2n} - a^{2n}}}$$

$$F = K + \frac{1}{na^{2n}} \times \text{Circ. Arc, rad. } a^n, \text{ fecant } x^n.$$

TABLE

---

T A B L E II.

CONTAINING

T H E O R E M S

FOR THE

CALCULATION of FLUENTS.

---

THEOREM I.

$$\dot{F} = \frac{x^m \dot{x}}{x+a}.$$

$m$  any positive integer.

$$F = K + a^m \times \frac{x^m}{m a^m} - \frac{x^{m-1}}{m-1 \cdot a^{m-1}} + \frac{x^{m-2}}{m-2 \cdot a^{m-2}} (m)^* \pm \text{Log. } x+a.$$

\* + or - according as  $m$  is even or odd.

---

THEOREM II.

$$\dot{F} = \frac{x^{-m} \dot{x}}{x+a}.$$

$m$  any positive integer.

$$F = K + \frac{1}{a^m} \times \frac{a^{m-1}}{m-1 \cdot x^{m-1}} - \frac{a^{m-2}}{m-2 \cdot x^{m-2}} + \frac{a^{m-3}}{m-3 \cdot x^{m-3}} (m-1)^* \pm \text{Log. } \frac{x+a}{x}.$$

\* + or - according as  $m-1$  is even or odd.

THEO.

T A B L E II.

9.

T H E O R E M III.

$$\dot{F} = \frac{x^m \dot{x}}{x^2 + a^2}.$$

$m$  any even positive number.

$$F = K + a^{m-2} \times \frac{x^{m-1}}{m-1 \cdot a^{m-2}} - \frac{x^{m-3}}{m-3 \cdot a^{m-4}} + \frac{x^{m-5}}{m-5 \cdot a^{m-6}} \left(\frac{m}{2}\right)^* \pm A.$$

$A = \text{Circ. Arc, rad. } a, \text{ tang. } x.$

\* + or — according as  $\frac{m}{2}$  is even or odd.

T H E O R E M IV.

$$\dot{F} = \frac{x^m \dot{x}}{x^2 - a^2}.$$

$m$  any even positive number.

$$F = K + a^{m-1} \times \frac{x^{m-1}}{m-1 \cdot a^{m-1}} + \frac{x^{m-3}}{m-3 \cdot a^{m-3}} + \frac{x^{m-5}}{m-5 \cdot a^{m-5}} \left(\frac{m}{2}\right) + \frac{1}{2} \text{Log. } \frac{a-x}{a+x}.$$

T H E O R E M V.

$$\dot{F} = \frac{x^m \dot{x}}{x^2 + a^2}.$$

$m$  any odd positive number.

$$F = K + a^{m-1} \times \frac{x^{m-1}}{m-1 \cdot a^{m-1}} - \frac{x^{m-3}}{m-3 \cdot a^{m-3}} + \frac{x^{m-5}}{m-5 \cdot a^{m-5}} \left(\frac{m-1}{2}\right)^* \pm \frac{1}{2} \text{Log. } \sqrt{x^2 + a^2}.$$

$a^2$  either positive or negative.

\* + or — according as  $\frac{m-1}{2}$  is even or odd.

b

T H E O.



## T H E O R E M VI.

$$\dot{F} = \frac{x^{-m} \dot{x}}{x^2 + a^2}$$

$m$  any even positive number.

$$F = K - \frac{1}{a^{m+1}} \times \frac{a^{m-1}}{m-1 \cdot x^{m-1}} - \frac{a^{m-3}}{m-3 \cdot x^{m-3}} + \frac{a^{m-5}}{m-5 \cdot x^{m-5}} \left(\frac{m}{2}\right)^* \pm A$$

$A = \text{Circ. Arc, rad. } 1, \text{ tang. } \frac{a}{x}$

\* + or - according as  $\frac{m}{2}$  is even or odd.

## T H E O R E M VII.

$$\dot{F} = \frac{x^{-m} \dot{x}}{x^2 - a^2}$$

$m$  any even positive number.

$$F = K + \frac{1}{a^{m+1}} \times \frac{a^{m-1}}{m-1 \cdot x^{m-1}} + \frac{a^{m-3}}{m-3 \cdot x^{m-3}} + \frac{a^{m-5}}{m-5 \cdot x^{m-5}} \left(\frac{m}{2}\right) + \frac{1}{2} \text{Log. } \frac{x-a}{x+a}$$

## T H E O R E M VIII.

$$\dot{F} = \frac{x^{-m} \dot{x}}{x^2 + a^2}$$

$m$  any odd positive number.

$$F = K - \frac{1}{a^{m+1}} \times \frac{a^{m-1}}{m-1 \cdot x^{m-1}} - \frac{a^{m-3}}{m-3 \cdot x^{m-3}} + \frac{a^{m-5}}{m-5 \cdot x^{m-5}} \left(\frac{m-1}{2}\right)^* \pm \text{Log. } \frac{x^2 + a^2}{x^2}$$

$a^2$  either positive or negative.

\* + or - according as  $\frac{m-1}{2}$  is even or odd.

TABLE

---

T A B L E III.

CONTAINING

T H E O R E M S

FOR THE

CALCULATION of FLUENTS.

---

THEOREM I.

$$\dot{F} = \frac{x^{-\frac{1}{2}}}{\sqrt{a^2 - x^2}}$$

$$F = K + \frac{4}{a^{\frac{1}{2}}} \times \overline{de - e' e''} = K + \frac{2}{a^{\frac{1}{2}}} \times \overline{de + DP - AD - L}$$

THEOREM II.

The fluent of  $\frac{x^{-\frac{1}{2}}}{\sqrt{a^2 - x^2}}$ , generated whilst  $x$  from 0 becomes equal to any quantity  $k$ , is equal to the fluent of the same fluxion, generated whilst  $x$  from  $a \times \frac{a-k}{a+k}$  becomes equal to  $a$ .

NOTE. All the theorems in this Table refer to the Scheme at the end of it, for the values of the quantities required.

## T H E O R E M III.

The fluent of  $\frac{x^{-\frac{1}{2}} \dot{x}}{\sqrt{a^2 - x^2}}$ , generated whilst  $x$  from  $o$  becomes equal to  $\sqrt{2^{\frac{1}{2}} - 1} \times a$ , is  $= \frac{M}{a^{\frac{1}{2}}}$ .

## T H E O R E M IV.

The whole fluent of  $\frac{x^{-\frac{1}{2}} \dot{x}}{\sqrt{a^2 - x^2}}$  is  $= \frac{2M}{a^{\frac{1}{2}}}$ .

## T H E O R E M V.

$$\dot{F} = \frac{x^{\frac{1}{2}} \dot{x}}{\sqrt{a^2 - x^2}}$$

$$F = K + \frac{2}{a^{\frac{1}{2}}} \times \overline{2e'e'' - de} = K + \frac{2}{a^{\frac{1}{2}}} \times \overline{L + AD - DP}.$$

## T H E O R E M VI.

The tangent  $eo$  ( $= \overline{ax}^{\frac{1}{2}} \times \overline{\frac{a-x}{a+x}}^{\frac{1}{2}}$ ) together with the fluent of  $\frac{\frac{1}{2} a^{\frac{1}{2}} x^{\frac{1}{2}} \dot{x}}{\sqrt{a^2 - x^2}}$ , generated whilst  $x$  from  $o$  becomes equal to any quantity  $k$ , is equal to the fluent of the same fluxion, generated whilst  $x$  from  $a \times \frac{a-k}{a+k}$  becomes equal to  $a$ .

T H E O-

T H E O R E M VII.

The fluent of  $\frac{x^{\frac{1}{2}}}{\sqrt{a^2 - x^2}}$ , generated whilst  $x$  from 0 becomes equal to  $2^{\frac{1}{2}} - 1 \times a$ , is  $= \frac{L}{a^{\frac{1}{2}}} - 2^{\frac{1}{2}} - 1 \times a^{\frac{1}{2}}$ .

T H E O R E M VIII.

The *whole* fluent of  $\frac{x^{\frac{1}{2}}}{\sqrt{a^2 - x^2}}$  is  $= \frac{2L}{a^{\frac{1}{2}}}$ .

T H E O R E M IX.

$$\dot{F} = \frac{y^{-\frac{1}{2}}}{\sqrt{y^2 - a^2}}$$

$$F = K + \frac{4}{a^{\frac{1}{2}}} \times \overline{ae - E'' + e'e''} = K + \frac{2}{a^{\frac{1}{2}}} \times \overline{ae + AD - DP}$$

$$x = \frac{a^2}{y}$$

T H E O R E M X.

The fluent of  $\frac{y^{-\frac{1}{2}}}{\sqrt{y^2 - a^2}}$ , generated whilst  $y$  from  $a$  becomes equal to  $2^{\frac{1}{2}} + 1 \times a$ , is  $= \frac{M}{a^{\frac{1}{2}}}$ .

T H E O R E M XI.

The *whole* fluent of  $\frac{y^{-\frac{1}{2}}}{\sqrt{y^2 - a^2}}$  is  $= \frac{2M}{a^{\frac{1}{2}}}$ .

T H E O -

## T H E O R E M XII.

$$\dot{F} = \frac{y^{\frac{1}{2}}j}{\sqrt{y^2 - a^2}}$$

$$F = K + \frac{2}{a^{\frac{1}{2}}} \times \overline{DP + ac + 2e'e'' - 2E''} = K + \frac{2}{a^{\frac{1}{2}}} \times AD.$$

$$x = \frac{a^2}{y}$$

## T H E O R E M XIII.

The fluent of  $\frac{y^{\frac{1}{2}}j}{\sqrt{y^2 - a^2}}$ , generated whilst  $y$  from  $a$  becomes equal to  $\overline{2^{\frac{1}{2}} + 1} \times a$ , is  $= \overline{2^{\frac{1}{2}} + 1} \times a^{\frac{1}{2}} - \frac{L}{a^{\frac{1}{2}}}$ .

NOTE. The *whole* fluent is infinite.

## T H E O R E M XIV.

$$\dot{F} = \frac{y^{-\frac{1}{2}}j}{\sqrt{a^2 + y^2}}$$

$$F = K + \frac{2^{\frac{1}{2}}}{a^{\frac{1}{2}}} \times \overline{ac + e'e'' - E''} = K + \frac{2^{\frac{1}{2}}}{a^{\frac{1}{2}}} \times \overline{ac + AD - DP}.$$

$$x = \overline{a^2 + y^2}^{\frac{1}{2}} - y$$

## T H E O R E M XV.

The fluent of  $\frac{y^{-\frac{1}{2}}j}{\sqrt{a^2 + y^2}}$ , generated whilst  $y$  from 0 becomes equal to  $a$ , is  $= \frac{2^{\frac{1}{2}}}{a^{\frac{1}{2}}} \times M$ .

T H E O-

T H E O R E M XVI.

The *whole* fluent of  $\frac{y^{-\frac{1}{2}}j}{\sqrt{a^2+y^2}}$  is  $= \frac{2^{\frac{1}{2}}}{a^{\frac{1}{2}}} \times M.$

T H E O R E M XVII.

$$\dot{F} = \frac{y^{\frac{1}{2}}j}{\sqrt{a^2+y^2}}.$$

$$F = K + \frac{2^{\frac{1}{2}}}{a^{\frac{1}{2}}} \times \overline{ac + 2e'e'' - 2E'' + \frac{1}{2}DP} = K + \frac{2^{\frac{1}{2}}}{a^{\frac{1}{2}}} \times \overline{AD - \frac{1}{2}DP}.$$

$$.x = \overline{a^2 + y^2}^{\frac{1}{2}} - y.$$

T H E O R E M XVIII.

The fluent of  $\frac{y^{\frac{1}{2}}j}{\sqrt{a^2+y^2}}$ , generated whilst  $y$  from  $\sigma$  becomes equal to  $a$ , is  $= \sqrt{2a} - \frac{2^{\frac{1}{2}}}{a} \times L.$

NOTE. The *whole* fluent is infinite.

T H E O R E M XIX.

$$\dot{F} = \frac{j}{\overline{a^2 - y^2}^{\frac{1}{2}}}.$$

$$F = K + \frac{2}{a^{\frac{1}{2}}} \times \overline{2E'' - 2e'e'' - ac} = K + \frac{2}{a^{\frac{1}{2}}} \times \overline{DP - AD}.$$

$$.x = \sqrt{a^2 - y^2}.$$

T H E O-

## T H E O R E M XX.

The fluent of  $\frac{y}{a^2 - y^2}^{\frac{1}{2}}$ , generated whilst  $y$  from  $o$  becomes equal to  $\sqrt{2^{\frac{1}{2}} - 2} \times a$ , is  $= \frac{L}{a^{\frac{1}{2}}} + 2^{\frac{1}{2}} - 1 \times a^{\frac{1}{2}}$ .

## T H E O R E M XXI.

The *whole* fluent of  $\frac{y}{a^2 - y^2}^{\frac{1}{2}}$  is  $= \frac{2L}{a^{\frac{1}{2}}}$ .

## T H E O R E M XXII.

$$\dot{F} = \frac{y^{\frac{1}{2}} y}{a^2 - y^2}^{\frac{1}{2}}$$

$$F = K + 2^{\frac{1}{2}} \times \frac{2e'' - 2e'e'' - ac - \overline{ay - y^2}}{\overline{ay - y^2}}^{\frac{1}{2}} \times \frac{\overline{a - y}}{\overline{a + y}}^{\frac{1}{2}}$$

$$= K + 2^{\frac{1}{2}} \times \frac{DP - AD - \overline{ay - y^2}}{\overline{ay - y^2}}^{\frac{1}{2}} \times \frac{\overline{a - y}}{\overline{a + y}}^{\frac{1}{2}}$$

$$x = a \times \frac{\overline{a - y}}{\overline{a + y}}^{\frac{1}{2}}$$

## T H E O R E M XXIII.

The fluent of  $\frac{y^{\frac{1}{2}} y}{a^2 - y^2}^{\frac{1}{2}}$ , generated whilst  $y$  from  $o$  becomes equal to  $\frac{a}{\sqrt{2}}$ , is  $= \frac{L}{\sqrt{2}}$ .

## T H E O R E M XXIV.

The *whole* fluent of  $\frac{y^{\frac{1}{2}} y}{a^2 - y^2}^{\frac{1}{2}}$  is  $= 2^{\frac{1}{2}} L$ .

T H E O.

T H E O R E M XXV.

$$\dot{F} = \frac{j}{y^2 - a^2}.$$

$$F = K + \frac{2^{\frac{1}{2}}}{a^{\frac{1}{2}}} \times \overline{ae + 2e'e'' - 2E'' + \frac{1}{2}DP} = K + \frac{2^{\frac{1}{2}}}{a^{\frac{1}{2}}} \times \overline{AD - \frac{1}{2}DP}.$$

$$x = y - \sqrt{y^2 - a^2}.$$

T H E O R E M XXVI.

The fluent of  $\frac{j}{y^2 - a^2}$ , generated whilst  $y$  from  $a$  becomes equal to  $2^{\frac{1}{2}}a$ , is  $= 2^{\frac{1}{2}} \times \overline{a^{\frac{1}{2}} - \frac{L}{a^{\frac{1}{2}}}}$ .

NOTE. The whole fluent is infinite.

T H E O R E M XXVII.

$$\dot{F} = \frac{y^{\frac{1}{2}}j}{y^2 - a^2}.$$

$$F = K + de - 2e'e'' + \frac{y^2 - a^2}{y^{\frac{1}{2}}} \quad \left. \begin{array}{l} \text{* A 3d proportional} \\ \text{to DP, CP.} \end{array} \right\}$$

$$= K + DP - AD - L + \frac{y^2 - a^2}{y^{\frac{1}{2}}}.$$

$$x = \frac{a\sqrt{y^2 - a^2}}{y}.$$



## T H E O R E M XXVIII.

The fluent of  $\frac{y^{\frac{1}{2}}j}{y^2 - a^2}^{\frac{1}{2}}$ , generated whilst  $y$  from  $a$  becomes equal to  $\frac{2^{\frac{1}{2}} + 1}{2}^{\frac{1}{2}} \times a$ , is  $= \frac{1}{2} \times a - L$ .

NOTE. The *whole* fluent is infinite.

## T H E O R E M XXIX.

$$\dot{F} = \frac{j}{a^2 + y^2}^{\frac{1}{2}}$$

$$F = K + \frac{2}{a^{\frac{1}{2}}} \times ac + ze'e'' - zE'' + DP = K + \frac{2}{a^{\frac{1}{2}}} \times AD.$$

$$x = \frac{a^2}{\sqrt{a^2 + y^2}}$$

## T H E O R E M XXX.

The fluent of  $\frac{j}{a^2 + y^2}^{\frac{1}{2}}$ , generated whilst  $y$  from  $o$  becomes equal to  $\sqrt{2 + \sqrt{2}} \times a$ , is  $= \frac{1}{2^{\frac{1}{2}} + 1} \times a^{\frac{1}{2}} - \frac{L}{a^{\frac{1}{2}}}$ .

NOTE. The *whole* fluent is infinite.

T H E O-

T H E O R E M XXXI.

$$\dot{F} = \frac{y^{\frac{1}{2}}j}{a^2 + y^2}^{\frac{1}{2}}$$

$$F = K + dc - 2e'e'' + \frac{y^{\frac{3}{2}}}{a^2 + y^2}^* \quad \left. \vphantom{\frac{y^{\frac{3}{2}}}{a^2 + y^2}^*} \right\} \text{* A 3d proportional to DP, CP.}$$

$$= K + DP - AD - L + \frac{y^{\frac{3}{2}}}{a^2 + y^2}^*$$

$$x = \frac{ay}{\sqrt{a^2 + y^2}}$$

T H E O R E M XXXII.

The fluent of  $\frac{y^{\frac{1}{2}}j}{a^2 + y^2}^{\frac{1}{2}}$ , generated whilst  $y$  from 0 becomes equal to  $\sqrt{\frac{1}{\sqrt{2}} - \frac{1}{2}} \times a$ , is  $= \frac{1}{2} \times a - L$ .

NOTE. The *whole* fluent is infinite.

T H E O R E M XXXIII.

$$\dot{F} = \frac{j}{a^2 - y^2}^{\frac{1}{2}}$$

$$F = K + \frac{4}{a^2} \times \overline{ae + e'e'' - E''} = K + \frac{2}{a^2} \times \overline{ae + AD - DP}$$

$$x = \sqrt{a^2 - y^2}$$

## T H E O R E M XXXIV.

The fluent of  $\frac{j}{a^2 - y^2}^{\frac{1}{2}}$ , generated whilst  $y$  from 0 becomes equal to  $\sqrt{2 - \sqrt{2}} \times a$ , is  $= \frac{M}{a^{\frac{1}{2}}}$ .

## T H E O R E M XXXV.

The *whole* fluent of  $\frac{j}{a^2 - y^2}^{\frac{1}{2}}$  is  $= \frac{2M}{a^{\frac{1}{2}}}$ .

## T H E O R E M XXXVI.

$$F = \frac{y^{-\frac{1}{2}} j}{a^2 - y^2}^{\frac{1}{2}}$$

$$F = K + \frac{2^{\frac{1}{2}}}{a^2} \times \overline{ae + e'e'' - E''} = K + \frac{2^{\frac{1}{2}}}{a^2} \times \overline{ae + AD - DP.}$$

$$x = a \times \left. \frac{a-y}{a+y} \right|^{\frac{1}{2}}$$

## T H E O R E M XXXVII.

The fluent of  $\frac{y^{-\frac{1}{2}} j}{a^2 - y^2}^{\frac{1}{2}}$ , generated whilst  $y$  from 0 becomes equal to  $\frac{a}{\sqrt{2}}$ , is  $= \frac{2^{\frac{1}{2}}}{a^2} \times M$ .

## T H E O R E M XXXVIII.

The *whole* fluent of  $\frac{y^{-\frac{1}{2}} j}{a^2 - y^2}^{\frac{1}{2}}$  is  $= \frac{2^{\frac{1}{2}}}{a^2} \times M$ .

T H E O -

T H E O R E M XXXIX.

$$\dot{F} = \frac{j}{y^2 - a^2)^{\frac{1}{2}}}$$

$$F = K + \frac{2^{\frac{1}{2}}}{a^{\frac{1}{2}}} \times \overline{ac + e'e'' - E''} = K + \frac{2^{\frac{1}{2}}}{a^{\frac{1}{2}}} \times \overline{ac + AD - DP.}$$

$$x = y - \sqrt{y^2 - a^2}.$$

T H E O R E M XL.

The fluent of  $\frac{j}{y^2 - a^2)^{\frac{1}{2}}}$ , generated whilst  $y$  from  $a$  becomes equal to  $2^{\frac{1}{2}} \times a$ , is  $= \frac{2^{\frac{1}{2}}}{a^{\frac{1}{2}}} \times M.$

T H E O R E M XLI.

The *whole* fluent of  $\frac{j}{y^2 - a^2)^{\frac{1}{2}}}$  is  $= \frac{2}{a} \Big)^{\frac{1}{2}} \times M.$

T H E O R E M XLII.

$$\dot{F} = \frac{y^{-\frac{1}{2}} j}{y^2 - a^2)^{\frac{1}{2}}}$$

$$F = K + \frac{4}{a^2} \times \overline{dc - e'e''} = K + \frac{2}{a^2} \times \overline{dc + DP - AD - L.}$$

$$y = \frac{a\sqrt{y^2 - a^2}}{y}.$$

T H E O-

## T H E O R E M XLIII.

The fluent of  $\frac{y^{-\frac{1}{2}}j}{y^2 - a^2}^{\frac{1}{2}}$ , generated whilst  $y$  from  $a$  becomes equal to  $\frac{1}{2^{\frac{1}{2}}} + \frac{1}{2}^{\frac{1}{2}} \times a$ , is  $= \frac{M}{a^2}$ .

## T H E O R E M XLIV.

The *whole* fluent of  $\frac{y^{-\frac{1}{2}}j}{y^2 - a^2}^{\frac{1}{2}}$  is  $= \frac{2M}{a^2}$ .

## T H E O R E M XLV.

$$F = \frac{j}{a^2 + y^2}^{\frac{1}{2}}$$

$$F = K + \frac{4}{a^{\frac{1}{2}}} \times \overline{ac + ce'' - E''} = K + \frac{2}{a^{\frac{1}{2}}} \times \overline{ac + AD - DP.}$$

$$x = \frac{a^2}{\sqrt{a^2 + y^2}}$$

## T H E O R E M XLVI.

The fluent of  $\frac{j}{a^2 + y^2}^{\frac{1}{2}}$ , generated whilst  $y$  from 0 becomes equal to  $\sqrt{2 + \sqrt{2}} \times a$ , is  $= \frac{M}{a^{\frac{1}{2}}}$ .

T H E O.

T H E O R E M XLVII.

The whole fluent of  $\frac{j}{a^2 + y^2)^{\frac{1}{2}}}$  is  $= \frac{2M}{a^{\frac{1}{2}}}$ .

---

T H E O R E M XLVIII.

$$\dot{F} = \frac{y^{-\frac{1}{2}}j}{a^2 + y^2)^{\frac{1}{2}}}$$

$$F = K + \frac{4}{a^2} \times \overline{de - e'e''} = K + \frac{2}{a^2} \times \overline{de + DP - AD - L}$$

$$x = \frac{ay}{\sqrt{a^2 + y^2}}$$

T H E O R E M XLIX.

The fluent of  $\frac{y^{-\frac{1}{2}}j}{a^2 + y^2)^{\frac{1}{2}}}$ , generated whilst  $y$  from  $a$  becomes equal to  $\sqrt{\frac{1}{\sqrt{2}} - \frac{1}{2}} \times a$ , is  $= \frac{M}{a^{\frac{1}{2}}}$ .

T H E O R E M L.

The whole fluent of  $\frac{y^{-\frac{1}{2}}j}{a^2 + y^2)^{\frac{1}{2}}}$  is  $= \frac{2M}{a^{\frac{1}{2}}}$ .

S C H E M E.

# S C H E M E

FOR

## T A B L E III.

$d$   $\left\{ \begin{array}{l} = \frac{1}{4} \text{ of the periphery of a circle whose radius is } 1. \\ = 1.57079632. \end{array} \right.$

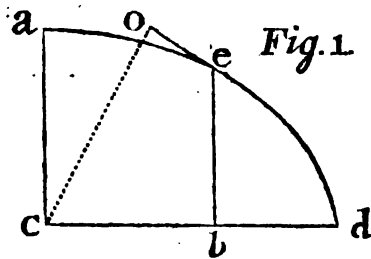
$aed$  (Fig. 1.) is a quadrantal arc of an ellipsis =  $E'$ .

Semi-transverse axis  $cd = 2^{\frac{1}{2}}a$ .

Semi-conjugate axis  $ac = a$ .

Abcissa  $cb = 2^{\frac{1}{2}}\sqrt{a^2 - ax}$ .

Ordinate  $be = \sqrt{ax}$ .

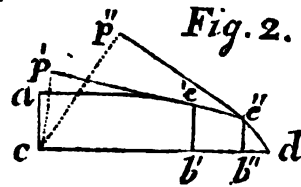


$e$   $\left\{ \begin{array}{l} = \text{the value of } E' \text{ when } a \text{ is } = 1. \\ = 1.91009889. \end{array} \right.$

$ae'e'd$  (Fig. 2.) is a quadrantal arc of another ellipsis =  $E''$ .

Semi-transverse axis  $cd = \frac{1}{\sqrt{2}} + \frac{1}{2} \times a$ .

Semi-conjugate axis  $ac = \frac{1}{\sqrt{2}} - \frac{1}{2} \times a$ .



$ep'$  and its equal  $e''p''$  (each =  $\sqrt{a^2 - ax}$ ) are tangents, to which  $cp'$ ,  $cp''$  are perpendiculars.

The abscissa  $cb'$ , or  $cb''$ , corresponding to the ordinate  $be$ ,

or  $be''$ , is =  $\frac{2^{\frac{1}{2}} + 1}{2^{\frac{1}{2}}} \times a^{\frac{1}{2}} \sqrt{2^{\frac{1}{2}}a + a - x \mp \sqrt{ax + x^2}}$ .

$f$   $\left\{ \begin{array}{l} = \text{the value of } E'' \text{ when } a \text{ is } = 1. \\ = 1.2545845059. \end{array} \right.$

SCHEME

S C H E M E for TABLE III. continued.

AD (Fig. 3.) is an equilateral hyperbola, whose vertex is A and center C.

DP is a tangent, to which CP is perpendicular.

$$AC = a.$$

$$CP = \sqrt{ax}.$$

$$DP = \frac{a}{x} \sqrt{x^2 - a^2}.$$

Abcissa CAB (corresponding to the ordinate BD)

$$= a \times \frac{a+x}{2x}.$$

L, the *limit* of DP-AD, is  $= 2E'' - E' = a \times 2f - e = \frac{a}{2} \times e - \sqrt{e^2 - 2d}$

$$= .5990701173 \times a.$$

$$M = 2 \times E' - E'' = 2a \times e - f = \frac{a}{2} \times e + \sqrt{e^2 - 2d}$$

$$= 1.3110287771 \times a.$$

NOTE. *All the Theorems in TABLE III. refer to this Scheme.*

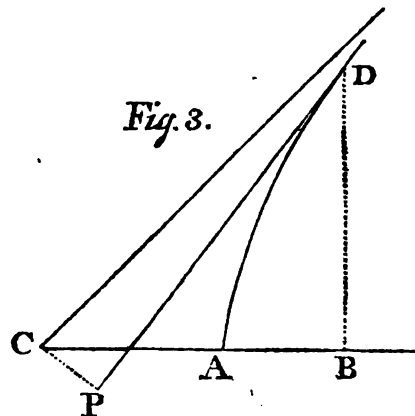


Fig. 3.



---

---

T A B L E IV.

CONTAINING

T H E O R E M S

FOR THE

CALCULATION of FLUENTS.

---

T H E O R E M I.

$$\dot{F} = \frac{x^{-i} \dot{x}}{a^2 - x^2} = \frac{\frac{1}{2} y^{-i} \dot{y}}{b^2 - y^2}$$

$$F = K - \frac{1}{2} a^{-i} B.$$

$$x = \frac{a^i - x^i}{x^i} = \frac{b - y}{y}.$$

T H E O R E M II.

The *whole* fluent of  $\frac{x^{-i} \dot{x}}{a^2 - x^2}$  is  $= \frac{3^i a^{-i} p}{3^i + 1}$ .

NOTE. The necessary explanation, respecting the values of the quantities concerned in the theorems in this Table, is given at the end of it.

T H E O -

T H E O R E M III.

$$\dot{F} = \frac{x^{-\frac{1}{2}}x}{a^2 - x^2} = \frac{\frac{1}{2}y}{b^2 - y^2}$$

$$F = K - \frac{1}{2}a^{-\frac{1}{2}}D.$$

$$z = \frac{a^{\frac{1}{2}} - x^{\frac{1}{2}}}{a^{\frac{1}{2}}} = \frac{b - y}{b}$$

T H E O R E M IV.

The whole fluent of  $\frac{x^{-\frac{1}{2}}x}{a^2 - x^2}$  is  $= \frac{3^{\frac{1}{2}}a^{-\frac{1}{2}}P}{3^{\frac{1}{2}} + 1}$ .

---

T H E O R E M V.

$$\dot{F} = \frac{x^{\frac{1}{2}}x}{a^2 - x^2} = \frac{\frac{1}{2}y}{b^2 - y^2}$$

$$F = K + \frac{1}{2}a^{\frac{1}{2}} \times C - D.$$

$$z = \frac{a^{\frac{1}{2}} - x^{\frac{1}{2}}}{a^{\frac{1}{2}}} = \frac{b - y}{b}$$

T H E O R E M VI.

The whole fluent of  $\frac{x^{\frac{1}{2}}x}{a^2 - x^2}$  is  $= 3^{\frac{1}{2}}a^{\frac{1}{2}}Q.$

## T H E O R E M VII.

$$\dot{F} = \frac{x^{\frac{1}{2}} \dot{x}}{a^2 - x^2} = \frac{\frac{1}{2} y^{\frac{1}{2}} \dot{y}}{b^2 - y^2}$$

$$F = K + \frac{1}{2} a^{\frac{1}{2}} \times A + B - \frac{1}{2} x^{-\frac{1}{2}} \sqrt{a^2 - x^2}$$

$$z = \frac{a^{\frac{1}{2}} - x^{\frac{1}{2}}}{x^{\frac{1}{2}}} = \frac{b - y}{y}$$

## T H E O R E M VIII.

The *whole* fluent of  $\frac{x^{\frac{3}{2}} \dot{x}}{a^2 - x^2}$  is  $= \frac{3^{\frac{1}{2}} a^{\frac{1}{2}} Q}{2}$ .

## T H E O R E M IX.

$$\dot{F} = \frac{x^{-\frac{1}{2}} \dot{x}}{x^2 - a^2} = \frac{\frac{1}{2} y^{-\frac{1}{2}} \dot{y}}{y^2 - b^2}$$

$$F = K + \frac{1}{2} a^{-\frac{1}{2}} D.$$

$$z = \frac{x^{\frac{1}{2}} - a^{\frac{1}{2}}}{x^{\frac{1}{2}}} = \frac{y - b}{y}$$

## T H E O R E M X.

The *whole* fluent of  $\frac{x^{-\frac{3}{2}} \dot{x}}{x^2 - a^2}$  is  $= \frac{3^{\frac{1}{2}} a^{-\frac{1}{2}} P}{3^{\frac{1}{2}} + 1}$ .

T H E O.

## T H E O R E M XI.

$$\dot{F} = \frac{x^{-\frac{1}{2}}}{x^2 - a^2} = \frac{\frac{1}{2}y}{y^2 - b^2}$$

$$F = K + \frac{1}{2}a^{-\frac{1}{2}}B.$$

$$x = \frac{x^{\frac{3}{2}} - a^{\frac{3}{2}}}{a^{\frac{3}{2}}} = \frac{y - b}{b}.$$

## T H E O R E M XII.

The whole fluent of  $\frac{x^{-\frac{1}{2}}}{x^2 - a^2}$  is  $= \frac{3^{\frac{1}{2}}a^{-\frac{1}{2}}P}{3^{\frac{1}{2}} + 1}$ .

---

## T H E O R E M XIII.

$$\dot{F} = \frac{x^{\frac{1}{2}}}{x^2 - a^2} = \frac{\frac{1}{2}y}{y^2 - b^2}$$

$$F = K + \frac{1}{2}a^{\frac{1}{2}} \times \overline{A + B}.$$

$$x = \frac{x^{\frac{3}{2}} - a^{\frac{3}{2}}}{a^{\frac{3}{2}}} = \frac{y - b}{b}.$$

NOTE. The whole fluent is infinite.

T H E O-

## T A B L E IV.

## T H E O R E M XIV.

$$\dot{F} = \frac{x^{\frac{3}{2}} \dot{x}}{(x^2 - a^2)^{\frac{3}{2}}} = \frac{\frac{3}{2} y^{\frac{3}{2}} \dot{y}}{(y^2 - b^2)^{\frac{3}{2}}}$$

$$F = K + \frac{3}{2} a^{\frac{3}{2}} \times \overline{C - D} + \frac{3}{2} x^{-\frac{1}{2}} \sqrt{x^2 - a^2}.$$

$$z = \frac{x^{\frac{3}{2}} - a^{\frac{3}{2}}}{x^{\frac{3}{2}}} = \frac{y - b}{y}.$$

NOTE. The *whole* fluent is infinite.

## T H E O R E M XV.

$$\dot{F} = \frac{x^{-\frac{3}{2}} \dot{x}}{(a^2 + x^2)^{\frac{3}{2}}} = \frac{\frac{3}{2} y^{-\frac{3}{2}} \dot{y}}{(b^2 + y^2)^{\frac{3}{2}}}$$

$$F = K - \frac{3}{2} a^{-\frac{3}{2}} D.$$

$$z = \frac{a^{\frac{3}{2}} + x^{\frac{3}{2}}}{x^{\frac{3}{2}}} = \frac{b + y}{y}.$$

## T H E O R E M XVI.

The *whole* fluent of  $\frac{x^{-\frac{3}{2}} \dot{x}}{(a^2 + x^2)^{\frac{3}{2}}}$  is  $= \frac{2 \cdot 3^{\frac{1}{2}} a^{-\frac{3}{2}} P}{3^{\frac{1}{2}} + 1}$ .

THEO-

## T H E O R E M XVII.

$$\dot{F} = \frac{x^{-\frac{1}{2}} \dot{x}}{a^2 + x^2} = \frac{\frac{1}{2} \dot{y}}{b^2 + y^2}$$

$$F = K + \frac{1}{2} a^{-1} D.$$

$$z = \frac{a^{\frac{1}{2}} + x^{\frac{1}{2}}}{a^{\frac{1}{2}}} = \frac{b + y}{b}$$

## T H E O R E M XVIII.

The *whole* fluent of  $\frac{x^{-\frac{1}{2}} \dot{x}}{a^2 + x^2}$  is  $= \frac{2 \cdot 3^{\frac{1}{2}} a^{-1} P}{3^{\frac{1}{2}} + 1}$ .

## T H E O R E M XIX.

$$\dot{F} = \frac{x^{\frac{1}{2}} \dot{x}}{a^2 + x^2} = \frac{\frac{1}{2} y \dot{y}}{b^2 + y^2}$$

$$F = K + \frac{1}{2} a^{\frac{1}{2}} \times C - D.$$

$$z = \frac{a^{\frac{1}{2}} + x^{\frac{1}{2}}}{a^{\frac{1}{2}}} = \frac{b + y}{b}$$

NOTE. The *whole* fluent is infinite.

THEO.

## T H E O R E M XX.

$$\dot{F} = \frac{x^{\frac{3}{2}} \dot{x}}{a^2 + x^2)^{\frac{3}{2}}} = \frac{\frac{3}{2} y^{\frac{1}{2}} \dot{y}}{b^2 + y^2)^{\frac{3}{2}}}$$

$$F = K - \frac{3}{2} a^{\frac{3}{2}} \times C - D + \frac{3}{2} x^{-\frac{1}{2}} \sqrt{a^2 + x^2}.$$

$$z = \frac{a^{\frac{3}{2}} + x^{\frac{3}{2}}}{x^{\frac{3}{2}}} = \frac{b + y}{y}.$$

NOTE. The *whole* fluent is infinite.

## T H E O R E M XXI.

$$\dot{F} = \frac{x^{-\frac{3}{2}} \dot{x}}{a^2 - x^2)^{\frac{3}{2}}} = \frac{\frac{3}{2} y^{-\frac{1}{2}} \dot{y}}{b^2 - y^2)^{\frac{3}{2}}}$$

$$F = - \text{fl.} \frac{w^{-\frac{1}{2}} \dot{w}}{4a^2 + w^2)^{\frac{3}{2}}}, \text{ to be found by theorem xvii.}$$

$$w = \frac{a^2 - x^2}{x} = \frac{b^2 - y^2}{y^{\frac{1}{2}}}.$$

## T H E O R E M XXII.

The *whole* fluent of  $\frac{x^{-\frac{3}{2}} \dot{x}}{a^2 - x^2)^{\frac{3}{2}}}$  is  $= \frac{2^{\frac{2}{3}} 3^{\frac{1}{3}} a^{-\frac{1}{3}} P}{3^{\frac{1}{3}} + 1}$ .

T H E O-

T H E O R E M XXIII.

$$\dot{F} = \frac{x^{\frac{1}{2}}}{a^2 - x^2} = \frac{\frac{1}{2}y^{\frac{1}{2}}}{b^2 - y^2}$$

$$F = -\text{fl.} \frac{w^{\frac{1}{2}}w}{a^2 - w^2}, \text{ to be found by theor. v.}$$

$$w = \frac{a^2 - x^2}{x} = \frac{b^2 - y^2}{y}$$

T H E O R E M XXIV.

The *whole* fluent of  $\frac{x^{\frac{1}{2}}}{a^2 - x^2}$  is  $= 3^{\frac{1}{2}}a^{\frac{1}{2}}Q.$

T H E O R E M XXV.

$$\dot{F} = \frac{x^{\frac{3}{2}}}{a^2 - x^2} = \frac{\frac{3}{2}y^{\frac{3}{2}}}{b^2 - y^2}$$

$$F = -\frac{1}{2}w^{\frac{3}{2}} + \frac{1}{2}\text{fl.} \frac{w^{\frac{1}{2}}w}{4a^2 + w^2}, \text{ to be found by theor. xx.}$$

$$w = \frac{a^2 - x^2}{x} = \frac{b^2 - y^2}{y}$$

T H E O R E M XXVI.

The *whole* fluent of  $\frac{x^{\frac{3}{2}}}{a^2 - x^2}$  is  $= \frac{3^{\frac{1}{2}}a^{\frac{1}{2}}Q}{2^{\frac{1}{2}}}.$



## T A B L E I V.

## T H E O R E M XXVII.

$$\dot{F} = \frac{x^{\frac{3}{2}} \dot{x}}{a^2 - x^2} = \frac{\frac{3}{2} y^{\frac{3}{2}} \dot{y}}{b^2 - y^2}$$

$$F = \frac{a^{\frac{1}{2}} w^{\frac{3}{2}}}{a^2 + w^2} - \frac{3}{2} a^{\frac{1}{2}} \text{fl.} \frac{w^{\frac{3}{2}} \dot{w}}{a^2 + w^2}, \text{ to be found by theor. xx.}$$

$$w = \frac{ax}{a^2 - x^2} = \frac{b^{\frac{3}{2}} y^{\frac{3}{2}}}{b^2 - y^2}$$

## T H E O R E M XXVIII.

The *whole* fluent of  $\frac{x^{\frac{3}{2}} \dot{x}}{a^2 - x^2}$  is  $= \frac{2^{\frac{1}{2}} a Q}{3^{\frac{1}{2}}}$ .

---

## T H E O R E M XXIX.

$$\dot{F} = \frac{x^{-\frac{1}{2}} \dot{x}}{x^2 - a^2} = \frac{\frac{1}{2} y^{-\frac{1}{2}} \dot{y}}{y^2 - b^2}$$

$$F = \text{fl.} \frac{w^{-\frac{1}{2}} \dot{w}}{4a^2 + w^2}, \text{ to be found by theor. xvii.}$$

$$w = \frac{x^2 - a^2}{x} = \frac{y^2 - b^2}{y^{\frac{1}{2}}}$$

## T H E O R E M XXX.

The *whole* fluent of  $\frac{x^{-\frac{1}{2}} \dot{x}}{x^2 - a^2}$  is  $= \frac{2^{\frac{1}{2}} 3^{\frac{1}{2}} a^{-\frac{1}{2}} P}{3^{\frac{1}{2}} + 1}$

T H E O-

T H E O R E M XXXI.

$$\dot{F} = \frac{x}{x^2 - a^2} = \frac{\frac{1}{2}y^{\frac{1}{2}}j}{y^2 - b^2}$$

F = fl.  $\frac{w^{\frac{1}{2}}\dot{w}}{a^2 + w^2}$ , to be found by theor. XIX.

$$w = \sqrt{x^2 - a^2} = \sqrt{y^2 - b^2}$$

NOTE. The *whole* fluent is infinite.

T H E O R E M XXXII.

$$\dot{F} = \frac{x^{\frac{1}{2}}x}{x^2 - a^2} = \frac{\frac{1}{2}yy}{y^2 - b^2}$$

F =  $\frac{1}{2}w^{\frac{3}{2}} + \frac{1}{2}$  fl.  $\frac{w^{\frac{3}{2}}\dot{w}}{a^2 + w^2}$ , to be found by theor. XX.

$$w = \frac{x^2 - a^2}{x} = \frac{y^2 - b^2}{y^{\frac{1}{2}}}$$

NOTE. The *whole* fluent is infinite.

## T H E O R E M XXXIII.

$$\dot{F} = \frac{x^{\frac{1}{2}} \dot{x}}{x^2 - a^2} = \frac{\frac{1}{2} y^{\frac{1}{2}} \dot{y}}{y^2 - b^2}$$

$$F = \frac{a^{\frac{1}{2}} w^{\frac{1}{2}}}{w^2 - a^2} - \frac{1}{2} a^{\frac{1}{2}} \text{fl. } \frac{w^{\frac{1}{2}} \dot{w}}{w^2 - a^2}, \text{ to be found by theor. xiv.}$$

$$w = \frac{ax}{x^2 - a^2} = \frac{b^{\frac{1}{2}} y^{\frac{1}{2}}}{y^2 - b^2}$$

NOTE. The *whole* fluent is infinite.

## T H E O R E M XXXIV.

$$\dot{F} = \frac{x^{-\frac{1}{2}} \dot{x}}{a^2 + x^2} = \frac{\frac{1}{2} y^{-\frac{1}{2}} \dot{y}}{b^2 + y^2}$$

$$F = \pm \text{fl. } \frac{w^{-\frac{1}{2}} \dot{w}}{w^2 - 4a^2}, \text{ to be found by theor. xi.}$$

$$w = \frac{a^2 + x^2}{x} = \frac{b^2 + y^2}{y^{\frac{1}{2}}}$$

## T H E O R E M XXXV.

$$\text{The } \textit{whole} \text{ fluent of } \frac{x^{-\frac{1}{2}} \dot{x}}{a^2 + x^2} \text{ is } = \frac{2^{\frac{1}{2}} 3^{\frac{1}{2}} a^{-\frac{1}{2}} P}{3^{\frac{1}{2}} + 1}$$

THEO-

## T H E O R E M XXXVI.

$$\dot{F} = \frac{x}{a^2 + x^2} = \frac{\frac{1}{2}j^{\frac{1}{2}}}{b^2 + y^2}.$$

$$F = \text{fl. } \frac{w^{\frac{1}{2}} \dot{w}}{w^2 - a^2}, \text{ to be found by theor. XIII.}$$

$$w = \sqrt{a^2 + x^2} = \sqrt{b^2 + y^2}.$$

NOTE. The *whole* fluent is infinite.

## T H E O R E M XXXVII.

$$\dot{F} = \frac{x^{\frac{1}{2}} \dot{x}}{a^2 + x^2} = \frac{\frac{1}{2}j}{b^2 + y^2}.$$

$$F = \frac{1}{2} w^{\frac{1}{2}} \pm \frac{1}{2} \text{fl. } \frac{w^{\frac{1}{2}} \dot{w}}{w^2 - 4a^2}, \text{ to be found by theor. XIV.}$$

$$w = \frac{a^2 + x^2}{x} = \frac{b^2 + y^2}{y^{\frac{1}{2}}}.$$

NOTE. The *whole* fluent is infinite.

THEO-

## T H E O R E M XXXVIII.

$$\dot{F} = \frac{x^{\frac{1}{2}} \dot{x}}{a^2 + x^2)^{\frac{1}{2}}} = \frac{\frac{1}{2} y^{\frac{1}{2}} \dot{y}}{b^2 + y^2)^{\frac{1}{2}}}$$

$$F = \frac{a^{\frac{1}{2}} w^{\frac{1}{2}}}{a^2 - w^2)^{\frac{1}{2}}} - \frac{1}{2} a^{\frac{1}{2}} \text{fl.} \frac{w^{\frac{1}{2}} \dot{w}}{a^2 - w^2)^{\frac{1}{2}}}, \text{ to be found by theor. VII.}$$

$$w = \frac{ax}{a^2 + x^2)^{\frac{1}{2}}} = \frac{b^{\frac{1}{2}} y^{\frac{1}{2}}}{b^2 + y^2)^{\frac{1}{2}}}$$

NOTE. The *whole* fluent is infinite.

---

## T H E O R E M XXXIX.

$$\dot{F} = \frac{x^{-\frac{1}{2}} \dot{x}}{a^2 - x^2)^{\frac{1}{2}}} = \frac{\frac{1}{2} y^{-\frac{1}{2}} \dot{y}}{b^2 - y^2)^{\frac{1}{2}}}$$

$$F = a^{-\frac{1}{2}} \text{fl.} \frac{w^{-\frac{1}{2}} \dot{w}}{a^2 + w^2)^{\frac{1}{2}}}, \text{ to be found by theor. xv.}$$

$$w = \frac{ax}{a^2 - x^2)^{\frac{1}{2}}} = \frac{b^{\frac{1}{2}} y^{\frac{1}{2}}}{b^2 - y^2)^{\frac{1}{2}}}$$

## T H E O R E M XL.

The *whole* fluent of  $\frac{x^{-\frac{1}{2}} \dot{x}}{a^2 - x^2)^{\frac{1}{2}}}$  is  $= \frac{2 \cdot 3^{\frac{1}{2}} a^{-1} P.}{3^{\frac{1}{2}} + 1}$ .

THEO.

T H E O R E M XLII.

$$\dot{F} = \frac{x^{-\frac{1}{2}} \dot{x}}{a^2 - x^2)^{\frac{1}{2}}} = \frac{\frac{1}{2} y \dot{y}}{b^2 - y^2)^{\frac{1}{2}}}$$

F = - fl.  $\frac{w^{-\frac{1}{2}} \dot{w}}{4a^2 + w^2)^{\frac{1}{2}}}$ , to be found by theor. xv.

$$w = \frac{a^2 - x^2}{x} = \frac{b^2 - y^2}{y^{\frac{1}{2}}}$$

T H E O R E M XLIII.

The whole fluent of  $\frac{x^{-\frac{1}{2}} \dot{x}}{a^2 - x^2)^{\frac{1}{2}}}$  is =  $\frac{2^{\frac{1}{2}} 3^{\frac{1}{2}} a^{-\frac{1}{2}} P}{3^{\frac{1}{2}} + 1}$ .

---

T H E O R E M XLIII.

$$\dot{F} = \frac{\dot{x}}{a^2 - x^2)^{\frac{1}{2}}} = \frac{\frac{1}{2} y^{\frac{1}{2}} \dot{y}}{b^2 - y^2)^{\frac{1}{2}}}$$

F = - fl.  $\frac{w^{-\frac{1}{2}} \dot{w}}{a^2 - w^2)^{\frac{1}{2}}}$ , to be found by theor. III.

$$w = \sqrt{a^2 - x^2} = \sqrt{b^2 - y^2}$$

T H E O R E M XLIV.

The whole fluent of  $\frac{\dot{x}}{a^2 - x^2)^{\frac{1}{2}}}$  is =  $\frac{3^{\frac{1}{2}} a^{-\frac{1}{2}} P}{3^{\frac{1}{2}} + 1}$ .

T H E O -

## T H E O R E M XLV.

$$\dot{F} = \frac{x^{\frac{1}{2}} \dot{x}}{a^2 - x^2)^{\frac{1}{2}}} = \frac{\frac{1}{2} y^{\frac{1}{2}} \dot{y}}{b^2 - y^2)^{\frac{1}{2}}}.$$

$$F = -\frac{1}{2} w^{\frac{1}{2}} + \frac{1}{2} \text{fl.} \frac{w^{\frac{1}{2}} \dot{w}}{4a^2 + w^2)^{\frac{1}{2}}}, \text{ to be found by theor. XIX.}$$

$$w = \frac{a^2 - x^2}{x} = \frac{b^2 - y^2}{y^{\frac{1}{2}}}.$$

## T H E O R E M XLVI.

The whole fluent of  $\frac{x^{\frac{1}{2}} \dot{x}}{a^2 - x^2)^{\frac{1}{2}}}$  is  $= 2^{\frac{1}{2}} 3^{\frac{1}{2}} a^{\frac{1}{2}} Q.$

## T H E O R E M XLVII.

$$\dot{F} = \frac{x^{-\frac{1}{2}} \dot{x}}{x^2 - a^2)^{\frac{1}{2}}} = \frac{\frac{1}{2} y^{-\frac{1}{2}} \dot{y}}{y^2 - b^2)^{\frac{1}{2}}}.$$

$$F = -a^{-\frac{1}{2}} \text{fl.} \frac{w^{-\frac{1}{2}} \dot{w}}{w^2 - a^2)^{\frac{1}{2}}}, \text{ to be found by theor. IX.}$$

$$w = \frac{ax}{x^2 - a^2)^{\frac{1}{2}}} = \frac{b^{\frac{1}{2}} y^{\frac{1}{2}}}{y^2 - b^2)^{\frac{1}{2}}}.$$

## T H E O R E M XLVIII.

The whole fluent of  $\frac{x^{-\frac{1}{2}} \dot{x}}{x^2 - a^2)^{\frac{1}{2}}}$  is  $= \frac{3^{\frac{1}{2}} a^{-1} P}{3^{\frac{1}{2}} + 1}.$

THEO.

T H E O R E M XLIX.

$$\dot{F} = \frac{x^{-\frac{1}{2}} \dot{x}}{x^2 - a^2)^{\frac{1}{2}}} = \frac{\frac{1}{2} \dot{y}}{y^2 - b^2)^{\frac{1}{2}}}.$$

F = fl.  $\frac{w^{-\frac{1}{2}} \dot{w}}{4a^2 + w^2)^{\frac{1}{2}}}$ , to be found by theorem xv.

$$w = \frac{x^2 - a^2}{x} = \frac{y^2 - b^2}{y^{\frac{1}{2}}}.$$

T H E O R E M L.

The *whole* fluent of  $\frac{x^{-\frac{1}{2}} \dot{x}}{x^2 - a^2)^{\frac{1}{2}}}$  is  $= \frac{2^{\frac{1}{2}} 3^{\frac{1}{2}} a^{-\frac{1}{2}} P}{3^{\frac{1}{2}} + 1}$ .

---

T H E O R E M LI.

$$\dot{F} = \frac{\dot{x}}{x^2 - a^2)^{\frac{1}{2}}} = \frac{\frac{1}{2} y^{\frac{1}{2}} \dot{y}}{y^2 - b^2)^{\frac{1}{2}}}.$$

F = fl.  $\frac{w^{-\frac{1}{2}} \dot{w}}{a^2 + w^2)^{\frac{1}{2}}}$ , to be found by theorem xvii.

$$w = \sqrt{x^2 - a^2} = \sqrt{y^2 - b^2}.$$

T H E O R E M LII.

The *whole* fluent of  $\frac{\dot{x}}{x^2 - a^2)^{\frac{1}{2}}}$  is  $= \frac{2 \cdot 3^{\frac{1}{2}} a^{-\frac{1}{2}} P}{3^{\frac{1}{2}} + 1}$ .



## T H E O R E M LIII.

$$\dot{F} = \frac{x^{\frac{3}{2}} \dot{x}}{x^2 - a^2)^{\frac{3}{2}}} = \frac{\frac{3}{2} y^{\frac{3}{2}} \dot{y}}{y^2 - b^2)^{\frac{3}{2}}}$$

$$F = \frac{1}{2} w^{\frac{1}{2}} + \frac{1}{2} \text{fl.} \frac{w^{\frac{1}{2}} \dot{w}}{4a^2 + w^2)^{\frac{3}{2}}}, \text{ to be found by theor. XIX.}$$

$$w = \frac{x^2 - a^2}{x} = \frac{y^2 - b^2}{y^{\frac{1}{2}}}$$

NOTE. The *whole* fluent is infinite.

## T H E O R E M LIV.

$$\dot{F} = \frac{x^{-\frac{3}{2}} \dot{x}}{a^2 + x^2)^{\frac{3}{2}}} = \frac{\frac{3}{2} y^{-\frac{3}{2}} \dot{y}}{b^2 + y^2)^{\frac{3}{2}}}$$

$$F = a^{-\frac{1}{2}} \text{fl.} \frac{w^{-\frac{3}{2}} \dot{w}}{a^2 - w^2)^{\frac{3}{2}}}, \text{ to be found by theor. I.}$$

$$w = \frac{ax}{a^2 + x^2)^{\frac{1}{2}}} = \frac{b^{\frac{3}{2}} y^{\frac{1}{2}}}{b^2 + y^2)^{\frac{1}{2}}}$$

## T H E O R E M LV.

$$\text{The whole fluent of } \frac{x^{-\frac{3}{2}} \dot{x}}{a^2 + x^2)^{\frac{3}{2}}} \text{ is } = \frac{3^{\frac{1}{2}} a^{-1} P}{3^{\frac{1}{2}} + 1}$$

THEO-

## T H E O R E M LVI.

$$\dot{F} = \frac{x^{-\frac{1}{2}} \dot{x}}{a^2 + x^2)^{\frac{3}{2}}} = \frac{\frac{3}{2}j}{(b^2 + y^2)^{\frac{3}{2}}}.$$

F = ± fl.  $\frac{w^{-\frac{3}{2}} \dot{w}}{w^2 - 4a^2)^{\frac{1}{2}}}$ , to be found by theorem IX.

$$w = \frac{a^2 + x^2}{x} = \frac{b^2 + y^2}{y^{\frac{1}{2}}}.$$

## T H E O R E M LVII.

The *whole* fluent of  $\frac{x^{-\frac{1}{2}} \dot{x}}{a^2 + x^2)^{\frac{3}{2}}}$  is =  $\frac{2^{\frac{1}{2}} 3^{\frac{1}{2}} a^{-\frac{1}{2}} P}{3^{\frac{1}{2}} + 1}$ .

## T H E O R E M LVIII.

$$\dot{F} = \frac{\dot{x}}{a^2 + x^2)^{\frac{3}{2}}} = \frac{\frac{1}{2} y^{\frac{1}{2}} \dot{y}}{(b^2 + y^2)^{\frac{3}{2}}}.$$

F = fl.  $\frac{w^{-\frac{1}{2}} \dot{w}}{w^2 - a^2)^{\frac{1}{2}}}$ , to be found by theorem XI.

$$w = \overline{a^2 + x^2)^{\frac{1}{2}}} = \overline{b^2 + y^2)^{\frac{1}{2}}}.$$

## T H E O R E M LIX.

The *whole* fluent of  $\frac{\dot{x}}{a^2 + x^2)^{\frac{3}{2}}}$  is =  $\frac{3^{\frac{1}{2}} a^{-\frac{1}{2}} P}{3^{\frac{1}{2}} + 1}$ .

## T H E O R E M LX.

$$\dot{F} = \frac{x^{\frac{1}{2}} \dot{x}}{a^2 + x^2]^{\frac{1}{2}}} = \frac{\frac{1}{2} y^{\frac{1}{2}} \dot{y}}{b^2 + y^2]^{\frac{1}{2}}}.$$

$F = \frac{1}{2} w^{\frac{1}{2}} \pm \frac{1}{2} \text{fl.} \frac{w^{\frac{1}{2}} \dot{w}}{w^2 - 4a^2]^{\frac{1}{2}}}$ , to be found by theorem XIII.

$$w = \frac{a^2 + x^2}{x} = \frac{b^2 + y^2}{y^{\frac{1}{2}}}.$$

NOTE. The *whole* fluent is infinite.

EXPLA-

## EXPLANATION OF TABLE IV.

A denotes the fluent of  $\frac{x^{\frac{1}{2}}z}{\sqrt{3+3x+z^2}}$ , B the fluent of  $\frac{x^{-\frac{1}{2}}z}{\sqrt{3+3x+z^2}}$ ,

C . . . . . the fluent of  $\frac{x^{\frac{1}{2}}z}{\sqrt{3-3x+z^2}}$ , D the fluent of  $\frac{x^{-\frac{1}{2}}z}{\sqrt{3-3x+z^2}}$ ;

whereof the general values are assigned in TABLE XII. by means of the arcs of the conic sections.

$$P \text{ is } = p + \sqrt{p^2 - 3^{\frac{1}{2}} + 1.d} = 4.366205,$$

$$Q = p - \sqrt{p^2 - 3^{\frac{1}{2}} + 1.d} = .982889;$$

$d$  being ( $= 1.570796$ ) the quadrantal arc of a circle whose radius is 1,

$p$  . . . . . ( $= 2.674547$ ) the quadrantal arc of an ellipsis whose semi-axes are  $3^{\frac{1}{2}}$  and  $\frac{3^{\frac{1}{2}} \cdot 3^{\frac{1}{2}} + 1}{2^{\frac{1}{2}}}$ .

The *whole* fluent (mentioned in any theorem) is generated whilst  $x$  from 0 becomes equal to  $a$ , whilst  $x$  from  $a$  becomes infinite, or whilst  $x$  from 0 becomes infinite; according as the denominator of the respective fluxional expression is some power or root of  $a^2 - x^2$ ,  $x^2 - a^2$ , or  $a^2 + x^2$ .

TABLE

---

T A B L E V.

CONTAINING

T H E O R E M S

FOR THE

CALCULATION of FLUENTS.

---

T H E O R E M I.

$$\dot{F} = \frac{x^{m-1} \dot{x}}{a^n + x^n}$$

$$F = \frac{2a^{m-n}}{n} \times \left\{ \begin{array}{l} \text{fl. } \frac{M' b \dot{x}}{b^2 + z^2} + \text{fl. } \frac{M'' c \dot{x}}{c^2 + z^2} + \text{fl. } \frac{M''' d \dot{x}}{d^2 + z^2} \text{ \&c.} \\ - \text{fl. } \frac{N' z \dot{x}}{b^2 + z^2} - \text{fl. } \frac{N'' z \dot{x}}{c^2 + z^2} - \text{fl. } \frac{N''' z \dot{x}}{d^2 + z^2} \text{ \&c.} \end{array} \right.$$

$m$  any positive integer less than the integer  $n$ .

$$z = \frac{x-a}{x+a}$$

$$\left. \begin{array}{l} M' \text{ and } N' \\ M'' \text{ and } N'' \\ M''' \text{ and } N''' \\ \text{\&c.} \end{array} \right\} \text{ sine and cosine of } \left\{ \begin{array}{l} \frac{m \times 180^\circ}{n} \\ \frac{3m \times 180^\circ}{n} \\ \frac{5m \times 180^\circ}{n} \\ \text{\&c.} \end{array} \right.$$

$$b = \text{tang. of } \frac{90^\circ}{n}, \quad c = \text{tang. of } \frac{3 \times 90^\circ}{n}, \quad d = \text{tang. of } \frac{5 \times 90^\circ}{n}, \text{ \&c.}$$

so long as these arcs are less than  $90^\circ$ .

Radius = 1.

THEO-

T H E O R E M II.

The fluent of  $\frac{x^{m-1}x}{a^n+x^n}$ , generated whilst  $x$  from being equal to any quantity  $k$  becomes equal to  $\frac{a^2}{k}$ , is

$$= \frac{a^{m-n}}{n} \times \left\{ \begin{array}{l} M' \times \text{Circ. Arc, rad. 1, tang. } \frac{g}{b} \\ + M'' \times \text{Circ. Arc, rad. 1, tang. } \frac{g}{c} \\ + M''' \times \text{Circ. Arc, rad. 1, tang. } \frac{g}{d} \\ \&c. \end{array} \right.$$

$m, n, M', M'', \&c. b, c, \&c.$  as in the preceding theor.  $g = \frac{a-k}{a+k}$ .

T H E O R E M III.

The fluent of  $\frac{x^{m-1}x}{a^n+x^n}$ , generated whilst  $x$  from being equal to any quantity  $k$  becomes equal to  $a$ , — the fluent of the same fluxion, generated whilst  $x$  from  $a$  becomes equal to  $\frac{a^2}{k}$ , is

$$= \frac{a^{m-n}}{n} \times N' \text{ Log. } \frac{b^2+g^2}{b^2} + N'' \text{ Log. } \frac{c^2+g^2}{c^2} \&c.$$

$m, n, N', N'', \&c. b, c, \&c.$  as in theor. I.  $g = \frac{a-k}{a+k}$ .

T H E O-

## T H E O R E M IV.

The fluent of  $\frac{x^{m-1} \dot{x}}{a^n + x^n}$ , generated whilst  $x$  from 0 becomes equal to  $a$ , is

$$= \frac{a^{m-n}}{n} \times \frac{Q}{s} + N' \text{Log.} \frac{1+b^2}{b^2} + N'' \text{Log.} \frac{1+c^2}{c^2} \&c.$$

$[m, n, N', N'', \&c. b, c, \&c. \text{ as in theorem I.}$

$Q = \text{quadrantal arc of the circ. rad. 1. } s = \text{fine of } \frac{2mQ}{n}.$

## T H E O R E M V.

The *whole* fluent of  $\frac{x^{m-1} \dot{x}}{a^n + x^n}$ , generated whilst  $x$  from 0 becomes infinite, is  $= \frac{2a^{m-n}}{n} \times Q.$

$m$  any positive integer or fraction less than the integer or fraction  $n.$

$Q$  and  $s$  as in the preceding theorem.

## T H E O R E M VI.

$$\dot{F} = \frac{x^{m-1} \dot{x}}{1 + a^{2m-n} x^{n-2m}} \times \frac{x^{m-1} \dot{x}}{a^n + x^n}.$$

$$F = \frac{4a^{m-n}}{n} \times \text{fl.} \frac{M' b \dot{x}}{b^2 + x^2} + \text{fl.} \frac{M'' c \dot{x}}{c^2 + x^2} \&c.$$

$m, n, x, M', M'', \&c. b, c, \&c. \text{ as in theorem I.}$

T H E O-

T H E O R E M VII.

The fluent of  $\frac{1}{1 + a^{2m-n}x^{n-2m}} \times \frac{x^{m-1}x}{a^2 + x^2}$ , generated whilst  $x$  from 0 becomes equal to  $a$ , or whilst  $x$  from  $a$  becomes infinite, is  $= \frac{2a^{m-n}}{n} \times Q$ .

$m$  and  $n$  as in theorem v.  $Q$  and  $s$  as in theorem iv.

T H E O R E M VIII.

$$\dot{F} = \frac{1}{1 - a^{2m-n}x^{n-2m}} \times \frac{x^{m-1}x}{a^2 + x^2}$$

$$F = -\frac{4a^{m-n}}{n} \times \text{fl.} \frac{N'xz}{b^2 + z^2} + \text{fl.} \frac{N''xz}{c^2 + z^2} \&c.$$

$m, n, z, N', N'', \&c. b, c, \&c.$  as in theorem i.

T H E O R E M IX.

The fluent of  $\frac{1}{1 - a^{2m-n}x^{n-2m}} \times \frac{x^{m-1}x}{a^2 + x^2}$ , generated whilst  $x$  from 0 becomes equal to  $a$ , is

$$= \frac{2a^{m-n}}{n} \times N' \text{Log.} \frac{1+b^2}{b^2} + N'' \text{Log.} \frac{1+c^2}{c^2} \&c.$$

$m, n, N', N'', \&c. b, c, \&c.$  as in theorem i.



## T H E O R E M X.

$$\dot{F} = \overline{a+x}^{n-r} \times \frac{x^{m-1} \dot{x}}{a^n + x^n}$$

$$F = \frac{2^{n-r+1}}{na^{r-m}} \times \left\{ \begin{array}{l} \text{fl. } \frac{N'bz}{b^2+z^2} - \text{fl. } \frac{N''cz}{c^2+z^2} + \&c. \\ - \text{fl. } \frac{M'xz}{b^2+z^2} + \text{fl. } \frac{M''xz}{c^2+z^2} - \&c. \end{array} \right.$$

$r$  any positive integer not greater than the integer  $n$ .

$m$  any positive integer less than  $r$ .

$$z = \frac{x-a}{x+a}$$

$$M' = P' \times \text{fine of } \frac{r-2m}{n} \cdot 90^\circ, \quad N' = P' \times \text{cofine of } \frac{r-2m}{n} \cdot 90^\circ.$$

$$M'' = P'' \times \text{fine of } 3 \cdot \frac{r-2m}{n} \cdot 90^\circ, \quad N'' = P'' \times \text{cofine of } 3 \cdot \frac{r-2m}{n} \cdot 90^\circ.$$

$$M''' = P''' \times \text{fine of } 5 \cdot \frac{r-2m}{n} \cdot 90^\circ, \quad N''' = P''' \times \text{cofine of } 5 \cdot \frac{r-2m}{n} \cdot 90^\circ.$$

&c.

&c.

&c.

&c.

$$b = \text{tang. of } \frac{90^\circ}{n}, \quad c = \text{tang. of } \frac{3 \times 90^\circ}{n}, \quad d = \text{tang. of } \frac{5 \times 90^\circ}{n}, \quad \&c.$$

so long as these arcs are less than  $90^\circ$ .

Radius = 1.

$$P' = \overline{1+b^2}^{\frac{r-m}{2}}, \quad P'' = \overline{1+c^2}^{\frac{r-m}{2}}, \quad P''' = \overline{1+d^2}^{\frac{r-m}{2}}, \quad \&c.$$

T H E O

T H E O R E M XI.

$$F = \overline{a+x}^{n-r} \times \frac{x^{r-1}z}{a^n + x^n}$$

$$F = \frac{2^{n-r+1}}{na^2} \times \text{fl.} \frac{P'bz}{b^2+z^2} - \text{fl.} \frac{P''cz}{c^2+z^2} + \&c.$$

$n, z, b, c, \&c. P', P'', \&c.$  as in the preceding theorem.  
 $r$  any even positive number not greater than  $n$ .

T H E O R E M XII.

$$F = \overline{a+x}^{n-r} \times \frac{x^{r-1}z}{a^n + x^n}$$

$$F = \text{fl.} \frac{z}{1-z} + \frac{2^{n-r+1}}{na} \times \left\{ \begin{array}{l} \text{fl.} \frac{N'bz}{b^2+z^2} - \text{fl.} \frac{N''cz}{c^2+z^2} + \&c. \\ + \text{fl.} \frac{M'xz}{b^2+z^2} - \text{fl.} \frac{M''xz}{c^2+z^2} + \&c. \end{array} \right.$$

$r=0$ , or any positive integer not greater than the integer  $n$ .

$z, b, c, \&c. P', P'', \&c.$  as in theorem x.

$$M' = P' \times \text{fine of } \frac{r}{n}.90^\circ. \quad N' = P' \times \text{cofine of } \frac{r}{n}.90^\circ.$$

$$M'' = P'' \times \text{fine of } \frac{3r}{n}.90^\circ. \quad N'' = P'' \times \text{cofine of } \frac{3r}{n}.90^\circ.$$

$$M''' = P''' \times \text{fine of } \frac{5r}{n}.90^\circ. \quad N''' = P''' \times \text{cofine of } \frac{5r}{n}.90^\circ.$$

$\&c. \quad \quad \quad \&c. \quad \quad \quad \&c. \quad \quad \quad \&c.$

## T H E O R E M XIII.

$$\dot{F} = \frac{x^{m-1} \dot{x}}{a^n - x^n}$$

$$F = \frac{2a^{m-1}}{n} \times \left\{ \begin{array}{l} \text{fl. } \frac{M'c\dot{x}}{c^2 + z^2} + \text{fl. } \frac{M''d\dot{x}}{d^2 + z^2} \&c. \\ - \text{fl. } \frac{\dot{x}}{z} - \text{fl. } \frac{N'z\dot{x}}{c^2 + z^2} - \text{fl. } \frac{N''z\dot{x}}{d^2 + z^2} \&c. \end{array} \right.$$

$m$  any positive integer less than the integer  $n$ .

$$z = \frac{x-a}{x+a}$$

$$\left. \begin{array}{l} M' \text{ and } N' \\ M'' \text{ and } N'' \\ M''' \text{ and } N''' \\ \&c. \end{array} \right\} \text{ sine and cosine of } \left\{ \begin{array}{l} \frac{2m \times 180^\circ}{n} \\ \frac{4m \times 180^\circ}{n} \\ \frac{6m \times 180^\circ}{n} \\ \&c. \end{array} \right.$$

$$c = \text{tang. of } \frac{180^\circ}{n}, d = \text{tang. of } \frac{2 \times 180^\circ}{n}, e = \text{tang. of } \frac{3 \times 180^\circ}{n}, \&c.$$

so long as these arcs are less than  $90^\circ$ .

Radius = 1.

## T H E O R E M XIV.

$$\dot{F} = \frac{x^{m-1} \dot{x}}{1 - a^{2m-n} x^{2m-n}} \times \frac{x^{m-1} \dot{x}}{a^n - x^n}$$

$$F = \frac{4a^{m-n}}{n} \times \text{fl. } \frac{M'c\dot{x}}{c^2 + z^2} + \text{fl. } \frac{M''d\dot{x}}{d^2 + z^2} \&c.$$

$m, n, z, M', M'', \&c. c, d, \&c.$  as in the preceding theorem.

T H E O

T H E O R E M XV.

The fluent of  $\frac{1 - a^{2m-n} x^{n-2m}}{a^n - x^n} \times \frac{x^{m-1} dx}{a^n - x^n}$ , generated whilst  $x$  from 0 becomes equal to  $a$ , or whilst  $x$  from  $a$  becomes infinite, is  $= \frac{2a^{m-n}}{nt} \times Q$ .

$m$  and  $n$  as in theor. v.  $Q$  as in theor. iv.  $t = \text{tang. of } \frac{2mQ}{n}$ .

T H E O R E M XVI.

$$F = \frac{1 - a^{2m-n} x^{n-2m}}{a^n - x^n} \times \frac{x^{m-1} dx}{a^n - x^n}$$

$$F = -\frac{4a^{m-n}}{n} \times \text{fl. } \frac{\frac{1}{2}x}{x} + \text{fl. } \frac{N'xz}{c^2 + x^2} + \text{fl. } \frac{N''xz}{d^2 + x^2} \&c.$$

$m, n, z, N', N'', \&c. c, d, \&c.$  as in theorem XIII.

T H E O R E M

## T H E O R E M XVII.

$$\dot{F} = \sqrt{a+x}^{n-r} \times \frac{x^{m-1} \dot{x}}{a^n - x^n}$$

$$F = \frac{2^{n-r+1}}{na^{r-m}} \times \left\{ \begin{array}{l} + \text{fl. } \frac{M' c \dot{x}}{c^2 + z^2} - \text{fl. } \frac{M'' d \dot{x}}{d^2 + z^2} + \&c. \\ - \text{fl. } \frac{\frac{1}{2} \dot{x}}{z} + \text{fl. } \frac{N' z \dot{x}}{c^2 + z^2} - \text{fl. } \frac{N'' z \dot{x}}{d^2 + z^2} + \&c. \end{array} \right.$$

$r$  any positive integer not greater than the integer  $n$ .

$m$  any positive integer less than  $r$ .

$$z = \frac{x-a}{x+a}$$

$$M' = P' \times \text{fine of } \frac{r-2m}{n} \cdot 180^\circ. \quad N' = P' \times \text{cofine of } \frac{r-2m}{n} \cdot 180^\circ.$$

$$M'' = P'' \times \text{fine of } 2 \cdot \frac{r-2m}{n} \cdot 180^\circ. \quad N'' = P'' \times \text{cofine of } 2 \cdot \frac{r-2m}{n} \cdot 180^\circ.$$

$$M''' = P''' \times \text{fine of } 3 \cdot \frac{r-2m}{n} \cdot 180^\circ. \quad N''' = P''' \times \text{cofine of } 3 \cdot \frac{r-2m}{n} \cdot 180^\circ.$$

$$\&c. \quad \&c. \quad \&c. \quad \&c.$$

$$c = \text{tang. of } \frac{180^\circ}{n}, \quad d = \text{tang. of } \frac{2 \times 180^\circ}{n}, \quad e = \text{tang. of } \frac{3 \times 180^\circ}{n}, \quad \&c.$$

so long as these arcs are less than  $90^\circ$ .

Radius = 1.

$$P' = \sqrt{1+c^2}^{\frac{r-n}{2}}, \quad P'' = \sqrt{1+d^2}^{\frac{r-n}{2}}, \quad P''' = \sqrt{1+e^2}^{\frac{r-n}{2}}, \quad \&c.$$

T H E O-

T H E O R E M XVIII.

$$\dot{F} = \overline{a+x}^{n-r} \times \frac{x^{\frac{1}{2}r-1}z}{a^n-x^n}$$

$$F = -\frac{2^{n-r+1}}{na^{\frac{1}{2}r}} \times \text{fl.} \frac{\frac{1}{2}z}{x} - \text{fl.} \frac{P'xz}{c^2+z^2} + \text{fl.} \frac{P''xz}{d^2+z^2} - \&c.$$

$n, z, c, d, \&c. P', P'', \&c.$  as in the preceding theorem.

$r$  any even positive number not greater than  $n$ .

T H E O R E M XIX.

$$\dot{F} = \overline{a+x}^{n-r} \times \frac{x^{\frac{1}{2}r-1}z}{a^n-x^n}$$

$$F = \text{fl.} \frac{z}{x-1} - \frac{2^{n-r+1}}{n} \times \left\{ \begin{array}{l} + \text{fl.} \frac{M'cz}{c^2+z^2} - \text{fl.} \frac{M''dz}{d^2+z^2} + \&c. \\ \text{fl.} \frac{\frac{1}{2}z}{x} - \text{fl.} \frac{N'xz}{c^2+z^2} + \text{fl.} \frac{N''xz}{d^2+z^2} - \&c. \end{array} \right.$$

$r$  any positive integer not greater than the integer  $n$ .

$z, c, d, \&c. P', P'', \&c.$  as in theorem xvii.

$$M' = P' \times \text{fine of } \frac{r}{n}.180^\circ. \quad N' = P' \times \text{cofine of } \frac{r}{n}.180^\circ.$$

$$M'' = P'' \times \text{fine of } \frac{2r}{n}.180^\circ. \quad N'' = P'' \times \text{cofine of } \frac{2r}{n}.180^\circ.$$

$$M''' = P''' \times \text{fine of } \frac{3r}{n}.180^\circ. \quad N''' = P''' \times \text{cofine of } \frac{3r}{n}.180^\circ.$$

&c.

&c.

&c.

&c.

T H E O.

T A B L E V.

T H E O R E M XX.

$$\dot{F} = \frac{y^{-1}j}{(gb^r + ky^n)^{\frac{m}{n}}}$$

$F = -\frac{\pi a^{m-n}}{rb^{\frac{m}{n}}} \times \text{fl. } \frac{x^{m-1}k}{a^m - gx^n}$ , to be found by Theor. I. or XIII.

$$x = \frac{ab^{\frac{r}{n}}}{(gb^r + ky^n)^{\frac{1}{n}}}, \quad y = h \times \frac{a^m - gx^n}{k^{\frac{1}{n}} x^{\frac{r}{n}}}$$

T H E O R E M XXI.

The fluent of  $\frac{y^{-1}j}{y^r - b^r}$ , generated whilst  $y$  from being

equal to  $h$  becomes equal to  $z^{\frac{1}{r}}h$ , is

$$= \frac{1}{rb^{\frac{m}{n}}} \times \frac{Q}{s} - N' \text{Log. } \frac{1+b^s}{b^s} - N'' \text{Log. } \frac{1+c^s}{c^s} \&c.$$

$m, n, N', N'', \&c. b, c, \&c.$  as in Theor. I.  $Q$  and  $s$  as in Theor. IV.

T H E O R E M XXII.

The *whole* fluent of  $\frac{y^{-1}j}{y^r - b^r}$ , generated whilst  $y$  from

being equal to  $h$  becomes infinite, is  $= \frac{zQ}{rb^{\frac{m}{n}}}$ .

$m, n, Q,$  and  $s$  as in Theorem IV.

T H E O

T H E O R E M XXIII.

$$F = \frac{y^{m-1} j}{(gb^r + ky^r)^{\frac{m}{r}}}$$

$F = \frac{na^{m-n}}{r} \times \text{fl. } \frac{x^{m-1} z}{a^n - kx^n}$ , to be found by Theorem I. of XIII.

$$x = \frac{ay^{\frac{r}{n}}}{(gb^r + ky^r)^{\frac{1}{n}}}, \quad y = h \times \frac{g^{\frac{1}{r}} x^{\frac{n}{r}}}{(a^n - kx^n)^{\frac{1}{r}}}$$

T H E O R E M XXIV.

The fluent of  $\frac{y^{m-1} j}{(b^r - y^r)^{\frac{m}{r}}}$ , generated whilst  $y$  from 0 becomes equal to  $\frac{b}{2^{\frac{1}{r}}}$ , is

$$= \frac{1}{r} \times \frac{Q}{s} + N' \text{Log. } \frac{1+b^2}{b^2} + N'' \text{Log. } \frac{1+c^2}{c^2} \&c.$$

$m, n, N', N'', \&c. b, c, \&c.$  as in Theor. I.  $Q$  and  $s$  as in Theor. IV.

T H E O R E M XXV.

The whole fluent of  $\frac{y^{m-1} j}{(b^r - y^r)^{\frac{m}{r}}}$ , generated whilst  $y$  from 0 becomes equal to  $h$ , is  $= \frac{2Q}{rs}$

$m, n, Q,$  and  $s$  as in Theorem IV.



---

T A B L E VI.

CONTAINING

T H E O R E M S

FOR THE

CALCULATION of FLUENTS.

---

T H E O R E M I.

$$\dot{F} = \overline{a + bx^n}^p \times x^{rn-1} \dot{x}.$$

$p = 0$ , or any positive integer.

$$\begin{aligned} F &= K + \frac{a^p x^{rn}}{n} \times \frac{1}{r} + \frac{pbx^n}{r+1.a} + \frac{p.p-1.b^2x^{2n}}{2.r+2.a^2} + \frac{p.p-1.p-2.b^3x^{3n}}{2.3.r+3.a^3} (p+1) \\ &= K + \frac{b^p x^{pn}}{n} \times \frac{1}{s} + \frac{pa}{s-1.bx^n} + \frac{p.p-1.a^2}{2.s-2.b^2x^{2n}} + \frac{p.p-1.p-2.a^3}{2.3.s-3.b^3x^{3n}} (p+1) \\ &= K + \frac{x^{rn} z^p}{ns} \times 1 + \frac{pa}{s-1.z} + \frac{p.p-1.a^2}{s-1.s-2.z^2} + \frac{p.p-1.p-2.a^3}{s-1.s-2.s-3.z^3} (p+1) \\ &= K + \frac{x^{rn} z^p}{nr} \times 1 - \frac{pbx^n}{r+1.z} + \frac{p.p-1.b^2x^{2n}}{r+1.r+2.z^2} - \frac{p.p-1.p-2.b^3x^{3n}}{r+1.r+2.r+3.z^3} (p+1) \end{aligned}$$

$$s = p + r. \quad z = a + bx^n.$$

THEO-

T H E O R E M II.

$$\dot{F} = \overline{a + bx^n}^p \times x^{rn-1}x.$$

$r$  any positive integer.

$$\begin{aligned} F &= K + \frac{z^i}{nb^r} \times \frac{1}{s} - \frac{r-1.a}{s-1.z} + \frac{r-1.r-2.a^2}{2.s-2.z^2} - (r) \\ &= K + \frac{\overline{-a}^{r-1}z^{p+1}}{nb^r} \times \frac{1}{p+1} - \frac{r-1.z}{p+2.a} + \frac{r-1.r-2.z^2}{2.p+3.a^2} - (r) \\ &= K + \frac{x^{rn-n}z^{p+1}}{bn^s} \times 1 - \frac{r-1.a}{s-1.bx^n} + \frac{r-1.r-2.a^2}{s-1.s-2.b^2x^{2n}} - (r) \\ &= K + \frac{x^{rn-n}z^{p+1}}{p+1.bn} \times 1 - \frac{r-1.z}{p+2.bx^n} + \frac{r-1.r-2.z^2}{p+2.p+3.b^2x^{2n}} - (r). \end{aligned}$$

$$s = p + r. \quad z = a + bx^n.$$

T H E O R E M III.

$$\dot{F} = \overline{a + bx^n}^p \times x^{rn-1}x.$$

$p + r$  any negative integer.

$$\begin{aligned} F &= K + \frac{x^{rn}}{na^t z^r} \times \frac{1}{r} - \frac{t-1.bx^n}{r+1.z} + \frac{t-1.t-2.b^2x^{2n}}{2.r+2.z^2} - (t) \\ &= K - \frac{\overline{-b}^{t-1}z^{p+1}}{na^t x^{rn+n}} \times \frac{1}{p+1} - \frac{t-1.z}{p+2.bx^n} + \frac{t-1.t-2.z^2}{2.p+3.b^2x^{2n}} - (t) \\ &= K + \frac{x^{rn}z^{p+1}}{rna} \times 1 + \frac{t-1.bx^n}{r+1.a} + \frac{t-1.t-2.b^2x^{2n}}{r+1.r+2.a^2} + (t) \\ &= K - \frac{x^{rn}z^{p+1}}{p+1.na} \times 1 - \frac{t-1.z}{p+2.a} + \frac{t-1.t-2.z^2}{p+2.p+3.a^2} - (t). \end{aligned}$$

$$t = -p - r. \quad z = a + bx^n.$$

h 2

NOTE

NOTE 1. When the exponent of the power of  $x$  or  $.x$  in any term (taken with its prefixed factor) in these theorems is  $= 0$ , and the denominator of the same term is also  $= 0$ , the equation will be fitly adjusted by taking the value of such term according to the Note to Theorem I. II. or III. in TABLE I.  $K$  being, at the same time, taken of a proper value.

NOTE 2. For the fluent of  $\overline{a + bx^n}^p \times x^{r-1}x$ , when  $p$  and  $r$  are such numbers that neither of the series in the values of  $F$ , in the three theorems in this Table, terminates, take any one of the values in which the series may be found to converge; and such value of  $F$ , considering the series therein as continued *ad infinitum*, will express the said fluent.

TABLE

---

T A B L E VII.

CONTAINING

T H E O R E M S

FOR THE

CALCULATION of FLUENTS.

---

T H E O R E M I.

$$\begin{aligned} \dot{F} &= \overline{a + bx^n}^p \times x^{rn-1} \dot{x}. \\ F &= \frac{x^{rn} z^p}{nt} \times \left( 1 + \frac{pa}{s-1.x} + \frac{p.p-1.a^2}{s-1.s-2.x^2} + (v) \right. \\ &\quad \left. + \frac{p.p-1.p-2(v) \times a^v}{s.s-1.s-2(v)} \times \text{fl. } z^{p-v} x^{rn-1} \dot{x}. \right) \\ &= \frac{x^{rn} z^p}{nr} \times \left( 1 - \frac{pbx^n}{r+1.x} + \frac{p.p-1.b^2 x^{2n}}{r+1.r+2.x^2} - (v) \right. \\ &\quad \left. + \frac{p.p-1.p-2(v) \times b^v}{r.r+1.r+2(v)} \times \text{fl. } z^{p-v} x^{rn+vn-1} \dot{x}. \right) \\ &= -\frac{x^{rn-n} z^{p+1}}{bnt} \times \left( 1 + \frac{r-1.a}{t+1.bx^n} + \frac{r-1.r-2.a^2}{t+1.t+2.b^2 x^{2n}} + (v) \right. \\ &\quad \left. + \frac{r-1.r-2.r-3(v) \times a^v}{t.t+1.t+2(v) \times b^v} \times \text{fl. } z^p x^{rn-vn-1} \dot{x}. \right) \end{aligned}$$

=

$$\begin{aligned}
&= \frac{x^{rn-n} z^{p+1}}{p+1.bn} \times \frac{1 - \frac{r-1.z}{p+2.bx^n} + \frac{r-1.r-2.z^2}{p+2.p+3.b^2x^{2n}}}{(v)} \\
&\quad \pm \frac{r-1.r-2.r-3(v)}{p+1.p+2.p+3(v) \times b^v} \times \text{fl. } x^{p+v} x^{rn-vn-1} x \\
&= \frac{x^{rn} z^{p+1}}{anr} \times \frac{1 + \frac{t-1.bx^n}{r+1.a} + \frac{t-1.t-2.b^2x^{2n}}{r+1.r+2.a^2}}{(v)} \\
&\quad + \frac{t-1.t-2.t-3(v) \times b^v}{r.r+1.r+2(v) \times a^v} \times \text{fl. } x^p x^{rn+vn-1} x \\
&= -\frac{x^{rn} z^{p+1}}{p+1.an} \times \frac{1 + \frac{s+1.z}{p+2.a} + \frac{s+1.s+2.z^2}{p+2.p+3.a^2}}{(v)} \\
&\quad + \frac{s+1.s+2.s+3(v)}{p+1.p+2.p+3(v) \times a^v} \times \text{fl. } x^{p+v} x^{rn-1} x.
\end{aligned}$$

$$s = -t = p + r. \quad z = a + bx^n.$$

\* + or - according as  $v$  is even or odd.

NOTE. This theorem being derived from the theorems in the preceding Table, it may sometimes be necessary to adjust the equation in the manner directed in Note 1. at the end of that Table: and the following theorems being deduced from this, the like correction in some of them may sometimes be requisite.

THE O-

T H E O R E M II.

$$\dot{F} = \overline{a + bx^n}^p \times x^{rn-1} \dot{x}. \quad \dot{G} = \overline{a + bx^n}^p \times x^{rn+vn-1} \dot{x}.$$

$$G = \frac{r.r + 1.r + 2(v) \times a^v}{t-1.t-2.t-3(v) \times b^v} \times$$

$$\overline{\frac{x^{rn} z^{p+1}}{anr} \times I + \frac{t-1.bx^n}{r+1.a} + \frac{t-1.t-2.b^2x^{2n}}{r+1.r+2.a^2}} (v).$$

$$t = -p - r. \quad z = a + bx^n.$$

T H E O R E M III.

The whole fluent of  $\overline{a + bx^n}^p \times x^{rn+vn-1} \dot{x}$ ,

generated whilst  $x$  from  $\left\{ \begin{array}{l} 0 \text{ becomes } = \frac{-a}{b} \Big| \frac{1}{n}, \text{ (1)} \\ \frac{-a}{b} \Big| \frac{1}{n} \text{ becomes infinite, (2)} \\ 0 \text{ becomes infinite, (3)} \end{array} \right.$

$$\text{is } = \frac{r.r + 1.r + 2(v) \times a^v}{t-1.t-2.t-3(v) \times b^v} \times F':$$

- (1)  $a, n, p + 1$ , and  $r$  being positive;  $b$  negative.
- (2)  $b, n$ , and  $p + 1$  being positive;  $a$ , and  $p + r + v$  negative.
- (3)  $a, b, n$ , and  $r$  being positive;  $p + r + v$  negative.

$F' =$  the contemporary fluent of  $\overline{a + bx^n}^p \times x^{rn-1} \dot{x}$ .

T H E O-

## T H E O R E M IV.

$$\dot{F} = \overline{a + bx^n}^p \times x^{rn-1} \dot{x}. \quad \dot{G} = \overline{a + bx^n}^p \times x^{rn-vn-1} \dot{x}.$$

$$G = \frac{t.t + 1.t + 2(v) \times b^v}{r - 1.r - 2.r - 3(v) \times a^v} \times$$

$$\overline{\frac{x^{rn-n} z^{p+1}}{bnz}} \times 1 + \frac{r-1.a}{t+1.bx^n} + \frac{r-1.r-2.a^2}{t+1.t+2.b^2x^{2n}} (v).$$

$$t = -p - r. \quad z = a + bx^n.$$

## T H E O R E M V.

The whole fluent of  $\overline{a + bx^n}^p \times x^{rn-vn-1} \dot{x}$ ,

generated whilst  $x$  from  $\left\{ \begin{array}{l} \circ \text{ becomes } = \overline{\frac{a}{b}}^{\frac{1}{n}}, \text{ (1)} \\ \overline{\frac{a}{b}}^{\frac{1}{n}} \text{ becomes infinite, (2)} \\ \circ \text{ becomes infinite, (3)} \end{array} \right.$

$$\text{is } = \frac{t.t + 1.t + 2(v) \times b^v}{r - 1.r - 2.r - 3(v) \times a^v} \times F' :$$

(1)  $a, n, p + 1$ , and  $r - v$  being positive;  $b$  negative.

(2)  $b, n$ , and  $p + 1$  being positive;  $a$ , and  $p + r$  negative.

(3)  $a, b, n$ , and  $r - v$  being positive;  $p + r$  negative.

$F' =$  the contemporary fluent of  $\overline{a + bx^n}^p \times x^{rn-1} \dot{x}$ .

T H E O-

T H E O R E M VI.

$$\dot{F} = \overline{a + bx^n}^p \times x^{rn-1} \dot{x}. \quad \dot{G} = \overline{a + bx^n}^{p+v} \times x^{rn-1} \dot{x}.$$

$$G = \frac{p + 1.p + 2.p + 3(v) \times a^v}{s + 1.s + 2.s + 3(v)} \times$$

$$\overline{F + \frac{x^{rn} z^{p+1}}{p + 1.an} \times 1 + \frac{s + 1.z}{p + 2.a} + \frac{s + 1.s + 2.z^2}{p + 2.p + 3.a^2} (v)}.$$

$$s = p + r. \quad z = a + bx^n.$$

T H E O R E M VII.

The *whole* fluent of  $\overline{a + bx^n}^{p+v} \times x^{rn-1} \dot{x}$ ,

generated whilst  $x$  from  $\left\{ \begin{array}{l} 0 \text{ becomes } = \frac{a}{b} \frac{1}{n}, \text{ (1)} \\ \frac{a}{b} \frac{1}{n} \text{ becomes infinite, (2)} \\ 0 \text{ becomes infinite, (3)} \end{array} \right.$

$$\text{is } = \frac{p + 1.p + 2.p + 3(v) \times a^v}{s + 1.s + 2.s + 3(v)} \times F' :$$

(1)  $a, n, p + 1$ , and  $r$  being positive;  $b$  negative.

(2)  $b, n$ , and  $p + 1$  being positive;  $a$ , and  $p + r + v$  negative.

(3)  $a, b, n$ , and  $r$  being positive;  $p + r + v$  negative.

$F'$  = the contemporary fluent of  $\overline{a + bx^n}^p \times x^{rn-1} \dot{x}$ .



## T H E O R E M VIII.

$$\dot{F} = \overline{a + bx^n}^p \times x^{r-1} \dot{x}. \quad \dot{G} = \overline{a + bx^n}^{p-v} \times x^{r-1} \dot{x}.$$

$$G = \frac{s.s - 1.s - 2(v)}{p.p - 1.p - 2(v) \times a^v} \times$$

$$F = \frac{x^r z^p}{ns} \times 1 + \frac{pa}{s-1.z} + \frac{p.p-1.a^2}{s-1.s-2.z^2} (v).$$

$$s = p + r. \quad z = a + bx^n.$$

## T H E O R E M IX.

The *whole* fluent of  $\overline{a + bx^n}^{p-v} \times x^{r-1} \dot{x}$ ,

generated whilst  $x$  from  $\begin{cases} 0 \text{ becomes } = \frac{\overline{a}}{b}^{\frac{1}{n}}, & (1) \\ \frac{\overline{a}}{b}^{\frac{1}{n}} \text{ becomes infinite,} & (2) \\ 0 \text{ becomes infinite,} & (3) \end{cases}$

$$\text{is } = \frac{s.s - 1.s - 2(v)}{p.p - 1.p - 2(v) \times a^v} \times F'$$

(1)  $a, n, p - v + 1$ , and  $r$  being positive;  $b$  negative.

(2)  $b, n$ , and  $p - v + 1$  being positive;  $a$ , and  $p + r$  negative.

(3)  $a, b, n$ , and  $r$  being positive;  $p + r$  negative.

$F'$  = the contemporary fluent of  $\overline{a + bx^n}^p \times x^{r-1} \dot{x}$ .

T H E O-

T H E O R E M X.

$$F = \sqrt[p]{a + bx^n} \times x^{r-1} \dot{x}. \quad G = \sqrt[p-v]{a + bx^n} \times x^{r+v-1} \dot{x}.$$

$$G = \pm \frac{r.r + 1.r + 2(v)}{p.p - 1.p - 2(v) \times b^v} \times$$

$$F = \frac{x^{rn} z^p}{nr} \times 1 - \frac{pbx^n}{r+1.z} + \frac{p.p - 1.b^2 x^{2n}}{r+1.r+2.z^2} - (v)$$

$$z = a + bx^n.$$

T H E O R E M XI.

The whole fluent of  $\sqrt[p-v]{a + bx^n} \times x^{r+v-1} \dot{x}$ ,

generated whilst  $x$  from  $\begin{cases} a \text{ becomes } = \frac{a}{b} \frac{1}{n}, & (1) \\ \frac{a}{b} \frac{1}{n} \text{ becomes infinite,} & (2) \\ a \text{ becomes infinite,} & (3) \end{cases}$

$$\text{is } = \pm \frac{r.r + 1.r + 2(v)}{p.p - 1.p - 2(v) \times b^v} \times F^v:$$

(1)  $a$ ,  $p - v + 1$ , and  $r$  being positive;  $b$  negative.

(2)  $b$ ,  $n$ , and  $p - v + 1$  being positive;  $a$ , and  $p + r$  negative.

(3)  $a$ ,  $b$ ,  $n$ , and  $r$  being positive;  $p + r$  negative.

$F^v =$  the contemporary fluent of  $\sqrt[p]{a + bx^n} \times x^{r-1} \dot{x}$ .

\* + or - according as  $v$  is even or odd.

## T H E O R E M XII.

$$\dot{F} = \overline{a + bx^n}^p \times x^{r-n-1} \dot{x}. \quad \dot{G} = \overline{a + bx^n}^{p+v} \times x^{r-n-v-1} \dot{x}.$$

$$G = * \pm \frac{p+1.p+2.p+3(v) \times b^v}{r-1.r-2.r-3(v)} \times$$

$$\frac{x^{r-n-n} z^{p+1}}{p+1.bn} \times \left( 1 - \frac{r-1.z}{p+2.bx^n} + \frac{r-1.r-2.z^2}{p+2.p+3.b^2 x^{2n}} - (v) \right).$$

$$z = a + bx^n.$$

## T H E O R E M XIII.

The *whole* fluent of  $\overline{a + bx^n}^{p+v} \times x^{r-n-v-1} \dot{x}$ ,

generated whilst  $x$  from  $\begin{cases} 0 \text{ becomes } = \left| \frac{-a}{b} \right|^{\frac{1}{n}}, & (1) \\ \left| \frac{-a}{b} \right|^{\frac{1}{n}} \text{ becomes infinite,} & (2) \\ 0 \text{ becomes infinite,} & (3) \end{cases}$

$$\text{is } = * \pm \frac{p+1.p+2.p+3(v) \times b^v}{r-1.r-2.r-3(v)} \times F' :$$

(1)  $a, n, p+1$ , and  $r-v$  being positive;  $b$  negative.

(2)  $b, n$ , and  $p+1$  being positive;  $a$ , and  $p+r$  negative.

(3)  $a, b, n$ , and  $r-v$  being positive;  $p+r$  negative.

$F'$  = the contemporary fluent of  $\overline{a + bx^n}^p \times x^{r-n-1} \dot{x}$ .

\* + or - according as  $v$  is even or odd.

T H E O-

T H E O R E M XIV.

$$\begin{aligned} \dot{F} &= \overline{a + bx^n}^p \times x^{rn-1} \dot{x}. & \dot{H} &= \overline{a + bx^n}^{p+w} \times x^{rn+vn-1} \dot{x}. \\ H &= \frac{p+1 \cdot p+2(w) \times a^w}{s+1 \cdot s+2(w)} \times \frac{x^{rn+vn} z^{p+1}}{p+1 \cdot an} \times \left( 1 + \frac{s+1 \cdot z}{p+2 \cdot a} + \frac{s+1 \cdot s+2 \cdot z^2}{p+2 \cdot p+3 \cdot a^2} (w) \right) \\ & * \pm \frac{r \cdot r+1(v) \times p+1 \cdot p+2(w) \times a^{v+w}}{p+r+1 \cdot p+r+2(v+w) \times b^v} \times \\ & \overline{\frac{x^{rn} z^{p+1}}{anr} \times \left( 1 + \frac{t-1 \cdot bx^n}{r+1 \cdot a} + \frac{t-1 \cdot t-2 \cdot b^2 x^{2n}}{r+1 \cdot r+2 \cdot a^2} (v) \right)}. \\ s &= p+r+v. & t &= -p-r. & z &= a+bx^n. \end{aligned}$$

T H E O R E M XV.

The whole fluent of  $\overline{a + bx^n}^{p+w} \times x^{rn+vn-1} \dot{x}$ ,

generated whilst  $x$  from  $\begin{cases} 0 \text{ becomes } = \frac{-a}{b} \frac{1}{n}, & (1) \\ \frac{-a}{b} \frac{1}{n} \text{ becomes infinite,} & (2) \\ 0 \text{ becomes infinite,} & (3) \end{cases}$

$$\text{is } = * \pm \frac{r \cdot r+1(v) \times p+1 \cdot p+2(w) \times a^{v+w}}{p+r+1 \cdot p+r+2(v+w) \times b^v} \times F'$$

- (1)  $a, n, p+1$ , and  $r$  being positive;  $b$  negative.
- (2)  $b, n$ , and  $p+1$  being positive;  $a$ , and  $p+r+v+w$  negative.
- (3)  $a, b, n$ , and  $r$  being positive;  $p+r+v+w$  negative.

$F' =$  the contemporary fluent of  $\overline{a + bx^n}^p \times x^{rn-1} \dot{x}$ .

\* + or - according as  $v$  is even or odd.

T H E O-

## T H E O R E M XVI.

$$\begin{aligned} \dot{F} &= \overline{a + bx^n}^p \times x^{r-1} \dot{x}. & \dot{H} &= \overline{a + bx^n}^{p-w} \times x^{r+v-1} \dot{x}. \\ H &= \frac{s.s-1(w)}{p.p-1(w) \times a^w} \times \frac{x^{rn+vn} z^p}{ns} \times \left( 1 + \frac{pa}{s-1.z} + \frac{p.p-1.a^2}{s-1.s-2.z^2} (w) \right) \\ &+ \frac{r.r+1(v) \times s.s-1(w) \times a^{v-w}}{t-1.t-2(v) \times p.p-1(w) \times b^v} \times \\ &\overline{\frac{x^{rn} z^{p+1}}{anr} \times \left( 1 + \frac{t-1.bx^n}{r+1.a} + \frac{t-1.t-2.b^2 x^{2n}}{r+1.r+2.a^2} (v) \right)}. \\ s &= p+r+v. & t &= -p-r. & z &= a + bx^n. \end{aligned}$$

## T H E O R E M XVII.

The whole fluent of  $\overline{a + bx^n}^{p-w} \times x^{r+v-1} \dot{x}$ ,

generated whilst  $x$  from  $\begin{cases} 0 \text{ becomes } = \frac{a}{b} \frac{1}{n}, & (1) \\ \frac{a}{b} \frac{1}{n} \text{ becomes infinite,} & (2) \\ 0 \text{ becomes infinite,} & (3) \end{cases}$

$$\text{is } = \frac{r.r+1(v) \times s.s-1(w) \times a^{v-w}}{t-1.t-2(v) \times p.p-1(w) \times b^v} \times F':$$

(1)  $a, n, p-w+1$ , and  $r$  being positive;  $b$  negative.

(2)  $b, n$ , and  $p-w+1$  being posit.;  $a, p+r$ , and  $p+r+v-w$  negat.

(3)  $a, b, n$ , and  $r$  being positive;  $p+r$ , and  $p+r+v-w$  negative.

$F'$  = the contemporary fluent of  $\overline{a + bx^n}^p \times x^{r-1} \dot{x}$ .

T H E O-

T H E O R E M XVIII.

$$\begin{aligned} \dot{F} &= \overline{a + bx^n}^p \times x^{rn-1} \dot{x}. & \dot{H} &= \overline{a + bx^n}^{p+w} \times x^{rn-vn-1} \dot{x}. \\ H &= \frac{p+1 \cdot p+2(w) \times a^w}{s+1 \cdot s+1(w)} \times \frac{x^{rn-vn} z^{p+1}}{p+1 \cdot an} \times \left( 1 + \frac{s+1 \cdot z}{p+2 \cdot a} + \frac{s+1 \cdot r+2 \cdot z^2}{p+2 \cdot p+3 \cdot a^2} (w) \right) \\ &+ \frac{t \cdot t+1(v) \times p+1 \cdot p+2(w) \times b^v}{r-1 \cdot r-2(v) \times s+1 \cdot s+2(w) \times a^{v-w}} \times \end{aligned}$$

$$\overline{F} + \frac{x^{rn-a} z^{p+1}}{bnz} \times \left( 1 + \frac{r-1 \cdot a}{s+1 \cdot bx^n} + \frac{r-1 \cdot r-2 \cdot a^2}{s+1 \cdot s+2 \cdot b^2 x^{2n}} (v) \right).$$

$$s = p + r - v. \quad t = -p - r. \quad z = a + bx^n.$$

T H E O R E M XIX.

The whole fluent of  $\overline{a + bx^n}^{p+w} \times x^{rn-vn-1} \dot{x}$ ,

generated whilst  $x$  from  $\begin{cases} \circ \text{ becomes } = \frac{a}{b} \frac{1}{n}, & (1) \\ \frac{a}{b} \frac{1}{n} \text{ becomes infinite,} & (2) \\ \circ \text{ becomes infinite,} & (3) \end{cases}$

$$\text{is } = \frac{t \cdot t+1(v) \times p+1 \cdot p+2(w) \times b^v}{r-1 \cdot r-2(v) \times s+1 \cdot s+2(w) \times a^{v-w}} \times F';$$

- (1)  $a, n, p+1$ , and  $r-v$  being positive;  $b$  negative.
- (2)  $b, n$ , and  $p+1$  being positive;  $a, p+r$ , and  $p+r-v+w$  negative.
- (3)  $a, b, n$ , and  $r-u$  being positive;  $p+r$ , and  $p+r-v+w$  negat-

$F'$  = the contemporary fluent of  $\overline{a + bx^n}^p \times x^{rn-1} \dot{x}$ .

T H E O-

## T H E O R E M XX.

$$\begin{aligned} \dot{F} &= \overline{a + bx^n}^p \times x^{rn-1} \dot{x}. & \dot{H} &= \overline{a + bx^n}^{p-w} \times x^{rn-w-1} \dot{x}. \\ H &= \frac{s.s-1(w)}{p.p-1(w) \times a^w} \times -\frac{x^{rn-w} z^p}{ns} \times I + \frac{pa}{s-1.z} + \frac{p.p-1.a^2}{s-1.s-2.z^2} (w) \\ &+ \frac{s.s+1(v) \times s.s-1(w) \times b^v}{r-1.r-2(v) \times p.p-1(w) \times a^{v+w}} \times \\ &\overline{F + \frac{x^{rn-s} z^{p+1}}{bns} \times I + \frac{r-1.a}{t+1.bx^s} + \frac{r-1.r-2.a^2}{t+1.s+2.b^2 x^{2s}} (v)} \\ s &= p + r - v. & t &= -p - r. & z &= a + bx^n. \end{aligned}$$

## T H E O R E M XXI.

The whole fluent of  $\overline{a + bx^n}^{p-w} \times x^{rn-w-1} \dot{x}$ ,

generated whilst  $x$  from  $\left\{ \begin{array}{l} 0 \text{ becomes } = \frac{a}{b} \Big|^{1/n}, \text{ (1)} \\ \frac{a}{b} \Big|^{1/n} \text{ becomes infinite, (2)} \\ 0 \text{ becomes infinite, (3)} \end{array} \right.$

$$Is = \frac{s.s+1(v) \times s.s-1(w) \times b^v}{r-1.r-2(v) \times p.p-1(w) \times a^{v+w}} \times F'$$

(1)  $a, n, p - w + 1$ , and  $r - v$  being positive;  $b$  negative.

(2)  $b, n$ , and  $p - w + 1$  being positive;  $a$ , and  $p + r$  negative.

(3)  $a, b, n$ , and  $r - v$  being positive;  $p + r$  negative.

$F$  = the contemporary fluent of  $\overline{a + bx^n}^p \times x^{rn-1} \dot{x}$ .

T H E O-

T H E O R E M XXII.

The whole fluent of  $\overline{a-x^n}^p \times x^{n-1} \dot{x} \times \overline{P+Qx^n+Rx^{2n}} \&c.$

generated whilst  $x$  from 0 becomes equal to  $a^{\frac{1}{n}}$ ,

$$is = F \times P + \frac{Qar}{s+1} + \frac{Ra^2r.r+1}{s+1.s+2} + \frac{Sa^3r.r+1.r+2}{s+1.s+2.s+3} \&c.$$

$a, n, p+1$ , and  $r$  being positive;  $s = p+r$ ;

$F =$  the contemporary fluent of  $\overline{a-x^n}^p \times x^{n-1} \dot{x}$ .

T H E O R E M XXIII.

The whole fluent of  $\overline{a-x^n}^p \times x^{n-1} \dot{x} \times \overline{P+Q.a-x^n+R.a-x^n}^2 \&c.$

generated whilst  $x$  from 0 becomes equal to  $a^{\frac{1}{n}}$ ,

$$is = F \times P + \frac{Qa.p+1}{s+1} + \frac{Ra^2.p+1.p+2}{s+1.s+2} + \frac{Sa^3.p+1.p+2.p+3}{s+1.s+2.s+3} \&c.$$

$a, n, p+1, r, s$ , and  $F$  being as in the preceding theorem.

T H E O R E M XXIV.

The whole fluent of  $\overline{x^n-a}^p \times x^{n-1} \dot{x} \times \overline{P+Qx^{-n}+Rx^{-2n}} \&c.$

generated whilst  $x$  from  $a^{\frac{1}{n}}$  becomes infinite,

$$is = F \times P + \frac{Qs}{a.r-1} + \frac{Rs.s-i}{a^2.r-1.r-2} + \frac{Ss.s-1.s-2}{a^3.r-1.r-2.r-3} \&c.$$

$a, n$ , and  $p+1$  being positive;  $s (= p+r)$  negative;

$F =$  the contemporary fluent of  $\overline{x^n-a}^p \times x^{n-1} \dot{x}$ .



## T H E O R E M XXV.

The whole fluent of  $\overline{x^r - a}^p \times x^{n-1} \dot{x} \times P + Q \frac{x^r - a}{x^n} + R \frac{x^r - a}{x^n}$  &c.

generated whilst  $x$  from  $a^{\frac{1}{n}}$  becomes infinite,

$$is = F \times P - \frac{Q.p+1}{r-1} + \frac{R.p+1.p+2}{r-1.r-2} - \frac{S.p+1.p+2.p+3}{r-1.r-2.r-3} \&c.$$

$a, n, p+1, p+r,$  and  $F$  being as in the preceding theorem.

## T H E O R E M XXVI.

The whole fluent of  $\overline{a+x^n}^p \times x^{n-1} \dot{x} \times P + \frac{Q}{a+x^n} + \frac{R}{a+x^n}$  &c.

generated whilst  $x$  from  $0$  becomes infinite,

$$is = F \times P + \frac{Qs}{ap} + \frac{R.s-1}{a^2 p.p-1} + \frac{S.s-1.s-2}{a^3 p.p-1.p-2} \&c.$$

$a, n,$  and  $r$  being positive;  $s (= p+r)$  negative;

$F =$  the contemporary fluent of  $\overline{a+x^n}^p \times x^{n-1} \dot{x}$ .

## T H E O R E M XXVII.

The whole fluent of  $\overline{a+x^n}^p \times x^{n-1} \dot{x} \times P + \frac{Qx^n}{a+x^n} + \frac{Rx^{2n}}{a+x^n}$  &c.

generated whilst  $x$  from  $0$  becomes infinite,

$$is = F \times P - \frac{Qr}{p} + \frac{Rr.r+1}{p.p-1} - \frac{Sr.r+1.r+2}{p.p-1.p-2} \&c.$$

$a, n, r, p+r,$  and  $F$  being as in the preceding theorem.

T A B L E

---

T A B L E VIII

CONTAINING

T H E O R E M S

FOR THE

CALCULATION of FLUENTS.

---

T H E O R E M I.

$$\frac{1.2.3(r-1) \times p + q + 1.p + q + 2.p + q + 3(r-1)}{p + 1.p + 2.p + 3(r-1) \times q + 1.q + 2.q + 3(r-1)} \text{ is } =$$

$$F \times 1 + \frac{pq}{r} + \frac{p.p - 1 \times q.q - 1}{1.2 \times r.r + 1} + \frac{p.p - 1.p - 2 \times q.q - 1.q - 2}{1.2.3 \times r.r + 1.r + 2} \&c.$$

F being =  $nq \times$  the whole fluent of  $(1-x^n)^p \times x^{r-1} x^q$ .

T H E O R E M II.

$$\frac{p + r.p + r + 1.p + r + 2(q)}{r.r + 1.r + 2.(q)} \text{ is } =$$

$$1 + \frac{pq}{r} + \frac{p.p - 1 \times q.q - 1}{1.2 \times r.r + 1} + \frac{p.p - 1.p - 2 \times q.q - 1.q - 2}{1.2.3 \times r.r + 1.r + 2} \&c.$$

## THEOREM III.

$$F = \overline{a - x^n}^p \times x^{rn-1} z.$$

$$F = \left\{ \begin{array}{l} K + \frac{x^{rn} z^{p+1}}{anr} \times 1 + \frac{s + 1.x^n}{r + 1.a} + \frac{s + 1.s + 2.x^{2n}}{r + 1.r + 2.a^2} (v) \\ + \frac{s + 1.s + 2(v)}{r.r + 1(v) \times a^v} \times \frac{x^{r'n} z^p}{nr''} \times 1 + \frac{p x^n}{r'' + 1.z} + \frac{p.p - 1.x^{2n}}{r'' + 1.r'' + 2.z^2} (w) \\ + \frac{s + 1.s + 2.s + 3(v)}{r.r + 1.r + 2(v) \times a^v} \times \frac{p.p - 1.p - 2(w)}{r''.r'' + 1.r'' + 2(w)} \times \frac{x^{r''+w} z^{p-w+1}}{an.r'' + w} \\ \times 1 + \frac{s'' + 1.x^n}{r'' + 1.a} + \frac{s'' + 1.s'' + 2.x^{2n}}{r'' + 1.r'' + 2.a^2} \&c. \end{array} \right.$$

$$= \left\{ \begin{array}{l} K + \frac{x^{rn} z^{p+1}}{anr} \times 1 + \frac{s + 1.x^n}{r + 1.a} + \frac{s + 1.s + 2.x^{2n}}{r + 1.r + 2.a^2} (v) \\ + \frac{s + 1.s + 2(v)}{r.r + 1(v) \times a^v} \times \frac{x^{r'n} z^p}{nr''} \times 1 + \frac{p x^n}{r'' + 1.z} + \frac{p.p - 1.x^{2n}}{r'' + 1.r'' + 2.z^2} (w) \\ + \frac{s + 1.s + 2.s + 3(v)}{r.r + 1.r + 2(v) \times a^v} \times \frac{p.p - 1.p - 2(w)}{r''.r'' + 1.r'' + 2(w)} \times \frac{2x^{2s''}}{n.2s'' + 1.z^{2s''}} \\ \times y + \frac{y^3}{2s'' + 3} + \frac{3y^5}{2s'' + 3.2s'' + 5} + \frac{3.5y^7}{2s'' + 3.2s'' + 5.2s'' + 7} \&c. \end{array} \right.$$

$$r'' = r + v. \quad r''' = r'' + w. \quad s = p + r. \quad s'' = p + r''. \quad y = \frac{x^n}{2a - x^n}. \quad z = a - x^n.$$

$v$  and  $w$  any positive integers ;

so that  $w - v$ , in the second value of  $F$ , be  $= 2p + r + 1$ .

THE O.

T H E O R E M IV.

The *whole* fluent of  $\sqrt[p]{a - x^n} \times x^{r-n-1} x$ ,

generated whilst  $x$  from 0 becomes equal to  $a^{\frac{1}{n}}$ ,

$$\begin{aligned} \text{is} &= \frac{1.2.3(r)}{p+1.p+2.p+3(r)} \times \frac{a^{p+r}}{nr} = \frac{1.2.3(p)}{r+1.r+2.r+3(p)} \times \frac{a^{p+r}}{nr} \\ &= \left\{ \frac{1.2.3(r+v+y) \times 1.1+1.1+2(w)}{p+1.p+2.p+3(r+w+y) \times r+1.r+2.r+3(v)} \times \frac{a^{p+r}}{nr} \right. \\ &\quad \left. \times 1 + \frac{p-v+w \times y}{r+v+1} + \frac{p-v+w.p-v+w-1 \times y.y-1}{1.2 \times r+v+1.r+v+2} \&c. \right. \\ &= \left\{ \frac{1.2.3(p+v+y) \times 1.1+1.1+2(w)}{p+1.p+2.p+3(v) \times r+1.r+2.r+3(p+w+y)} \times \frac{a^{p+r}}{nr} \right. \\ &\quad \left. \times 1 + \frac{r-v+w \times y}{p+v+1} + \frac{r-v+w.r-v+w-1 \times y.y-1}{1.2 \times p+v+1.p+v+2} \&c. \right. \\ &= \text{the Limit of } \frac{1.2.3(z) \times 1.1+1.1+2(z)}{p+1.p+2.p+3(z) \times r+1.r+2.r+3(z)} \times \frac{a^{p+r}}{nr} \\ &\quad \text{z increasing ad infinitum.} \end{aligned}$$

$a, n, p+1,$  and  $r$  being positive.  $t = p+r+1,$

$p+y,$  or  $r+y,$  any positive integer.

$v$  and  $w$  any positive integers..

T H E O

## T H E O R E M V.

$$F = \overline{x^s - a}^p \times x^{rn-1} z$$

$$F = \left\{ \begin{array}{l} K + \frac{x^{rn} z^{p+a}}{an.p+1} \times 1 - \frac{s+1.z}{p+2.a} + \frac{s+1.s+2.z^2}{p+2.p+3.a^2} - (v) \\ \pm * \frac{s+1.s+2(v)}{p+1.p+2(v) \times a^v} \times \frac{x^{rn-n} z^{p'+1}}{n.p'+1} \times 1 - \frac{r-1.z}{p''+2.z^n} + \frac{r-1.r-2.z^2}{p''+2.p''+3.z^{2n}} \&c. \end{array} \right.$$

$$= \left\{ \begin{array}{l} K + \frac{x^{rn} z^{p+1}}{an.p+1} \times 1 - \frac{s+1.z}{p+2.a} + \frac{s+1.s+2.z^2}{p+2.p+3.a^2} - (v) \\ \pm * \frac{s+1.s+2(v)}{p+1.p+2(v) \times a^v} \times \frac{x^{rn-n} z^{p'+1}}{n.p'+1} \times 1 - \frac{r-1.z}{p''+2.z^n} + \frac{r-1.r-2.z^2}{p''+2.p''+3.z^{2n}} - (w) \\ \pm + \frac{s+1.s+2.s+3(v)}{p+1.p+2.p+3(v) \times a^v} \times \frac{r-1.r-2.r-3(w)}{p''+1.p''+2.p''+3(w)} \times \frac{2z^{2r}}{n.2s'+1.z^{2n}} \\ \times y + \frac{y^3}{2s'+3} + \frac{3y^5}{2s'+3.2s'+5} + \frac{3.5y^7}{2s'+3.2s'+5.2s'+7} \&c. \end{array} \right.$$

$$p'' = p + v. \quad s = p + r. \quad s'' = p'' + r. \quad y = \frac{x^r - a}{x^r + a}. \quad z = x^r - a.$$

$v$  and  $w$  any positive integers; so that  $w - v$  be  $= p + 2r$ .

\* + or - according as  $v$  is even or odd.

† + or - according as  $v + w$  is even or odd.

T H E O-

T H E O R E M VI.

The whole fluent of  $\sqrt[p]{x^2 - a^2} \times x^{-n-1} x^{\frac{1}{2}}$ ,

generated whilst  $x$  from  $a^{\frac{1}{2}}$  becomes infinite,

$$\begin{aligned} \text{is} &= \frac{1.2.3(t)}{p+1.p+2.p+3(t)} \times \frac{a^{p-r}}{nr} = \frac{1.2.3(p)}{t+1.t+2.t+3(p)} \times \frac{a^{p-r}}{nr} \\ &= \left\{ \frac{1.2.3(t+v+y) \times r.r+1.r+2(w)}{p+1.p+2.p+3(t+w+y) \times t+1.t+2.t+3(v)} \times \frac{a^{p-r}}{nr} \right. \\ &\quad \left. \times 1 + \frac{p-v+w \times y}{t+v+1} + \frac{p-v+w.p-v+w-1 \times y.y-1}{1.2 \times t+v+1.t+v+2} \&c. \right. \\ &= \left\{ \frac{1.2.3(p+v+y) \times r.r+1.r+2(w)}{p+1.p+2.p+3(v) \times t+1.t+2.t+3(p+w+y)} \times \frac{a^{p-r}}{nr} \right. \\ &\quad \left. \times 1 + \frac{t-v+w \times y}{p+v+1} + \frac{t-v+w.t-v+w-1 \times y.y-1}{1.2 \times p+v+1.p+v+2} \&c. \right. \\ &= \text{the Limit of } \frac{1.2.3(z) \times r.r+1.r+2(z)}{p+1.p+2.p+3(z) \times t+1.t+2.t+3(z)} \times \frac{a^{p-r}}{nr} \\ &\quad x \text{ increasing ad infinitum.} \end{aligned}$$

$a, n, p+1,$  and  $r-p$  being positive.  $t = r - p - 1.$

$p+y,$  or  $t+y,$  any positive integer.

$v$  and  $w$  any positive integers.

T H E O.

## T H E O R E M VII.

$$\dot{F} = \overline{a + x^n}^p \times x^{rn-1} \dot{x}.$$

$$F = \left\{ \begin{array}{l} K + \frac{x^{rn} z^{p+1}}{a n r} \times I + \frac{t - 1 \cdot x^n}{r + 1 \cdot a} + \frac{t - 1 \cdot t - 2 \cdot x^{2n}}{r + 1 \cdot r + 2 \cdot a^2} (v) \\ + \frac{t - 1 \cdot t - 2 (v)}{r \cdot r + 1 (v) \times a^v} \times \frac{x^{rn+vn} z^p}{n \cdot r + v} \times I - \frac{p x^n}{r'' + 1 \cdot z} + \frac{p \cdot p - 1 \cdot x^{2n}}{r'' + 1 \cdot r'' + 2 \cdot z^2} - \&c. \end{array} \right.$$

$$= \left\{ \begin{array}{l} K + \frac{x^{rn} z^{p+1}}{a n r} \times I + \frac{t - 1 \cdot x^n}{r + 1 \cdot a} + \frac{t - 1 \cdot t - 2 \cdot x^{2n}}{r + 1 \cdot r + 2 \cdot a^2} (v) \\ + \frac{t - 1 \cdot t - 2 (v)}{r \cdot r + 1 (v) \times a^v} \times \frac{x^{r''n} z^p}{n r''} \times I - \frac{p x^n}{r'' + 1 \cdot z} + \frac{p \cdot p - 1 \cdot x^{2n}}{r'' + 1 \cdot r'' + 2 \cdot z^2} - (w) \\ \pm * \frac{t - 1 \cdot t - 2 \cdot t - 3 (v)}{r \cdot r + 1 \cdot r + 2 (v) \times a^v} \times \frac{p \cdot p - 1 \cdot p - 2 (w)}{r'' \cdot r'' + 1 \cdot r'' + 2 (w)} \times \frac{2 x^{2n}}{n \cdot 2 t + 1 \cdot z^2} \end{array} \right.$$

$$\times y + \frac{y^3}{2s+3} + \frac{3y^5}{2s+3 \cdot 2s+5} + \frac{3 \cdot 5y^7}{2s+3 \cdot 2s+5 \cdot 2s+7} \&c.$$

$$r'' = r + v. \quad s = p + r''. \quad t = -p - r. \quad y = \frac{x^n}{2a + x^n}. \quad z = a + x^n.$$

$v$  and  $w$  any positive integers; so that  $w - v$  be  $= 2p + r + 1$ .

\* + or - according as  $w$  is even or odd.

T H E O -

T H E O R E M VIII.

The whole fluent of  $\sqrt{a+x^2}^{-p} \times x^{r-1} x$ ,  
generated whilst  $x$  from 0 becomes infinite,

$$\begin{aligned} \text{is} &= \frac{1.2.3(t)}{r+1.r+2.r+3(t)} \times \frac{a^{r-p}}{nr} = \frac{1.2.3(r)}{t+1.t+2.t+3(r)} \times \frac{a^{r-p}}{nr} \\ &= \left\{ \frac{1.2.3(t+v+y) \times p.p+1.p+2(w)}{r+1.r+2.r+3(t+w+y) \times t+1.t+2.t+3(v)} \times \frac{a^{r-p}}{nr} \right. \\ &\quad \left. \times \left[ 1 + \frac{r-v+w \times y}{t+v+1} + \frac{r-v+w.r-v+w-1 \times y.y-1}{1.2 \times t+v+1.t+v+2} \&c. \right] \right. \\ &= \left\{ \frac{1.2.3(r+v+y) \times p.p+1.p+2(w)}{r+1.r+2.r+3(v) \times t+1.t+2.t+3(r+w+y)} \times \frac{a^{r-p}}{nr} \right. \\ &\quad \left. \times \left[ 1 + \frac{t-v+w \times y}{r+v+1} + \frac{t-v+w.t-v+w-1 \times y.y-1}{1.2 \times r+v+1.r+v+2} \&c. \right] \right. \\ &= \text{the Limit of } \frac{1.2.3(z) \times p.p+1.p+2(z)}{r+1.r+2.r+3(z) \times t+1.t+2.t+3(z)} \times \frac{a^{r-p}}{nr} \end{aligned}$$

$z$  increasing ad infinitum.

$a, n, r,$  and  $p - r$  being positive.  $t = p - r - 1.$

$r + y,$  or  $t + y,$  any positive integer.

$v$  and  $w$  any positive integers.



## T H E O R E M IX.

The *Limit* of  $1 - m^2 \cdot 2^z - m^2 \cdot 3^z - m^2 (z) \times \frac{N^{2z}}{z^{2z+1}}$  is  $= \frac{2}{m}$  fine of  $2md$ .  
 $z$  increasing ad infinitum.

---

## T H E O R E M X.

The *Limit* of  $1 \cdot 2 \cdot 3 (z) \times \frac{N^z}{z^{z+\frac{1}{2}}}$  is  $= 2d^{\frac{1}{2}}$ .

---

## T H E O R E M XI.

The *Limit* of  $bz + a \cdot bz + a + b \cdot bz + a + 2b (z) \times \frac{N^z}{2^{2z} b^z z^z}$  is  $= 2^{\frac{a}{b} - \frac{1}{2}}$ .

---

## T H E O R E M XII.

The *Limit* of  $1 \cdot 3 \cdot 5 (z) \times \frac{N^z}{2^z z^z} = z + 1 \cdot z + 2 \cdot z + 3 (z) \times \frac{N^z}{2^{2z} z^z}$   
 is  $= 2^{\frac{1}{2}}$ .

---

## T H E O R E M XIII.

The *Limit* of  $1 \cdot 5 \cdot 9 (z) \times \frac{N^z}{2^{2z} z^{z-\frac{1}{2}}}$  is  $= \frac{\sqrt{e - \sqrt{e^2 - 2d}}}{2^{\frac{1}{2}} d^{\frac{1}{2}}}$ .

## T H E O R E M XIV.

The *Limit* of  $3 \cdot 7 \cdot 11 (z) \times \frac{N^z}{2^{2z} z^{z+\frac{1}{2}}}$  is  $= \frac{2^{\frac{1}{2}} \sqrt{e + \sqrt{e^2 - 2d}}}{d^{\frac{1}{2}}}$ .

T H E O R E M

T A B L E VIII. 83

T H E O R E M XV.

The *Limit* of 1.4.7  $(z) \times \frac{N^x}{3^x z^{x-\frac{1}{2}}}$  is  $= \frac{3^{\frac{1}{2}} Q^{\frac{1}{2}}}{2^{\frac{1}{2}} d^{\frac{1}{2}}}$ .

T H E O R E M XVI.

The *Limit* of 2.5.8  $(z) \times \frac{N^x}{3^x z^{x+\frac{1}{2}}}$  is  $= \frac{2^{\frac{1}{2}} 3^{\frac{1}{2}} d^{\frac{1}{2}}}{Q^{\frac{1}{2}}}$ .

T H E O R E M XVII.

The *Limit* of 1.7.13  $(z) \times \frac{N^x}{6^x z^{x-\frac{1}{2}}}$  is  $= \frac{Q^{\frac{1}{2}}}{2^{\frac{1}{2}} 3^{\frac{1}{2}} d^{\frac{1}{2}}}$ .

T H E O R E M XVIII.

The *Limit* of 5.11.17  $(z) \times \frac{N^x}{6^x z^{x+\frac{1}{2}}}$  is  $= \frac{2^{\frac{1}{2}} 3^{\frac{1}{2}} d^{\frac{1}{2}}}{Q^{\frac{1}{2}}}$ .

$N = 2.718281 =$  the number whose hyp. log. is 1.

$Q = .982889 = f - \sqrt{f^2 - 3^{\frac{1}{2}} + 1.d.}$

$d = 1.570796 =$  the quadrantal arc of a circ. whose rad. is 1.

$e = 1.910098 =$  the quadrantal arc of an ellipsis whose semi-axes are  $2^{\frac{1}{2}}$  and 1.

$f = 2.674547 =$  the quadrantal arc of another ellipsis whose semi-axes are  $3^{\frac{1}{2}}$  and  $\frac{3^{\frac{1}{2}} \cdot 3^{\frac{1}{2}} + 1}{2^{\frac{1}{2}}}$ .

---

T A B L E IX.  
CONTAINING  
T H E O R E M S  
FOR THE  
CALCULATION of FLUENTS.

---

T H E O R E M I.

The *whole* fluent of  $\frac{x^{p-r} \dot{x}}{a+x^2}^p$ ,

generated whilst  $x$  from 0 becomes infinite,

$$\text{is} = \frac{p-2r}{na^{p-r}} \times \frac{\text{fine of } p \times 90^\circ}{\text{fine of } p-2r \times 90^\circ} \times \frac{1}{r \cdot p-r} + \frac{p}{r+1 \cdot p-r+1} + \frac{p''}{r+2 \cdot p-r+2} \&c.$$

$a$ ,  $n$ ,  $r$ , and  $p-r$  being positive.

$$p'' = \frac{p \cdot p + 1}{1 \cdot 2}, \quad p''' = \frac{p \cdot p + 1 \cdot p + 2}{1 \cdot 2 \cdot 3}, \quad p^{iv} = \frac{p \cdot p + 1 \cdot p + 2 \cdot p + 3}{1 \cdot 2 \cdot 3 \cdot 4}, \quad \&c.$$

T H E O R E M II.

The *whole* fluent of  $\frac{x^{2p-1} \dot{x}}{\sqrt{1-x^2}}$ ,

generated whilst  $x$  from 0 becomes equal to 1,

$$\text{is} = \frac{2^{p+1} \times \text{fine of } pd}{nd} \times \frac{1}{p^2} + \frac{p}{p+2} + \frac{p \cdot p + 1}{2 \cdot p + 4} + \frac{p \cdot p + 1 \cdot p + 2}{2 \cdot 3 \cdot p + 6} \&c.$$

$p$ , and  $2-p$  being positive.

$d$  = the quadrantal arc of a circle whose radius is 1.

T H E O-

T H E O R E M III.

$$F = * \pm 2^p \times \frac{q-p-q-p-2-q-p-4(q)}{p \cdot p + 1 \cdot p + 2(q)} \times \text{cofine of } \overline{p+q} \cdot 90^\circ \times G,$$

$p+q$ , and  $1-p$  being positive. Radius = 1.

$$\left. \begin{array}{l} F = \\ G = \end{array} \right\} \text{the whole fluent of } \left\{ \begin{array}{l} \frac{x^{\frac{1}{2}p + \frac{1}{2}q - 1} x}{\sqrt{1-x^n}} \\ \frac{x^{\frac{1}{2}p + \frac{1}{2}q - 1} x}{(1-x^n)^p} \end{array} \right.$$

generated whilst  $x$  from 0 becomes equal to 1.

\* + or - according as  $q$  is even or odd.

T H E O R E M IV.

$$F = * \pm \frac{q-p-q-p-2-q-p-4(q)}{p \cdot p + 1 \cdot p + 2(q)} \times$$

$$2^p H = \frac{2}{n \cdot q - p} \times 1 - \frac{p}{q-p-2} + \frac{p \cdot p + 1}{q-p-2 \cdot q - p - 4} - (q) \pm$$

$q$  being = 0, or any positive integer;

$p+q$  any positive integer, or fraction.

$$F = \text{the whole fluent of } \dots \frac{x^{\frac{1}{2}p + \frac{1}{2}q - 1} x}{\sqrt{1-x^n}}$$

$$H = \text{the contemporary fluent of } \frac{x^{\frac{1}{2}p + \frac{1}{2}q - 1} x}{1+x^n^p}$$

\* + or - according as  $q$  is even or odd.

T H E O-

## T H E O R E M V.

The *whole* fluent of  $\frac{x^{r^n-1}x}{\sqrt{1-x^n}}$ ,

generated whilst  $x$  from 0 becomes equal to 1,

$$\text{is} = \frac{2}{n} \times \frac{1}{r} + \frac{1}{2.r+1} + \frac{1.3}{2.4.r+2} + \frac{1.3.5}{2.4.6.r+3} (r);$$

the sum of  $r$  terms of the series being exactly equal to *half* the sum of the whole series continued ad infinitum,  $r$  being any positive number whatever!

## T H E O R E M VI.

The *whole* fluent of  $\frac{x^{r^n-1}x}{1-x^n|^p}$ ,

generated whilst  $x$  from 0 becomes equal to 1;

the *whole* fluent of  $\frac{x^{r^n-1}x}{x^n-1|^p}$ ,

generated whilst  $x$  from 1 becomes infinite;

and the *whole* fluent of  $\frac{x^{r^n-1}x}{1+x^n|^p}$ ,

generated whilst  $x$  from 0 becomes infinite;

are to each other

as the sines of  $\overline{p-r} \times 180^\circ$ ,  $r \times 180^\circ$ , and  $p \times 180^\circ$  respectively;

$\overline{p-r}$ ,  $p-r$ , and  $r$  being positive.

TABLE

---

T A B L E X.

CONTAINING

T H E O R E M S

FOR THE

CALCULATION of FLUENTS.

---

T H E O R E M I.

The *whole* fluent of  $\sqrt{a - x^n}^p \times x^{r-1} \dot{x}$ , generated whilst  $x$  from 0 becomes equal to  $a^{\frac{1}{n}}$ , is to the fluent of the same fluxion, generated whilst  $x$  from 0 becomes equal to any quantity  $k$ , as

$$a^r \text{ to } k^{rn} + q k^{rn-n} \cdot \overline{a - k^n} + \frac{q \cdot q - 1}{1 \cdot 2} k^{rn-2n} \cdot \overline{a - k^n}^2 (p + 1) \dot{x}$$

$a, n, p + 1,$  and  $r$  being positive :

$$q = p + r.$$

When  $r$  is  $= p + 1$ , the *whole* fluent is to the part generated whilst  $x$  from 0 becomes equal to  $\frac{a}{2}^{\frac{1}{n}}$  as 2 to 1, whether  $p + 1$  be an integer or not.

T H E O

## T H E O R E M II.

The *whole* fluent of  $x^n - a$  multiplied by  $x^{-n-1}$ , generated whilst  $x$  from  $a^{\frac{1}{n}}$  becomes infinite, is to the fluent of the same fluxion, generated whilst  $x$  from  $a^{\frac{1}{n}}$  becomes equal to any quantity  $k$ , as

$$k^n \text{ to } \overline{k^n - a} + r a \overline{k^n - a}^{r-1} + \frac{r \cdot r - 1}{1 \cdot 2} a^2 \overline{k^n - a}^{r-2} (r - p);$$

$a$ ,  $n$ ,  $p + 1$ , and  $r - p$  being positive.

When  $r$  is  $= 2p + 1$ , the *whole* fluent is to the part generated whilst  $x$  from  $a^{\frac{1}{n}}$  becomes equal to  $\overline{2a}^{\frac{1}{n}}$  as 2 to 1, whether  $r - p$  be an integer or not.

## T H E O R E M III.

The *whole* fluent of  $\frac{x^{r-n-1}}{a + x^n}$ , generated whilst  $x$  from 0 becomes infinite, is to the fluent of the same fluxion, generated whilst  $x$  from 0 becomes equal to any quantity  $k$ , as

$$\overline{a + k^n}^p \text{ to } k^{2n} + p a k^{2n-2} + \frac{p \cdot p - 1}{1 \cdot 2} a^2 k^{2n-4} (p - r + 1);$$

$a$ ,  $n$ ,  $p - r + 1$ , and  $r$  being positive.

When  $2r$  is  $= p + 1$ , the *whole* fluent is to the part generated whilst  $x$  from 0 becomes equal to  $a^{\frac{1}{n}}$  as 2 to 1, whether  $p - r + 1$  be an integer or not.

T A B L E

---

T A B L E XI.

CONTAINING

T H E O R E M S

FOR THE

CALCULATION of FLUENTS.

---

T H E O R E M I.

$$F = \frac{x}{a + bx \times d + ex} \quad \frac{d}{e} \text{ greater or less than } \frac{a}{b}$$

$$F = \frac{1}{bd - ac} \times \text{Log.} \frac{d + ex}{a + bx} \times \frac{a + bx}{d + ex}$$

$c$  the value of  $x$  when  $F = 0$ .

T H E O R E M II.

$$F = \frac{xx}{a + bx \times d + ex} \quad \frac{d}{e} \text{ greater or less than } \frac{a}{b}$$

$$F = \frac{1}{bd - ac} \times \frac{d}{e} \text{Log.} \frac{d + ex}{d + ec} - \frac{a}{b} \text{Log.} \frac{a + bx}{a + bc}$$

$c$  the value of  $x$  when  $F = 0$ .

T H E O R E M III.

$$F = \frac{x^m x}{a + bx \times d + ex} \quad \frac{d}{e} \text{ greater or less than } \frac{a}{b}$$

$$F = \frac{1}{bd - ac} \times \text{fl.} \frac{x^m x}{x + \frac{a}{b}} - \text{fl.} \frac{x^m x}{x + \frac{d}{e}}$$

m

T H E O-



## T H E O R E M IV.

$$F = \frac{x}{a^2 - 2kax - x^2} \quad k \text{ any value positive or negative.}$$

$$F = \frac{1}{2b} \times \text{Log. } \frac{b-c}{b+c} \times \frac{b+y}{b-y}$$

$$y = x + ka, \quad b = a\sqrt{k^2 + 1}. \quad c \text{ the value of } y \text{ when } F = 0.$$


---

## T H E O R E M V.

$$F = \frac{x^2}{a^2 - 2kax - x^2} \quad k \text{ any value positive or negative.}$$

$$F = \frac{1}{2b} \times \overline{b + ka} \cdot \text{Log. } \frac{b+c}{b+y} + \overline{b - ka} \cdot \text{Log. } \frac{b-c}{b-y}$$

$$y, b, \text{ and } c \text{ as in the preceding theorem.}$$


---

## T H E O R E M VI.

$$F = \frac{x^2}{a^2 - 2kax - x^2} \quad k \text{ any value positive or negative.}$$

$$F = \frac{1}{2b} \times \text{fl. } \frac{x^2}{x+p} - \text{fl. } \frac{x^2}{x+q}$$

$$b = a\sqrt{k^2 + 1}. \quad p = ka + b. \quad q = ka - b.$$

T H E O.

T H E O R E M VII.

$$\dot{F} = \frac{x}{a^2 - 2kax + x^2} \quad k \text{ greater than } 1, \text{ or less than } -1.$$

$$F = \frac{1}{2b} \times \text{Log.} \frac{b+c}{b-c} \times \frac{b-y}{b+y}$$

$$y = x - ka. \quad b = a\sqrt{k^2 - 1}. \quad c \text{ the value of } y \text{ when } F = 0.$$


---

T H E O R E M VIII.

$$\dot{F} = \frac{x^2}{a^2 - 2kax + x^2} \quad k \text{ greater than } 1, \text{ or less than } -1.$$

$$F = \frac{1}{2b} \times \frac{b-y}{b+ka} \cdot \text{Log.} \frac{b-y}{b-c} + \frac{b+y}{b-ka} \cdot \text{Log.} \frac{b+y}{b+c}$$

$y, b,$  and  $c$  as in the preceding theorem.

---

T H E O R E M IX.

$$\dot{F} = \frac{x^m}{a^2 - 2kax + x^2} \quad k \text{ greater than } 1, \text{ or less than } -1.$$

$$F = \frac{1}{2b} \times \text{fl.} \frac{x^m}{x-p} - \text{fl.} \frac{x^m}{x-q}$$

$$b = a\sqrt{k^2 - 1}. \quad p = ka + b. \quad q = ka - b.$$

## T H E O R E M X.

$$\dot{F} = \frac{x}{a^2 - 2kax + x^2} \quad k \text{ less than } 1, \text{ but greater than } -1.$$

$$F = K + \frac{1}{2ab} \times \text{circ. arc whose cosine is } \frac{2ax - k.a^2 + x^2}{a^2 + x^2 - 2kax}$$

$$= K + \frac{1}{ab} \times \text{circ. arc whose tangent is } \frac{y}{b}$$

$b$  = the tangent of *half* the arc whose sine is  $b$  and cosine  $k$ .

$$\text{Radius} = 1. \quad h = \sqrt{1 - k^2}. \quad y = \frac{x - a}{x + a}.$$

NOTE. The value of the fluent  $F$ , generated whilst  $x$  from 0 becomes equal to  $g$ , is equal to

$$\frac{1}{2ab} \times \text{the arc, sine } k; \quad \frac{1}{2ab} \times \overline{Q + \text{the arc, sine } k};$$

$$\text{or } \frac{1}{2ab} \times \overline{2Q + \text{the arc, sine } k};$$

according as  $g$  is equal to  $\frac{1-b}{k} \times a$ ,  $a$ , or  $\frac{1+b}{k} \times a$ :

and the *whole* fluent  $F$ , generated whilst  $x$  from 0 becomes infinite, is =  $\frac{1}{ab} \times \overline{Q + \text{the arc, sine } k}$ :

$Q$  being the quadrantal arc.

## T H E O R E M XI.

$$\dot{F} = \frac{x}{a^2 - 2kax + x^2} \quad k \text{ less than } 1, \text{ but greater than } -1.$$

$$F = \frac{1}{2} \text{Log. } \frac{a^2 - 2kax + x^2}{a^2 - 2kad + d^2} + ka \times \text{fl. } \frac{x}{a^2 - 2kax + x^2}.$$

$d$  the value of  $x$  when  $F = 0$ .

THEO-

T H E O R E M XII.

$$F = \frac{x^{m+1}}{a^2 - 2kax + x^2} \quad k \text{ less than } 1, \text{ but greater than } -1.$$

$m$  any positive integer.

$$F = \frac{1}{b'} \text{ into } \left\{ \begin{array}{l} a^{(m+1)} h^{(m+1)} \times \text{fl.} \frac{x}{a^2 - 2kax + x^2} - a^{m+1} h^{(m)} \times \text{fl.} \frac{x}{a^2 - 2kax + x^2} \\ + \frac{b' \cdot x^m - d^m}{m} + \frac{ab' \cdot x^{m-1} - d^{m-1}}{m-1} + \frac{a^2 b'' \cdot x^{m-2} - d^{m-2}}{m-2} (m). \end{array} \right.$$

$B$  the arc whose radius is 1 and cosine  $k$ .

$h$  = the sine of  $B$ .       $h^{(m)}$  = the sine of  $mB$ .

$h''$  = the sine of  $2B$ ;       $h^{(m+1)}$  = the sine of  $(m+1)B$ .

&c.

$d$  the value of  $x$  when  $F = 0$ .

T H E O R E M XIII.

$$F = \frac{x^{-m}}{a^2 - 2kax + x^2} \quad k \text{ less than } 1, \text{ but greater than } -1.$$

$m$  any positive integer.

$$F = \frac{1}{b'} \text{ into } \left\{ \begin{array}{l} \frac{b^{(m)}}{a^{m+1}} \text{Log.} \frac{x}{d} - \frac{b^{(m)}}{a^{m+1}} \times \text{fl.} \frac{x}{a^2 - 2kax + x^2} + \frac{b^{(m+1)}}{a^m} \times \text{fl.} \frac{x}{a^2 - 2kax + x^2} \\ + \frac{b' \cdot d^{1-m} - x^{1-m}}{m-1 \cdot a^2} + \frac{b'' \cdot d^{2-m} - x^{2-m}}{m-2 \cdot a^2} (m-1). \end{array} \right.$$

$d$ , and  $h, h', h'', \&c.$   $h^{(m)}, h^{(m+1)}$  as in the preceding theorem.

T H E O-

## T H E O R E M XIV.

$$\dot{F} = \frac{x^{m-1} \dot{x}}{P + 2kx^n + x^{2n}} = \frac{x^{m-1} \dot{x}}{R' + x^n \times R'' + x^n} \quad k^2 \text{ greater than } P.$$

$$F = \frac{1}{2\sqrt{k^2 - P}} \times \text{fl.} \frac{x^{m-1} \dot{x}}{R' + x^n} - \text{fl.} \frac{x^{m-1} \dot{x}}{R'' + x^n}$$

$$R' = k - \sqrt{k^2 - P}. \quad R'' = k + \sqrt{k^2 - P}.$$

## T H E O R E M XV.

The *whole* fluent of  $\frac{x^{m-1} \dot{x}}{P + 2kx^n + x^{2n}}$ , generated whilst  $x$

from 0 becomes infinite, is  $= \frac{R' \frac{m-n}{n} - R'' \frac{m-n}{n}}{np\sqrt{k^2 - P}} \times Q.$

$k$  and  $P$  being both positive; and  $k^2$  greater than  $P$ .

$m$  any positive integer or fraction less than the integer or fraction  $n$ .

$Q$  = the quadrantal arc of the circle whose radius is 1.

$$p = \text{the sine of the arc } \frac{2m}{n} Q.$$

NOTE. When  $k^2$  is  $= P$ ,  $R'$  and  $R''$  being each  $= k$ , the expression for the value of the *whole* fluent becomes

$$\text{equal to } \frac{2.n - m.k^{\frac{m}{n}-2}}{n^2 p} \times Q.$$

T H E O.

T H E O R E M XVI.

$$\dot{F} = \frac{x^{n+m-1} \dot{x}}{P + 2kx^n + x^{2n}} = \frac{x^{n+m-1} \dot{x}}{R' + x^n \times R'' + x^n} \quad k^2 \text{ greater than } P.$$

$$F = \frac{1}{2\sqrt{k^2 - P}} \times \text{fl.} \frac{R' x^{m-1} \dot{x}}{R' + x^n} - \text{fl.} \frac{R'' x^{m-1} \dot{x}}{R'' + x^n}$$

R' and R'' as in theorem xiv.

T H E O R E M XVII.

The *whole* fluent of  $\frac{x^{n+m-1} \dot{x}}{P + 2kx^n + x^{2n}}$ , generated whilst  $x$ :

from 0 becomes infinite, is  $= \frac{R''^{\frac{n}{2}} - R'^{\frac{n}{2}}}{np\sqrt{k^2 - P}} \times Q.$

$k, P, m, n, p,$  and  $Q$  as in theorem xv.

NOTE. When  $k^2$  is  $= P$ ,  $R'$  and  $R''$  being each  $= k$ , the expression for the value of the *whole* fluent becomes.

equal to  $\frac{2mk^{\frac{n-1}{2}}}{n^2 p} \times Q.$

T H E O-

## T H E O R E M XVIII.

$$\dot{F} = \frac{x^{m-1} \dot{x}}{a^{2n} - 2ka^n x^n + x^{2n}} \quad k \text{ less than } 1, \text{ but greater than } -1.$$

$$F = \frac{a^{m-2n}}{bn} \times \left\{ \begin{array}{l} \text{fl. } \frac{N'b\dot{x}}{b^2+x^2} + \text{fl. } \frac{N''c\dot{x}}{c^2+x^2} + \text{fl. } \frac{N'''d\dot{x}}{d^2+x^2} \&c. \\ -\text{fl. } \frac{M'x\dot{x}}{b^2+x^2} - \text{fl. } \frac{M''x\dot{x}}{c^2+x^2} - \text{fl. } \frac{M'''x\dot{x}}{d^2+x^2} \&c. \end{array} \right.$$

$m$  any positive integer less than the even number  $2n$ .

$$z = \frac{x-a}{x+a}$$

$$\left. \begin{array}{l} M' \text{ and } N' \\ M'' \text{ and } N'' \\ M''' \text{ and } N''' \\ \&c. \end{array} \right\} \text{ sine and cosine of } \left\{ \begin{array}{l} \frac{n-m}{n} \times B. \\ \frac{n-m}{n} \times \overline{B + 360^\circ}. \\ \frac{n-m}{n} \times \overline{B + 2.360^\circ}. \\ \&c. \end{array} \right.$$

$$\left. \begin{array}{l} b = \\ c = \\ d = \\ \&c. \end{array} \right\} \text{ tangent of } \left\{ \begin{array}{l} \frac{A}{n} \\ \frac{A + 180^\circ}{n} \\ \frac{A + 2.180^\circ}{n} \\ \text{to } \frac{A + n - 1.180^\circ}{n} \end{array} \right.$$

( $n$ ).

$A =$  half the arc  $B$  whose sine is  $h$  and cosine  $k$ .

Radius = 1.

T H E O-

T H E O R E M XIX.

The *whole* fluent of  $\frac{x^{n-1}x}{a^{2n} + 2ka^n x^n + x^{2n}}$ , generated whilst  
 $x$  from 0 becomes infinite, is  $= \frac{2qa^{n-2n}}{nbp} \times Q$ .

$k$  less than 1, but greater than  $-1$ .

$m, n, p$ , and  $Q$  as in theorem xv.

$q =$  the sine of  $\frac{n-m}{n}B$ ;  $B$  being the arc whose sine is  $h$  and  
 cosine  $k$ .

Radius  $= r$ .

T H E O R E M XX.

The *whole* fluent of  $\frac{x^{n+m-1}x}{a^{2n} + 2ka^n x^n + x^{2n}}$ , generated whilst  
 $x$  from 0 becomes infinite, is  $= \frac{2qa^{n-m}}{nbp} \times Q$ .

$h, k, m, n, p, B, Q$ , and radius as in the preceding theorem.

$q =$  the sine of  $\frac{m}{n}B$ .

NOTE. When  $m$  is  $= 0$ , the expression for the value of  
 the *whole* fluent becomes equal to  $\frac{B}{nba^n}$ .

n

T H E O.



T H E O R E M XXI.

$$F = \frac{(a+x)^{2n-r} \times x^{m-1}}{a^{2n} - 2ka^n x^n + x^{2n}} \quad k \text{ less than } 1, \text{ but greater than } -1.$$

$$F = \frac{2^{2n-r}}{nba^{r-2n}} \times \left\{ \begin{array}{l} \text{fl. } \frac{N'bz}{b^2+x^2} + \text{fl. } \frac{N''cz}{c^2+x^2} \text{ \&c.} \\ + \text{fl. } \frac{M'xz}{b^2+x^2} + \text{fl. } \frac{M''xz}{c^2+x^2} \text{ \&c.} \end{array} \right.$$

$r$  any positive integer not greater than the even number  $2n$ .

$m$  any positive integer less than  $r$ .

$$z = \frac{x-a}{x+a}$$

$$M' = P' \times \text{fine of } \frac{r-2m}{n}.A \quad N' = P' \times \text{cof. of } \frac{r-2m}{n}.A.$$

$$M'' = P'' \times \text{fine of } \frac{r-2m}{n}.A + 180^\circ. N'' = P'' \times \text{cof. of } \frac{r-2m}{n}.A + 180^\circ.$$

$$M''' = P''' \times \text{fine of } \frac{r-2m}{n}.A + 2.180^\circ. N''' = P''' \times \text{cof. of } \frac{r-2m}{n}.A + 2.180^\circ.$$

&c.

&c.

&c.

&c.

$$\left. \begin{array}{l} b = \\ c = \\ d = \end{array} \right\} \text{tangent of } \left\{ \begin{array}{l} \frac{A}{n} \\ \frac{A + 180^\circ}{n} \\ \frac{A + 2.180^\circ}{n} \\ \frac{A + n - 1.180^\circ}{n} \end{array} \right.$$

&c.

$$\text{to } \frac{A + n - 1.180^\circ}{n}$$

(n)

$A =$  half the arc whose sine is  $b$  and cosine  $k$ .

Radius = 1.

$$P' = \sqrt{1+b^2}^{1/r-1}, P'' = \sqrt{1+c^2}^{1/r-1}, P''' = \sqrt{1+d^2}^{1/r-1}, \text{ \&c.}$$

T H E O-

T H E O R E M XXII.

$$\dot{F} = \frac{(a+x)^{2n-r} \times x^{\frac{1}{2}r-1} \dot{x}}{a^{2n} - 2ka^n x^n + x^{2n}}$$

$$F = \frac{2^{2n-r}}{nba^{\frac{1}{2}r}} \times \text{fl. } \frac{P'b\dot{x}}{b^2+z^2} + \text{fl. } \frac{P''c\dot{x}}{c^2+z^2} \text{ \&c.}$$

*h, k, n, x, b, c, \&c. P', P'', \&c. as in the preceding theorem.  
r any even positive number not greater than 2n.*

T H E O R E M XXIII.

$$\dot{F} = \frac{(a+x)^{2n-r} \times x^{r-1} \dot{x}}{a^{2n} - 2ka^n x^n + x^{2n}}$$

$$F = \text{fl. } \frac{\dot{x}}{1-x} + \frac{2^{2n-r}}{nb} \times \left\{ \begin{array}{l} \text{fl. } \frac{N'b\dot{x}}{b^2+z^2} + \text{fl. } \frac{N''c\dot{x}}{c^2+z^2} \text{ \&c.} \\ + \text{fl. } \frac{M'z\dot{x}}{b^2+z^2} + \text{fl. } \frac{M''z\dot{x}}{c^2+z^2} \text{ \&c.} \end{array} \right.$$

*h, k, n, r, z, b, c, \&c. A, P', P'', \&c. as in theorem XXI.  
M' = P' × sine of  $\frac{r}{n}.A$ .      N' = P' × cof. of  $\frac{r}{n}.A$ .  
M'' = P'' × sine of  $\frac{r}{n}.A + 180^\circ$ .      N'' = P'' × cof. of  $\frac{r}{n}.A + 180^\circ$ .  
M''' = P''' × sine of  $\frac{r}{n}.A + 2.180^\circ$ .      N''' = P''' × cof. of  $\frac{r}{n}.A + 2.180^\circ$ .  
&c.                      &c.                      &c.                      &c.*

## T H E O R E M XXIV.

$$\dot{F} = \frac{x^{v+1} x^{m-1}}{a^{2v} - 2ka^v x^v + x^{2v}} \quad k \text{ less than } 1, \text{ but greater than } -1.$$

$$F = \frac{1}{b^v} \times \left\{ \begin{array}{l} \text{fl. } \frac{a^{v-1} b^{(v)} x^{v-1} x^{m-1}}{a^{2v} - 2ka^v x^v + x^{2v}} - \text{fl. } \frac{a^{v-1} b^{(v-1)} x^{v-1} x^{m-1}}{a^{2v} - 2ka^v x^v + x^{2v}} \\ + \text{fl. } x^{v-1} x \times \frac{h' x^{v-2v} + a^v h'' x^{v-3v} (v-1)}{a^{2v} - 2ka^v x^v + x^{2v}} \end{array} \right.$$

$v$  any positive integer.

$B$  the circ. arc whose rad. is 1 and cosine  $k$ .

$h'$  = the sine of  $B$ .       $h^{(v-1)}$  = the sine of  $\overline{v-1}B$ .

$h''$  = the sine of  $2B$ .       $h^{(v)}$  = the sine of  $vB$ .

$h'''$  = the sine of  $3B$ .       $h^{(v+1)}$  = the sine of  $\overline{v+1}B$ .

&c.

&c.

## T H E O R E M XXV.

$$\dot{F} = \frac{x^{-v+1} x^{m-1}}{a^{2v} - 2ka^v x^v + x^{2v}} \quad k \text{ less than } 1, \text{ but greater than } -1.$$

$$F = \frac{1}{b^v} \times \left\{ \begin{array}{l} \text{fl. } \frac{a^{-v+1} b^{(v+1)} x^{-v+1} x^{m-1}}{a^{2v} - 2ka^v x^v + x^{2v}} - \text{fl. } \frac{a^{-v+1} b^{(v)} x^{-v+1} x^{m-1}}{a^{2v} - 2ka^v x^v + x^{2v}} \\ + \text{fl. } x^{-v+1} x \times \frac{b' x^{-v-2v} + b'' x^{-v-3v} + b''' x^{-v-4v} (v)}{a^{2v} - 2ka^v x^v + x^{2v}} \end{array} \right.$$

$v, B, h', h''$  &c.  $h^{(v)}, h^{(v+1)}$ , as in the preceding theorem.

T H E O

T A B L E XI.

101.

T H E O R E M XXVI.

$$\dot{F} = \frac{x}{\sqrt{a^2 + 2bdx + d^2x^2}}$$

$$F = K + \frac{1}{d} \times \text{Log. } b + dx + \sqrt{a^2 + 2bdx + d^2x^2}.$$

T H E O R E M XXVII.

$$\dot{F} = \frac{x}{\sqrt{a^2 + 2bdx - d^2x^2}}$$

$$F = K + \frac{1}{d} \times \text{Circ. arc, rad. 1, sine } \frac{dx - b}{\sqrt{a^2 + b^2}}.$$

T H E O R E M XXVIII.

$$\dot{F} = \frac{x^2}{\sqrt{a^2 + 2bdx \pm d^2x^2}}$$

$$F = \frac{\sqrt{a^2 + 2bdx \pm d^2x^2}}{\pm d} \mp \frac{b}{d} \times \text{fl. } \frac{x}{\sqrt{a^2 + 2bdx \pm d^2x^2}}$$

T H E O R E M XXIX.

$$\dot{F} = \frac{x^{n-1}}{\sqrt{a^2 + 2bdx^n \pm d^2x^{2n}}}$$

$$F = \frac{1}{n} \times \text{fl. } \frac{y \mp b^{n-1}j}{\sqrt{a^2 \mp b^2 \pm d^2y^2}} = -\frac{2}{n} \times \text{fl. } \frac{b + 2kv + cv^2}{v^2 \mp d^2}^{n-1} v$$

$$h = 2bd \pm cd^2. \quad k = \sqrt{a^2 + 2bcd^2 \pm c^2d^2}.$$

$c=0$ , or any quantity, so that the value of  $k$  be not imaginary.

$$v = \frac{k + \sqrt{a^2 + 2bdx^n \pm d^2x^{2n}}}{x^n - c}$$

$$x^n = \frac{b + 2kv + cv^2}{v^2 \mp d^2} = y \mp \frac{b}{d}$$

T H E O.

## T H E O R E M XXX.

$$\dot{F} = \frac{y^{r-1}j}{b + ky^r \times p + qy^{r \cdot n}}$$

$$F = - \frac{na^{n-m}}{r \cdot bq - kp} \times \int \frac{x^{m-1} \dot{x}}{ka^n + x^n}$$

$$x = \frac{a \cdot bq - kp}{p + qy^r} \quad y = \frac{bq - kp \cdot a^n - px^n}{q^{\frac{1}{r}} x^{\frac{n}{r}}}$$

## T H E O R E M XXXI.

The fluent of  $\frac{y^{r-1}j}{b + y^r \times y^r - p}$ , generated whilst  $y$  from

being equal to  $p^{\frac{1}{r}}$  becomes equal to  $\overline{h + 2p}^{\frac{1}{r}}$ , is

$$= \frac{1}{r \cdot b + p} \times \frac{Q}{s} - N' \text{Log.} \frac{1+b^2}{b^2} - N'' \text{Log.} \frac{1+c^2}{c^2} \&c.$$

$m, n, N', N'', \&c. b, c, \&c. Q, s$  and  $s$  as in Theor. iv. Tab. V.  
 $p$  and  $h + p$  any positive quantities.

## T H E O R E M XXXII.

The whole fluent of  $\frac{y^{r-1}j}{b + y^r \times y^r - p}$ , generated whilst  $y$

from being equal to  $p^{\frac{1}{r}}$  becomes infinite, is  $= \frac{2Q}{r \cdot b + p}$

$m, n, Q$  and  $s$  as in Theorem iv. Tab. V  
 $p$  and  $h + p$  any positive quantities.

T H E O-

T A B L E XI.

103

T H E O R E M XXXIII.

$$F = \frac{y^{\frac{m-1}{n}} j}{(b + ky^r \times p + qy^r)^{\frac{n}{m}}}$$

$$F = \frac{na^{n-m}}{r.kp - bq} \times A. \frac{x^{n-1} z}{ba^x + x^x}$$

$$x = \frac{a.kp - bq \cdot y^{\frac{r}{n}}}{p + qy^r} \quad y = \frac{p^{\frac{1}{r}} x^{\frac{n}{r}}}{kp - bq.a^x - qx^x}$$

T H E O R E M XXXIV.

The fluent of  $\frac{y^{\frac{m-1}{n}} j}{(1 + ky^r \times p - y^r)^{\frac{n}{m}}}$ , generated whilst  $y$  from

$a$  becomes equal to  $\frac{p^{\frac{1}{r}}}{kp + 2}$ , is

$$= \frac{1}{r.kp + 1} \times \frac{Q}{s} + N' \text{Log.} \frac{1 + b^s}{b^s} + N'' \text{Log.} \frac{1 + c^s}{c^s} \&c.$$

$m, n, N', N'', \&c. b, c, \&c. Q, s$  as in Theor. iv. Tab. V.  
 $p$  and  $kp + 1$  any positive quantities.

T H E O R E M XXXV.

The whole fluent of  $\frac{y^{\frac{m-1}{n}} j}{(1 + ky^r \times p - y^r)^{\frac{n}{m}}}$ , generated whilst  $y$

from 0 becomes equal to  $p^{\frac{1}{r}}$ , is  $= \frac{2Q}{r.kp + 1}$ .

$m, n, Q, s$  as in Theorem iv. Tab. V.  
 $p$  and  $kp + 1$  any positive quantities.

T H E O

## T H E O R E M XXXVI.

$$\dot{F} = \frac{y^{r-1}j}{b + 2ky' + ly^{2r} \times p + qy'^{\frac{m}{n}}}$$

$$F = -\frac{nqa^{m-n}}{rt^{\frac{m+n}{n}}} \times \text{fl.} \frac{x^{n+m-1}z}{la^{2n} + 2\frac{kq-lp}{t}a^n x^n + x^{2n}}$$

$$t = \sqrt{hq^2 - 2kpq + lp^2}$$

$$x = \frac{at^{\frac{1}{n}}}{p + qy'^{\frac{1}{n}}} \quad y = \frac{ta^n - px^n}{q^{\frac{1}{r}}x^{\frac{n}{r}}}$$

## T H E O R E M XXXVII.

The whole fluent of  $\frac{y^{r-1}j}{b + 2ky' + y^{2r} \times y' - p}$ , generated whilst  $y$  from being equal to  $p^{\frac{1}{r}}$  becomes infinite, is

$$= \frac{k+p + \sqrt{k^2-b}}{rt^{\frac{2m}{n}} \sqrt{k^2-b} \times \text{the sine of } \frac{2m}{n} Q} - \frac{k+p - \sqrt{k^2-b}}{rt^{\frac{2m}{n}} \sqrt{k^2-b} \times \text{the sine of } \frac{2m}{n} Q} \times Q.$$

$k^2$  greater than  $b$ .  $p$  and  $k+p - \sqrt{k^2-b}$  any positive quantities.

$m$ ,  $n$ , and  $Q$  as in Theorem xv.

$$t = \sqrt{h + 2kp + p^2}$$

NOTE. When  $h$  is  $= k^2$  the expression for the value of the

$$\text{whole fluent becomes} = \frac{2m}{nr \cdot k + p} \times \frac{2m}{n} Q \times \text{the sine of } \frac{2m}{n} Q$$

T H E O.

T H E O R E M XXXVIII.

The whole fluent of  $\frac{y^{r-1}j}{(b+2ky'+y^{2r} \times y^r - p)^{\frac{m}{n}}}$ , generated

whilst  $y$  from being equal to  $p^{\frac{1}{r}}$  becomes infinite, is

$$= \frac{\text{fine of } \frac{m}{n} B}{r t^{\frac{m}{n}} \sqrt{b-k^2} \times \text{the fine of } \frac{m}{n} A} \times A.$$

$k^2$  less than  $h$ .  $p$  any positive quantity.

$m$  any positive integer less than the integer  $n$ .

$$t = \sqrt{h + 2kp + p^2}.$$

$A$  = the semi-circumference of the circle whose radius is  $r$ ;

$B$  = an arc of the same circle whose fine is  $\frac{\sqrt{h-k^2}}{t}$ .

NOTE. If  $m$  be  $= 0$ , the fluent, generated whilst  $y$  from being equal to any quantity  $p^{\frac{1}{r}}$  becomes infinite,

$$\text{will be} = \frac{B}{r \sqrt{h-k^2}}$$

T H E O R E M XXXIX.

The whole fluent of  $\frac{y^{r-1}j}{(b-2py'+y^{2r} \times y^r - p)^{\frac{m}{n}}}$  is

$$= \frac{Q}{r(b-p^2)^{\frac{m+n}{2n}} \times \text{the cosine of } \frac{m}{n} Q}$$

$h$  greater than  $p^2$ .  $m$ ,  $n$ , and  $Q$  as in Theorem xv.

NOTE. If  $m$  be  $= 0$ , the fluent, generated whilst  $y$  from  $p^{\frac{1}{r}}$  becomes infinite, will be  $= \frac{Q}{r \sqrt{h-p^2}}$ .

o

T H E O-



## T H E O R E M XL.

$$\dot{F} = \frac{y^{\frac{r+m}{n}-1} j}{b + 2ky^r + ly^{2r} \times p + qy^r}^{\frac{m}{n}}$$

$$F = \frac{npa^{n-m}}{rt^{\frac{n+s}{n}}} \times \text{fl.} \frac{x^{n+m-1} j}{ba^{2n} + 2 \frac{kp-bq}{t} a^n x^n + x^{2n}}$$

$$t = \sqrt{hq^2 - 2kpg + lp^2}.$$

$$x = \frac{at^{\frac{1}{n}} y^{\frac{r}{n}}}{p + qy^r}^{\frac{1}{n}} \quad y = \frac{p^{\frac{1}{n}} x^{\frac{n}{r}}}{ta^n - qx^n}^{\frac{1}{r}}$$

## T H E O R E M XLI.

The *whole* fluent of  $\frac{y^{\frac{r+m}{n}-1} j}{b + 2ky^r + y^{2r} \times p - y^r}^{\frac{m}{n}}$ , generated whilst  $y$  from 0 becomes equal to  $p^{\frac{1}{r}}$ , is

$$= \frac{\sqrt{b+kp+p\sqrt{k^2-b}}^{\frac{m}{n}} - \sqrt{b+kp-p\sqrt{k^2-b}}^{\frac{m}{n}}}{rt^{\frac{2m}{n}} \sqrt{k^2-b} \times \text{the sine of } \frac{2m}{n} Q} \times Q.$$

$h, p$ , and  $h+kp$  any positive quantities, so that  $k^2$  be greater than  $h$ .  
 $m, n$ , and  $Q$  as in Theorem xv.

$$t = \sqrt{h + 2kp + p^2}.$$

NOTE. When  $h$  is  $= k^2$  the expression for the value of the

$$\text{whole fluent becomes} = \frac{2mpk^{\frac{m-n}{n}}}{nr.k + pl^{\frac{n+s}{n}} \times \text{the sine of } \frac{2m}{n} Q} \times Q.$$

T H E O-

T H E O R E M XLII.

The whole fluent of  $\frac{y^{r+\frac{m}{n}-1}}{b+2ky^r+y^{2r} \times p-y^r t^n}$ , generated

whilst  $y$  from 0 becomes equal to  $p^{\frac{1}{r}}$ , is

$$\frac{b^{\frac{m}{n}} \times \text{the sine of } \frac{m}{n} B}{r t^n \sqrt{b-k^2} \times \text{the sine of } \frac{m}{n} A} \times A.$$

$h, p$ , and  $h+kp$  any positive quantities, so that  $k^2$  be less than  $h$ .  
 $m$  any positive integer less than the integer  $n$ .

$$t = \sqrt{h+2kp+p^2}.$$

$A$  = the semi-circumference of the circle whose radius is 1;

$B$  = an arc of the same circle whose sine is  $\frac{p\sqrt{h-k^2}}{b^{\frac{1}{2}t}}$ .

NOTE. If  $m$  be = 0, the fluent, generated whilst  $y$  from 0 becomes equal to any quantity  $p^{\frac{1}{r}}$ , will be =  $\frac{B}{r\sqrt{b-k^2}}$ .

T H E O R E M XLIII.

The whole fluent of  $\frac{y^{r+\frac{m}{n}-1}}{h-\frac{2b}{p}y^r+y^{2r} \times p-y^r t^n}$  is

$$\frac{p h^{\frac{m-n}{2n}} Q}{r p^2 - b^{\frac{m+n}{2n}} \times \text{the cosine of } \frac{m}{n} Q}$$

$p^2$  greater than  $h$ .  $m, n$ , and  $Q$  as in theorem xv.

NOTE. If  $m$  be = 0, the fluent, generated whilst  $y$  from

0 becomes equal to  $p^{\frac{1}{r}}$ , will be =  $\frac{pQ}{r b^{\frac{1}{2}} \sqrt{p^2-b}}$ .

## T H E O R E M XLIV.

$$\dot{F} = y^{-1} \dot{y} \times \frac{b + ky'}{p + qy'}^{\frac{n}{r}}$$

$$F = \frac{n.bq - kp}{pqr} \times \text{fl.} \frac{x^{n+1} - a}{x^n - \frac{b}{p} \times x^n - \frac{k}{q}}$$

$$x = \frac{b + ky'}{p + qy'}^{\frac{1}{r}} \quad y = \frac{px^n - b}{k - qx^n}^{\frac{1}{r}}$$

## T H E O R E M XLV.

The *whole* fluent of  $y^{-1} \dot{y} \times \frac{y' - b}{p - y'}^{\frac{n}{r}}$ , generated whilst  $y$   
from being equal to  $h^{\frac{1}{r}}$  becomes equal to  $p^{\frac{1}{r}}$ , is

$$= \frac{p^{\frac{n}{r}} - b^{\frac{n}{r}}}{r p^{\frac{n}{r}} \times \text{the sine of } \frac{m}{r} A} \times A.$$

$h$  and  $p$  any positive quantities, so that  $h$  be less than  $p$ .

$m$  any positive integer less than the integer  $n$ .

$A$  as in theorem XLII.

T H E O

THEOREM XLVI.

$$F = \frac{y^{-1}j}{l+y} \times \frac{b+ky^n}{p+qy^n}$$

$$F = \frac{n.bq - kp}{pn.lq - p} \times \text{fl.} \frac{x^{n+m-1}}{x^n - \frac{b}{p} \times x^2 - \frac{b-kl}{p-4}}$$

$x$  and  $y$  as in Theorem XLIV.

THEOREM XLVII.

The whole fluent of  $\frac{y^{-1}j}{l+y} \times \frac{y^n - b}{p - y^n}$ , generated whilst  $y$  from being equal to  $h^{\frac{1}{n}}$  becomes equal to  $p^{\frac{1}{n}}$ , is

$$= \frac{p\sqrt[n]{b+l} - b\sqrt[n]{p+l}}{r.p\sqrt[n]{p+l} \times \text{the sine of } \frac{m}{n}A} \times A.$$

$A, m, n, h,$  and  $p$  as in Theorem XLV.

$b$  positive or negative, so that it be greater than  $-h$ .

NOTE. If  $l$  be  $= -h$ , the whole fluent (of  $\frac{y^{-1}j}{y^n - b} \times \frac{y^n - b}{p - y^n}$ )

$$\text{will be} = \frac{A}{r.h \sqrt[n]{p} \times \text{the sine of } \frac{m}{n}A}$$

THEO-

## T H E O R E M XLVIII.

$$F = \frac{(r \pm z)^n z}{(1+z)^n + (1-z)^n} = \frac{a}{2^{n-1}} \times \frac{(x+a)^{n-1} \times r \pm 1 \cdot x + r \mp 1 \cdot a}{a^2 + z^2}$$

$$F = \frac{1}{n} \times \left\{ \begin{array}{l} \text{fl. } \frac{M'bz}{b^2+z^2} - \text{fl. } \frac{M''cz}{c^2+z^2} + \text{fl. } \frac{M'''dz}{d^2+z^2} - \text{&c.} \\ \pm \text{fl. } \frac{N'z}{b^2+z^2} \mp \text{fl. } \frac{N''z}{c^2+z^2} \pm \text{fl. } \frac{N'''z}{d^2+z^2} \mp \text{&c.} \end{array} \right.$$

$m=0$ , or any even positive number less than the integer  $n-1$ ,  
or any odd positive number less than the integer  $n$ .

$$x = a \times \frac{1+z}{1-z} \quad z = \frac{x-a}{x+a}$$

$$M' = P' \times \text{cofine of } mA'. \quad N' = P' \times \text{fine of } mA'.$$

$$M'' = P'' \times \text{cofine of } mA''. \quad N'' = P'' \times \text{fine of } mA''.$$

$$M''' = P''' \times \text{cofine of } mA'''. \quad N''' = P''' \times \text{fine of } mA'''.$$

&c.                      &c.                      &c.                      &c.

$A', A'', A''', \text{ &c.}$  are circular arcs whose radius is 1 and

tangents  $\frac{b}{r}, \frac{c}{r}, \frac{d}{r}, \text{ &c.}$  respectively.

$b, c, d, \text{ &c.}$  as in Theorem 1, Table V.

$$P' = \frac{r^2 + b^2}{1 + b^2} \frac{1}{2^{n-1}}, \quad P'' = \frac{r^2 + c^2}{1 + c^2} \frac{1}{2^{n-1}}, \quad P''' = \frac{r^2 + d^2}{1 + d^2} \frac{1}{2^{n-1}}, \quad \text{&c.}$$

T H E O-

T H E O R E M. XLIX.

$$F = \frac{r \pm z^m}{1+z} - \frac{r \mp z^m}{1-z} = -\frac{a}{z^{n-1}} \times \frac{(x+a)^{n-m-1} \times r \pm 1 \cdot x + r \mp 1 \cdot a}{a^n - x^n}$$

$$F = \frac{1}{n} \times \left\{ \begin{array}{l} \pm \text{fl. } \frac{M' c z}{c^2 + z^2} \mp \text{fl. } \frac{M' d z}{d^2 + z^2} \pm \&cc. \\ \mp \text{fl. } \frac{\frac{1}{2} r^{m+1} z}{z} - \text{fl. } \frac{N' z z}{c^2 + z^2} + \text{fl. } \frac{N'' z z}{d^2 + z^2} - \&cc. \end{array} \right.$$

$m = 0$ , or any even positive number less than the integer  $n$ ,  
 or any odd positive number less than the integer  $n - 1$ .  
 $a$  and  $z$  as in the preceding theorem.

$$M' = P' \times \text{fine of } mA'. \quad N' = P' \times \text{cofine of } mA'.$$

$$M'' = P'' \times \text{fine of } mA''. \quad N'' = P'' \times \text{cofine of } mA''.$$

$$M''' = P''' \times \text{fine of } mA'''. \quad N''' = P''' \times \text{cofine of } mA'''.$$

$$\&cc. \quad \&cc. \quad \&cc. \quad \&cc.$$

$A', A'', A''', \&cc.$  are circular arcs whose radius is 1 and  
 tangents  $\frac{c}{r}, \frac{d}{r}, \frac{e}{r}, \&cc.$  respectively.

$c, d, e, \&cc.$  as in Theorem XIII. Table V.

$$P' = \frac{r^2 + c^2}{1 + c^2} \pm \frac{r^2 + d^2}{1 + d^2} \pm \frac{r^2 + e^2}{1 + e^2} \pm \&cc.$$

T H E O-

## T H E O R E M L.

$$F = \frac{r \pm z^{2n}}{1 + z^{2n} - 2k \sqrt{1 + z^{2n}} \sqrt{1 - z^{2n}} + 1 - z^{2n}}$$

$$= \frac{a}{2^{2n-1}} \times \frac{(a+b)^{2n-1} \times r \pm 1 \cdot x + r \pm 1 \cdot a}{a^{2n} - 2ka^2 x^2 + x^{2n}}$$

$$F = \frac{1}{2nb} \times \left\{ \begin{array}{l} \text{fl. } \frac{M'bz}{b^2 + z^2} + \text{fl. } \frac{M'cz}{c^2 + z^2} + \text{fl. } \frac{M'dz}{d^2 + z^2} \text{ \&c.} \\ \pm \text{fl. } \frac{N'bz}{b^2 + z^2} \pm \text{fl. } \frac{N'cz}{c^2 + z^2} \pm \text{fl. } \frac{N'dz}{d^2 + z^2} \text{ \&c.} \end{array} \right.$$

$z$  less than 1, but greater than  $-1$ .

$m=0$ , or any positive integer less than the even number  $2n$ .

$x$  and  $z$  as in the preceding theorem.

$$\begin{array}{ll} M' = P' \times \text{cofine of } mA'. & N' = P' \times \text{fine of } mA'. \\ M'' = P'' \times \text{cofine of } mA''. & N'' = P'' \times \text{fine of } mA''. \\ M''' = P''' \times \text{cofine of } mA'''. & N''' = P''' \times \text{fine of } mA'''. \\ \&c. & \&c. \end{array}$$

$A'$ ,  $A''$ ,  $A'''$ , &c. are circular arcs whose radius is 1 and

tangents  $\frac{b}{r}$ ,  $\frac{c}{r}$ ,  $\frac{d}{r}$ , &c. respectively.

$h$ ,  $b$ ,  $c$ ,  $d$ , &c. as in theorem XVIII.

$$P' = \frac{(r^2 + b^2)^{2n}}{1 + b^{2n}}, P'' = \frac{(r^2 + c^2)^{2n}}{1 + c^{2n}}, P''' = \frac{(r^2 + d^2)^{2n}}{1 + d^{2n}}, \text{ \&c.}$$

T A B L E

---

T A B L E XII.

CONTAINING

T H E O R E M S

FOR THE

CALCULATION of FLUENTS.

---

T H E O R E M I.

$$\dot{F} = \frac{x^{-\frac{1}{2}}}{\sqrt{b^2 + 2fx - x^2}} = \frac{x^{-\frac{1}{2}}}{\sqrt{a - x} \times \frac{b^2}{a} + x}$$

$$F = K + \frac{4a^{\frac{1}{2}}}{b^2} \times dc - c^{\frac{1}{2}} = K + \frac{2a^{\frac{1}{2}}}{b^2} \times dc + DP - AD - L.$$

$$F = \frac{a^2 - b^2}{2a}$$

T H E O R E M II.

The fluent of  $\frac{x^{-\frac{1}{2}}}{\sqrt{b^2 + 2fx - x^2}}$ , generated whilst  $x$  from  $o$  becomes equal to any quantity  $k$ , is equal to the fluent of the same fluxion, generated whilst  $x$  from  $\frac{b^2}{a} \times \frac{a-k}{\frac{b^2}{a} + k}$  becomes equal to  $a$ .

NOTE. *All the theorems in this Table refer to the Scheme at the end of it, for the values of the quantities required.*



T A B L E XII.

T H E O R E M III.

$$\dot{F} = \frac{x^{\frac{1}{2}} \dot{x}}{\sqrt{b^2 + 2fx - x^2}} = \frac{x^{\frac{1}{2}} \dot{x}}{\sqrt{\frac{b^2}{a-x} \times \frac{b^2}{a} + x}}$$

$$F = K + \frac{2}{a^{\frac{1}{2}}} \times \overline{2e'e'' - de} = K + \frac{2}{a^{\frac{1}{2}}} \times \overline{L + AD - DP.}$$

*f* as in the preceding theorems.

T H E O R E M IV.

The tangent co  $(= \overline{ak})^{\frac{1}{2}} \times \left( \frac{\frac{a-k}{b^2}}{\frac{a}{a} + k} \right)^{\frac{1}{2}}$  together with the

fluent of  $\frac{\frac{1}{2} a^{\frac{1}{2}} x^{\frac{1}{2}} \dot{x}}{\sqrt{b^2 + 2fx - x^2}}$ , generated whilst *x* from *o* becomes equal to any quantity *k*, is equal to the fluent of the same fluxion, generated whilst *x* from  $\frac{b^2}{a} \times \frac{a-k}{\frac{b^2}{a} + k}$  becomes equal to *a*.

T H E O R E M V.

$$\dot{F} = \frac{y^{-\frac{1}{2}} \dot{y}}{\sqrt{y^2 + 2fy - b^2}} = \frac{y^{-\frac{1}{2}} \dot{y}}{\sqrt{\frac{b^2}{y+a} \times y - \frac{b^2}{a}}}$$

$$F = K + \frac{4a^{\frac{1}{2}}}{b^2} \times \overline{ac + e'e'' - E''} = K + \frac{2a^{\frac{1}{2}}}{b^2} \times \overline{ac + AD - DP.}$$

$x = \frac{b^2}{y}$  *f* as in the preceding theorems.

T H E O R E M VI.

The fluent of  $\frac{y^{-\frac{1}{2}} \dot{y}}{\sqrt{y^2 + 2fy - b^2}}$ , generated whilst *y* from  $\frac{b^2}{a}$  becomes equal to any quantity *k*, is equal to the fluent of the same fluxion, generated whilst *y* from  $\frac{b^2}{a} \times \frac{k+a}{k-\frac{a}{a}}$  becomes infinite.

T H E O-

T H E O R E M VII.

$$\dot{F} = \frac{y^{\frac{1}{2}}j}{\sqrt{y^2 + 2fy - b^2}} = \frac{y^{\frac{1}{2}}j}{\sqrt{y + a} \times y + \frac{b^2}{a}}$$

$F = K + \frac{2}{a^{\frac{1}{2}}} \times AD.$   $f$  and  $x$  as in the two preceding theorems.

T H E O R E M VIII.

$$\dot{F} = \frac{y^{-\frac{1}{2}}j}{\sqrt{y^2 + 2gy + b^2}} = \frac{y^{-\frac{1}{2}}j}{\sqrt{y + m} \times y + \frac{b^2}{m}}$$

$$F = K + \frac{4m^{\frac{1}{2}}}{b^2} \times ac + \frac{2m^{\frac{1}{2}}}{b^2} \times \overline{AD - DP}.$$

$$a = \sqrt{m^2 - b^2}. \quad g = \frac{m^2 + b^2}{2m} \quad x = \frac{ab^2}{my + b^2}$$

T H E O R E M IX.

The fluent of  $\frac{y^{-\frac{1}{2}}j}{\sqrt{y^2 + 2gy + b^2}}$ , generated whilst  $y$  from 0 becomes equal to any quantity  $k$ , is equal to the fluent of the same fluxion generated whilst  $y$  from  $\frac{b^2}{k}$  becomes infinite.

T H E O R E M X.

$$\dot{F} = \frac{y^{\frac{1}{2}}j}{\sqrt{y^2 + 2gy + b^2}} = \frac{y^{\frac{1}{2}}j}{\sqrt{y + m} \times y + \frac{b^2}{m}}$$

$$F = K + \frac{2}{m^{\frac{1}{2}}} \times \overline{DP - ac}. \quad a, g, \text{ and } x \text{ as in theorem VIII.}$$

## T A B L E XII.

## T H E O R E M XI.

$$\dot{F} = \frac{y^{-\frac{1}{2}}j}{\sqrt{2gy - y^2 - b^2}} = \frac{y^{-\frac{1}{2}}j}{\sqrt{\frac{m-y}{m} \times y - \frac{b^2}{m}}}$$

$$F = K + \frac{4m^{\frac{1}{2}}}{b^2} \times dc - e'e'' = K + \frac{2m^{\frac{1}{2}}}{b^2} \times dc + DP - AD - L.$$

$$a = \sqrt{m^2 - b^2}. \quad g = \frac{m^2 + b^2}{2m}. \quad x = \frac{my - b^2}{a}.$$

## T H E O R E M XII.

The fluent of  $\frac{y^{-\frac{1}{2}}j}{\sqrt{2gy - y^2 - b^2}}$ , generated whilst  $y$  from  $\frac{b^2}{m}$  becomes equal to any quantity  $k$ , is equal to the fluent of the same fluxion, generated whilst  $y$  from  $\frac{b^2}{k}$  becomes equal to  $m$ .

## T H E O R E M XIII.

$$\dot{E} = \frac{y^{\frac{1}{2}}j}{\sqrt{2gy - y^2 - b^2}} = \frac{y^{\frac{1}{2}}j}{\sqrt{\frac{m-y}{m} \times y - \frac{b^2}{m}}}$$

$$E = K + \frac{2}{m^{\frac{1}{2}}} \times dc. \quad a, g, \text{ and } x \text{ as in theorem XI.}$$

## T H E O R E M XIV.

The tangent  $eo$  ( $= \frac{m}{k}^{\frac{1}{2}} \times \sqrt{2gk - k^2 - b^2}$ ) together with the fluent of  $\frac{\frac{1}{2}m^{\frac{1}{2}}y^{\frac{1}{2}}j}{\sqrt{2gy - y^2 - b^2}}$ , generated whilst  $y$  from  $\frac{b^2}{m}$  becomes equal to any quantity  $k$ , is equal to the fluent of the same fluxion, generated whilst  $y$  from  $\frac{b^2}{k}$  becomes equal to  $m$ .

T H E O-

T H E O R E M XV.

$$\bar{F} = \frac{y^{-\frac{1}{2}}}{\sqrt{y^2 - 2gy + a^2}} = \frac{y^{-\frac{1}{2}}}{\sqrt{m-y} \times \frac{a^2}{m} - y} \quad y \text{ less than } \frac{a^2}{m}$$

$$E = K + \frac{4m^{\frac{1}{2}}}{b^2} \times ac + e'e'' - E'' = K + \frac{2m^{\frac{1}{2}}}{b^2} \times ac + AD - DP.$$

$$b = \sqrt{m^2 - a^2}. \quad g = \frac{m^2 + a^2}{2m}. \quad x = \frac{a^2 - my}{a}.$$

T H E O R E M XVI.

The fluent of  $\frac{y^{-\frac{1}{2}}}{\sqrt{y^2 - 2gy + a^2}}$ , generated whilst  $y$  from 0 becomes equal to any quantity  $k$ , is equal to the fluent of the same fluxion, generated whilst  $y$  from  $\frac{a^2 - mk}{m - k}$  becomes equal to  $\frac{a^2}{m}$ .

T H E O R E M XVII.

$$\bar{F} = \frac{y^{\frac{1}{2}}}{\sqrt{y^2 - 2gy + a^2}} = \frac{y^{\frac{1}{2}}}{\sqrt{m-y} \times \frac{a^2}{m} - y} \quad y \text{ less than } \frac{a^2}{m}$$

$$E = K + \frac{4m^{\frac{1}{2}}}{b^2} \times g'ac + m.e'e'' - E'' = K + \frac{2m^{\frac{1}{2}}}{b^2} \times \frac{a^2}{m}ac + m.AD - DP.$$

$b, g,$  and  $x$  as in the two preceding theorems.

T H E O R E M XVIII.

The fluent of  $\frac{\frac{1}{2}m^{\frac{1}{2}}y^{\frac{1}{2}}}{\sqrt{y^2 - 2gy + a^2}}$ , generated whilst  $y$  from 0 becomes equal to any quantity  $k$ , is equal to the tangent

to  $(= mk)^{\frac{1}{2}} \times \left. \frac{\frac{a^2 - k}{m}}{m - k} \right)^{\frac{1}{2}}$  together with the fluent of the same fluxion, generated whilst  $y$  from  $\frac{a^2 - mk}{m - k}$  becomes equal to  $\frac{a^2}{m}$ .

T H E O --

## T H E O R E M XIX.

$$\dot{F} = \frac{y^{-\frac{1}{2}j}}{\sqrt{y^2 - 2gy + a^2}} = \frac{y^{-\frac{1}{2}j}}{\sqrt{y-m} \times y - \frac{a^2}{m}} \quad y \text{ greater than } m.$$

$$F = K + \frac{4m^{\frac{1}{2}}}{b^2} \times \overline{ac + e'e'' - E''} = K + \frac{2m^{\frac{1}{2}}}{b^2} \times \overline{ac + AD - DP}.$$

$$b = \sqrt{m^2 - a^2}. \quad g = \frac{m^2 + a^2}{2m}. \quad x = \frac{ab^2}{my - a^2}.$$

## T H E O R E M XX.

The fluent of  $\frac{y^{-\frac{1}{2}j}}{\sqrt{y^2 - 2gy + a^2}}$ , generated whilst  $y$  from  $m$  becomes equal to any quantity  $k$ , is equal to the fluent of the same fluxion, generated whilst  $y$  from  $\frac{mk - a^2}{k - m}$  becomes infinite.

## T H E O R E M XXI.

$$\dot{F} = \frac{y^{\frac{1}{2}j}}{y^2 - 2gy + a^2} = \frac{y^{\frac{1}{2}j}}{\sqrt{y-m} \times y - \frac{a^2}{m}} \quad y \text{ greater than } m.$$

$$F = K + \frac{4m^{\frac{1}{2}}}{b^2} \times \overline{\frac{b^2}{2m} DP + g.ac + m.e'e'' - E''}.$$

$b$ ,  $g$ , and  $x$  as in the two preceding theorems.

## T H E O R E M XXII.

$$\dot{F} = \frac{y^{-\frac{1}{2}j}}{\sqrt{y^2 - 2fy + g^2}}$$

$$F = K + \frac{2^{\frac{1}{2}}a^{\frac{1}{2}}}{b^2} \times \overline{ac + e'e'' - E''} = K + \frac{2^{\frac{1}{2}}a^{\frac{1}{2}}}{b^2} \times \overline{ac + AD - DP}.$$

$$f = \frac{a^2 - b^2}{2a}. \quad g = \frac{a^2 + b^2}{2a}. \quad x = f - y + \sqrt{y^2 - 2fy + g^2}. \quad DP = \sqrt{2ay}.$$

T H E O-

T H E O R E M XXIII.

The fluent of  $\frac{y^{\frac{1}{2}}j}{\sqrt{y^2 - 2fy + g^2}}$ , generated whilst  $y$  from 0 becomes equal to any quantity  $k$ , is equal to the fluent of the same fluxion, generated whilst  $y$  from  $\frac{g^2}{k}$  becomes infinite.

T H E O R E M XXIV.

$$\dot{F} = \frac{y^{\frac{1}{2}}j}{\sqrt{y^2 - 2fy + g^2}}$$

$$F = K + \frac{2^{\frac{1}{2}}a^{\frac{1}{2}}}{b^2} \times \frac{a}{2} \cdot ac + g \cdot e' e'' - E'' + \frac{b^2}{4a} \cdot DP.$$

$f, g, x,$  and  $DP$  as in the two preceding theorems.

T H E O R E M XXV.

$$\dot{F} = \frac{yj}{\sqrt{f \pm y \times g^2 - y^2}} = \frac{-z}{\sqrt{f \pm \sqrt{g^2 - z^2}}}$$

$$F = K + \frac{4a^{\frac{1}{2}}}{b^2} \times g \cdot e' e'' - \frac{a}{2} \cdot de.$$

$$f = \frac{a^2 - b^2}{2a} \quad g = \frac{a^2 + b^2}{2a} \quad x = f \pm y = f \pm \sqrt{g^2 - z^2}$$

T H E O R E M XXVI.

The tangent co ( $= a^{\frac{1}{2}} \times \sqrt{f+k}^{\frac{1}{2}} \times \left| \frac{g-k}{g+k} \right|^{\frac{1}{2}}$ ) together with the fluent of  $\frac{\frac{1}{2}b^{\frac{1}{2}}yj}{\sqrt{f+y \times g^2 - y^2}}$ , generated whilst  $y$  from  $-f$  becomes equal to any quantity  $k$ , is equal to the fluent of the same fluxion, generated whilst  $y$  from  $g \times \frac{g-2f-k}{g+k}$  becomes equal to  $g$ .

T H E O.

## T H E O R E M XXVII.

$$\dot{F} = \frac{yy}{\sqrt{b \pm y \times \frac{1}{4}a^2 - y^2}} = \frac{-z}{\sqrt{b \pm \sqrt{\frac{1}{4}a^2 - z^2}}}$$

$$F = K + \frac{4a^{\frac{1}{2}}}{b^2} \times h.e.e'' - g.d.c.$$

$$g = \frac{1}{2}a + \frac{b^2}{2a} \quad h = \frac{1}{2}a + \frac{b^2}{a} \quad x = \frac{1}{2}a \pm y = \frac{1}{2}a \pm \sqrt{\frac{1}{4}a^2 - z^2}$$

## T H E O R E M XXVIII.

The tangent eo ( $= a^{\frac{1}{2}} \times \left[ \frac{\frac{1}{2}a^2 - k^2}{b+k} \right]^{\frac{1}{2}}$ ) together with the fluent of  $\frac{\frac{1}{2}a^{\frac{1}{2}}yy}{\sqrt{b+y \times \frac{1}{4}a^2 - y^2}}$ , generated whilst  $y$  from  $-\frac{1}{2}a$  becomes equal to any quantity  $k$ , is equal to the fluent of the same fluxion, generated whilst  $y$  from  $-\frac{\frac{1}{2}a^2 + bk}{b+k}$  becomes equal to  $\frac{1}{2}a$ .

## T H E O R E M XXIX.

$$\dot{F} = \frac{yy}{\sqrt{b-y \times y^2 - \frac{b^4}{4a^2}}} = \frac{z}{\sqrt{b - \sqrt{\frac{b^4}{4a^2} + z^2}}}$$

$$F = K + \frac{2}{a^{\frac{1}{2}}} \times e.e'' \quad h = a + \frac{b^2}{2a} \quad x = y - \frac{b^2}{2a} = \left[ \frac{b^4}{4a^2} + z^2 \right]^{\frac{1}{2}} - \frac{b^2}{2a}$$

## T H E O R E M XXX.

The tangent eo ( $= a^{\frac{1}{2}} \times \left[ k - \frac{b^2}{2a} \right]^{\frac{1}{2}} \times \left[ \frac{b-k}{\frac{b^2}{2a} + k} \right]^{\frac{1}{2}}$ ) together with the fluent of  $\frac{\frac{1}{2}a^{\frac{1}{2}}yy}{\sqrt{b-y \times y^2 - \frac{b^4}{4a^2}}}$ , generated whilst  $y$  from  $\frac{b^2}{2a}$  becomes equal to any quantity  $k$ , is equal to the fluent of the same fluxion, generated whilst  $y$  from  $\frac{b^2}{2a} \times \frac{b^2 + 4ab - 2ak}{b^2 + 2ak}$  becomes equal to  $h$ .

T H E O-

T H E O R E M XXXI.

$$\dot{F} = \frac{yj}{\sqrt{b+y \times y^2 - \frac{b^2}{4a^2}}} = \frac{z}{\sqrt{b + \sqrt{\frac{b^2}{4a^2} + z^2}}}$$

$$F = K + \frac{2}{a^{\frac{1}{2}}} \times \overline{DP + e'e'' - E''}.$$

$$h = a + \frac{b^2}{2a} \quad x = \frac{2ab^2}{b^2 + 2ay} = \frac{2ab^2}{b^2 + \sqrt{b^2 + 4a^2z^2}}$$

T H E O R E M XXXII.

$$\dot{F} = \frac{yj}{\sqrt{y-f \times y^2 - g^2}} = \frac{z}{\sqrt{g^2 + z^2}^{\frac{1}{2}} - f}^{\frac{1}{2}}$$

$$F = K + \frac{2a^{\frac{1}{2}}}{b^2} \times a.ac + 2g.e'e'' - E'' + \frac{b^2}{a}.DP.$$

$$f = \frac{a^2 - b^2}{2a} \quad g = \frac{a^2 + b^2}{2a} \quad x = \frac{b^2}{y-f} = \frac{b^2}{g^2 + z^2}^{\frac{1}{2}} - f$$

T H E O R E M XXXIII.

$$\dot{F} = \frac{yj}{\sqrt{y-b \times y^2 - \frac{1}{4}a^2}} = \frac{z}{\sqrt{\frac{1}{4}a^2 + z^2}^{\frac{1}{2}} - b}^{\frac{1}{2}}$$

$$F = K + \frac{4a^{\frac{1}{2}}}{b^2} \times g.ac + h.e'e'' - E'' + \frac{b^2}{2a}.DP.$$

$$g = \frac{1}{2}a + \frac{b^2}{2a} \quad h = \frac{1}{2}a + \frac{b^2}{a} \quad x = \frac{2b^2}{2y-a} = \frac{2b^2}{a^2 + 4z^2}^{\frac{1}{2}} - a$$

T H E O R E M XXXIV.

$$\dot{F} = \frac{yj}{\sqrt{f \pm y \times b^2 + y^2}} = \frac{z}{\sqrt{f \pm \sqrt{z^2 - b^2}}}$$

$$F = K + \left[\frac{2}{a}\right]^{\frac{1}{2}} \times \overline{2AD - DP}.$$

$$f = \frac{a^2 - b^2}{2a} \quad x = \sqrt{b^2 + y^2}^{\frac{1}{2}} \mp y = z \mp \sqrt{z^2 - b^2}. \quad DP = \sqrt{2a.f \pm y}.$$

q

T H E O-



## T H E O R E M XXXV.

$$\dot{F} = x^{-\frac{1}{2}} \dot{x} \times \sqrt{\frac{b^2}{a-x}}. \quad F = K + \frac{2}{a^{\frac{1}{2}}} \times de.$$

## T H E O R E M XXXVI.

The tangent co  $(= \overline{ak})^{\frac{1}{2}} \times \sqrt{\frac{a-k}{\frac{b^2}{a} + k}})^{\frac{1}{2}}$  together with the fluent of  $\frac{1}{2} a^{\frac{1}{2}} x^{-\frac{1}{2}} \dot{x} \times \sqrt{\frac{b^2}{a-x}}$ , generated whilst  $x$  from 0 becomes equal to any quantity  $k$ , is equal to the fluent of the same fluxion, generated whilst  $x$  from  $\frac{b^2}{a} \times \frac{a-k}{\frac{b^2}{a} + k}$  becomes equal to  $a$ .

## T H E O R E M XXXVII.

$$\dot{F} = x^{-\frac{1}{2}} \dot{x} \times \sqrt{\frac{a-x}{\frac{b^2}{a} + x}}.$$

$$F = K + \frac{2}{a^{\frac{1}{2}} b^2} \times \overline{2a^2 + b^2} \cdot de - \overline{2a^2 + 2b^2} \cdot e' e''.$$

## T H E O R E M XXXVIII.

The fluent of  $\frac{1}{2} a^{\frac{1}{2}} x^{-\frac{1}{2}} \dot{x} \times \sqrt{\frac{a-x}{\frac{b^2}{a} + x}}$ , generated whilst  $x$  from 0 becomes equal to any quantity  $k$ , is equal to the tangent co  $(= \overline{ak})^{\frac{1}{2}} \times \sqrt{\frac{a-k}{\frac{b^2}{a} + k}})^{\frac{1}{2}}$  together with the fluent of the same fluxion, generated whilst  $x$  from  $\frac{b^2}{a} \times \frac{a-k}{\frac{b^2}{a} + k}$  becomes equal to  $a$ .

T H E O-

T H E O R E M XXXIX.

$$\dot{F} = y^{-\frac{1}{2}} \times \frac{\sqrt{\frac{y-b^2}{a}}}{y+a} \cdot F = K + \frac{2}{a^{\frac{1}{2}}} \times \overline{DP-ac}.$$

$$x = \frac{b^2}{y} \quad DP = \frac{2}{y} \times \sqrt{\frac{y+a}{y-b^2/a}}$$


---

T H E O R E M XL.

$$\dot{F} = y^{-\frac{1}{2}} \times \frac{\sqrt{\frac{y+a}{b^2}}}{y-\frac{b^2}{a}}.$$

$$F = K + \frac{2}{a^{\frac{1}{2}}b^2} \times \overline{b^2 \cdot DP + 2a^2 + b^2 \cdot ac + 2a^2 + 2b^2 \cdot e' e'' - E''}.$$

$x$ , and  $DP$  as in the preceding theorem.

---

T H E O R E M XLI.

$$\dot{F} = y^{-\frac{1}{2}} \times \frac{\sqrt{\frac{y+\frac{b^2}{m}}{y+m}}}{y+m} \cdot F = K + \frac{2}{m^{\frac{1}{2}}} \times AD.$$

$$a = \sqrt{m^2 - b^2} \quad x = \frac{ab^2}{my + b^2}.$$


---

T H E O R E M XLII.

$$\dot{F} = y^{-\frac{1}{2}} \times \frac{\sqrt{\frac{y+m}{b^2}}}{y+\frac{m}{m}}.$$

$$F = K + \frac{2}{m^{\frac{1}{2}}b^2} \times \overline{b^2 \cdot DP + 2m^2 - b^2 \cdot ac + 2m^2 \cdot e' e'' - E''}.$$

$a$ , and  $x$  as in the preceding theorem.  $DP = \overline{my}^{\frac{1}{2}} \times \frac{\sqrt{\frac{y+m}{b^2}}}{y+\frac{m}{m}}.$

## T A B L E XII.

## T H E O R E M XLIII.

$$\dot{F} = y^{-\frac{1}{2}} \times \frac{\sqrt{\frac{b^2}{m}}}{m-y} \cdot F = K + \frac{2}{m^{\frac{1}{2}}} \times \overline{2c'e'' - dc.}$$

$$a = \sqrt{m^2 - b^2} \quad x = \frac{my - b^2}{a}$$

## T H E O R E M XLIV.

The tangent eo  $(= \frac{m}{k} \times \overline{m - k} \times \overline{k - \frac{b^2}{m}})^{\frac{1}{2}}$  together with the fluent of  $\frac{1}{2} m^{\frac{1}{2}} y^{-\frac{1}{2}} \times \frac{\sqrt{\frac{b^2}{m}}}{m-y}$ , generated whilst  $y$  from  $\frac{b^2}{m}$  becomes equal to any quantity  $k$ , is equal to the fluent of the same fluxion, generated whilst  $y$  from  $\frac{b^2}{k}$  becomes equal to  $m$ .

## T H E O R E M XLV.

$$\dot{F} = y^{-\frac{1}{2}} \times \frac{\frac{m-y}{k^2}}{y - \frac{m}{k}}$$

$$F = K + \frac{2}{m^{\frac{1}{2}} k^2} \times \overline{2m^2 - b^2} \cdot dc - 2m^2 \cdot c'e''.$$

$a$ , and  $x$  as in the two preceding theorems.

## T H E O R E M XLVI.

The fluent of  $\frac{1}{2} m^{\frac{1}{2}} y^{-\frac{1}{2}} \times \frac{\sqrt{\frac{m-y}{b^2}}}{y - \frac{m}{k}}$ , generated whilst from  $\frac{b^2}{m}$  becomes equal to any quantity  $k$ , is equal to the tangent eo  $(= \frac{m}{k} \times \overline{m - k} \times \overline{k - \frac{b^2}{m}})^{\frac{1}{2}}$  together with the fluent of the same fluxion, generated whilst  $y$  from  $\frac{b^2}{k}$  becomes equal to  $m$ .

T H E O-

T A B L E XII.

125

T H E O R E M XLVII.

$$\dot{F} = y^{-\frac{1}{2}} \dot{y} \times \sqrt{\frac{a^2 - y}{m - y}} \cdot y \text{ less than } \frac{a^2}{m}$$

$$F = K + \frac{2}{m^{\frac{1}{2}}} \times \overline{DP - AD} = K + \frac{2}{m^{\frac{1}{2}}} \times \overline{2E'' - 2.e'e'' - ac.}$$

$$b = \sqrt{m^2 - a^2}. \quad x = \frac{a^2 - my}{a}$$

T H E O R E M XLVIII.

The tangent eo  $\left( = \overline{mk}^{\frac{1}{2}} \times \sqrt{\frac{a^2 - k}{m - k}} \right)$  together with the fluent of  $\frac{1}{2} m^{\frac{1}{2}} y^{-\frac{1}{2}} \dot{y} \times \sqrt{\frac{a^2 - y}{m - y}}$ , generated whilst  $y$  from  $o$  becomes equal to any quantity  $k$ , is equal to the fluent of the same fluxion, generated whilst  $y$  from  $\frac{a^2 - mk}{m - k}$  becomes equal to  $\frac{a^2}{m}$ .

T H E O R E M XLIX.

$$\dot{F} = y^{-\frac{1}{2}} \dot{y} \times \sqrt{\frac{m - y}{a^2 - y}} \cdot y \text{ less than } \frac{a^2}{m} \quad F = K + \frac{2}{m^{\frac{1}{2}}} \times ac.$$

$b$ , and  $x$  as in the two preceding theorems.

T H E O R E M L.

The tangent eo  $\left( = \overline{mk}^{\frac{1}{2}} \times \sqrt{\frac{a^2 - k}{m - k}} \right)$  together with the fluent of  $\frac{1}{2} m^{\frac{1}{2}} y^{-\frac{1}{2}} \dot{y} \times \sqrt{\frac{m - y}{a^2 - y}}$ , generated whilst  $y$  from  $o$  becomes equal to any quantity  $k$ , is equal to the fluent of the same fluxion, generated whilst  $y$  from  $\frac{a^2 - mk}{m - k}$  becomes equal to  $\frac{a^2}{m}$ .

T H E C-

## T H E O R E M LI.

$$\dot{F} = y^{-\frac{1}{2}} \times \frac{\sqrt{y - \frac{a^2}{m}}}{y - m}. \quad y \text{ greater than } m. \quad F = K + \frac{2}{m^{\frac{1}{2}}} \times AD.$$

$$b = \sqrt{m^2 - a^2}. \quad x = \frac{ab^2}{my - a^2}.$$

## T H E O R E M LII.

$$\dot{F} = y^{-\frac{1}{2}} \times \frac{\sqrt{\frac{y-m}{a^2}}}{y - \frac{m}{a^2}}. \quad y \text{ greater than } m. \quad F = K + \frac{2}{m^{\frac{1}{2}}} \times \overline{DP - ac}.$$

$$b, \text{ and } x \text{ as in the preceding theorem. } DP = \overline{my}^{\frac{1}{2}} \times \frac{\sqrt{\frac{y-m}{a^2}}}{y - \frac{m}{a^2}}.$$

## T H E O R E M LIII.

$$\dot{F} = y^{-\frac{1}{2}} \times \frac{\sqrt{p^2 + 2 \cdot 1 - q \cdot py + q^2 - q \cdot y^2}}{\sqrt{2p - qy}}.$$

$$F = K + VW. \quad (\text{Fig. 4.})$$

When  $q$  is negative,  $VW$  is an *hyperbola* whose semi-axes are  $\frac{p}{-q}$  and  $\frac{p}{\sqrt{-q}}$ .

When  $q$  is = 0, . . .  $VW$  is a *parabola* whose semi-parameter is  $p$ .

When  $q$  is positive,  $VW$  is an *ellipse* whose semi-axes are  $\frac{p}{q}$  and  $\frac{p}{\sqrt{q}}$ ;

which becomes a *circle* when  $q$  is = 1.

T H E O-

T H E O R E M LIV.

$$\dot{F} = \frac{v^{-\frac{1}{2}} \dot{v}}{v + f \times \sqrt{cv^2 + dv + e}} = \frac{-w^{\frac{1}{2}} \dot{w}}{1 + fw \times \sqrt{c + dw + ew^2}}$$

$$F = \left\{ \begin{array}{l} \frac{1}{2f} \times \text{fl.} \frac{x^{-\frac{1}{2}} \dot{x}}{\sqrt{x^2 - 2dx + d^2 - 4ce}} - \text{fl.} \frac{x^{-\frac{1}{2}} \dot{x}}{x - d + cf + \frac{e}{f}} \\ + \frac{cf^2 - e}{2f^2} \times \text{fl.} \frac{x^{-\frac{1}{2}} \dot{x}}{x - d + cf + \frac{e}{f} \times \sqrt{x^2 - 2dx + d^2 - 4ce}} \end{array} \right.$$

$$x = \frac{cv^2 + dv + e}{v} \quad v = \frac{1}{w} = \frac{x - d + \sqrt{x^2 - 2dx + d^2 - 4ce}}{2c}$$

$$w = \frac{1}{v} = \frac{x - d - \sqrt{x^2 - 2dx + d^2 - 4ce}}{2e}$$

Case 1.  $e = cf^2$ .

$$F = \frac{1}{2f} \times \text{fl.} \frac{x^{-\frac{1}{2}} \dot{x}}{\sqrt{x^2 - 2dx + d^2 - 4ce}} - \text{fl.} \frac{x^{-\frac{1}{2}} \dot{x}}{x - d + 2cf}$$

Case 2.  $e'' = c''f''^2$ .

$$F = \left\{ \begin{array}{l} \frac{1}{2f} \times \text{fl.} \frac{x^{-\frac{1}{2}} \dot{x}}{\sqrt{x^2 - 2dx + d^2 - 4ce}} - \text{fl.} \frac{x^{-\frac{1}{2}} \dot{x}}{x - d + cf + \frac{e}{f}} \\ + \frac{cf^2 - e}{4f^2 f''} \times \text{fl.} \frac{y^{-\frac{1}{2}} \dot{y}}{\sqrt{y^2 - 2d'y + d''^2 - 4c''e''}} - \text{fl.} \frac{y^{-\frac{1}{2}} \dot{y}}{y - d' + 2c''f''} \end{array} \right.$$

$$c'' = 1. \quad d'' = -2d. \quad e'' = d^2 - 4ce. \quad f'' = cf + \frac{e}{f} - d. \quad y = \frac{d'x^2 + d'x + e'}{x}$$

In which cases F will be assigned by the arcs of the conic sections.

NOTE. By substituting repeatedly in the resulting term \* similar to the original value of  $\dot{F}$ , the value of F may be assigned, by means of the arcs of the conic sections, in other particular cases, though not in general.

T H E O-

## T H E O R E M LV.

$$F = \frac{v^{-1} \dot{v}}{\sqrt{2gv+b} \times \sqrt{cv^2+2dv+e}} = \frac{-w^{\frac{1}{2}} \dot{w}}{\sqrt{2g+bw} \times \sqrt{c+2dw+ew^2}}$$

$$F = \begin{cases} -\frac{1}{2b^{\frac{1}{2}}} \times \text{fl.} \frac{g^{\frac{1}{2}} x}{\sqrt{x+e} \times \sqrt{g^2 x^2 + cb^2 - 2dgh + 2eg^2 \cdot x + db - eg}} - \frac{x^{-1} \dot{x}}{\sqrt{x+e}} \\ + \frac{dh-eg}{2b^{\frac{1}{2}}} \times \text{fl.} \frac{x^{-1} \dot{x}}{\sqrt{x+e} \times \sqrt{g^2 x^2 + cb^2 - 2dgh + 2eg^2 \cdot x + db - eg}} \end{cases}$$

$$x = h \times \frac{cv^2 + 2dv + e}{2gv + b} - e. \quad v = \frac{1}{w} = \frac{\sqrt{p^2 x^2 + px + q^2} + gx - q}{cb}$$

$$w = \frac{1}{v} = \frac{\sqrt{g^2 x^2 + px + q^2} - gx + q}{bx}$$

$$p = ch^2 - 2dgh + 2eg^2. \quad q = dh - eg.$$

Case 1.  $dh = eg. \quad x = \frac{cbv^2}{2gv + b}$

$$F = -\frac{1}{2b^{\frac{1}{2}}} \times \text{fl.} \frac{g^{\frac{1}{2}} x^{-\frac{1}{2}}}{\sqrt{x+e} \times \sqrt{g^2 x + cb^2}} - \text{fl.} \frac{x^{-1} \dot{x}}{\sqrt{x+e}}$$

Case 2.  $d''h'' = e''g''. \quad x = \frac{cbv^2 + 2 \cdot db - eg \cdot v}{2gv + b}$

$$F = \begin{cases} -\frac{1}{2b^{\frac{1}{2}}} \times \text{fl.} \frac{g^{\frac{1}{2}}}{\sqrt{x+b''} \times \sqrt{d''x^2 + 2d''x + e''}} - \text{fl.} \frac{x^{-1} \dot{x}}{\sqrt{x+e}} \\ -\frac{dh-eg}{4b^{\frac{1}{2}}b''^{\frac{1}{2}}} \times \text{fl.} \frac{g''y^{-\frac{1}{2}}}{\sqrt{y+e''} \times \sqrt{g''^2 y + e''b''^2}} - \text{fl.} \frac{y^{-1} \dot{y}}{\sqrt{y+e''}} \end{cases}$$

$$e'' = g^2. \quad d'' = \frac{1}{2}ch^2 - dgh + eg^2. \quad e'' = [dh - eg]^2. \quad g'' = \frac{1}{2}. \quad h'' = e.$$

$$y = e \times \frac{d''x^2 + 2d''x + e''}{x+e} - e'' = \frac{eg^2 x^2}{x+e}$$

In which cases F will be assigned by the arcs of the conic sections.

NOTE. The value of F will be so assigned in other particular cases (though not in general) by proceeding as intimated in the Note to the preceding theorem.

T H E O-

T H E O R E M LVI.

$$\dot{F} = \frac{v^{-2}\dot{v}}{\sqrt{2gv+b}\times\sqrt{cv^2+2dv+e}} = \frac{-w^{\frac{1}{2}}\dot{w}}{\sqrt{2g+bw}\times\sqrt{c+2dw+ew^2}}$$

$$F = -\frac{1}{cb} \times \left\{ \begin{array}{l} \frac{\sqrt{2gv+b}\times\sqrt{cv^2+2dv+e}}{v} - \text{fl.} \frac{cgv\dot{v}}{\sqrt{2gv+b}\times\sqrt{cv^2+2dv+e}} \\ + dh+eg \times \text{fl.} \frac{v^{-1}\dot{v}}{\sqrt{2gv+b}\times\sqrt{cv^2+2dv+e}} \end{array} \right.$$

$$w = \frac{1}{v}$$

NOTE. The value of F will be assigned by the arcs of the conic sections, not only when  $dh+eg$  is = 0, but likewise when the fluent of the last written fluxion \* can be so assigned.

T H E O R E M LVII.

$$\dot{F} = \frac{\dot{v}}{\sqrt{c+2dv+ev^2}\times\sqrt{f+2gv+bv^2}} = \frac{-\dot{w}}{\sqrt{cw^2+2dw+e}\times\sqrt{fw^2+2gw+b}}$$

$$F = -\frac{1}{2} \text{fl.} \frac{x^{-\frac{1}{2}}\dot{x}}{\sqrt{p^2+qx+rx^2}}$$

$$p^2 = d^2 - ce. \quad q = ch+ef-2dg. \quad r = g^2 - fh.$$

$$x = \frac{c+2dv+ev^2}{f+2gv+bv^2}. \quad v = \frac{1}{w} = \frac{d-gx+\sqrt{p^2+qx+rx^2}}{bx-e}$$

$$\text{Or } F = -\frac{b^{\frac{1}{2}}}{2} \text{fl.} \frac{j}{\sqrt{e+y}\times\sqrt{P^2+Qy+ry^2}}$$

$$P^2 = [dh-eg]^2. \quad Q = ch^2 - efh - 2dgh + 2eg^2 = hq + 2er.$$

$$y = h \times \frac{c+2dv+ev^2}{f+2gv+bv^2} - e. \quad v = \frac{1}{w} = \frac{1}{b} \times \frac{P-gy+\sqrt{P^2+Qy+ry^2}}{y}$$

NOTE. The value of F will always be assigned by the arcs of the conic sections.

T H E O-



## T H E O R E M LVIII.

$$\dot{F} = \frac{v\dot{v}}{\sqrt{c+2dv+ev^2} \times \sqrt{f+2gv+bv^2}} = \frac{-w^{-1}\dot{w}}{\sqrt{cw^2+2dw+e} \times \sqrt{fw^2+2gw+b}}$$

$$F = \frac{1}{2b^{\frac{1}{2}}} \times \left\{ \begin{array}{l} \text{fl. } \frac{gy}{\sqrt{e+y} \times \sqrt{P^2+Qy+ry^2}} - \text{fl. } \frac{y^{-1}j}{\sqrt{e+y}} \\ - \overline{dh-eg} \times \text{fl. } \frac{y^{-1}j}{\sqrt{e+y} \times \sqrt{P^2+Qy+ry^2}}^* \end{array} \right.$$

$P, Q, r, v, w,$  and  $y$  as in the preceding theorem.

NOTE. The value of  $F$  will be assigned by the *arcs of the conic sections* when  $dh$  is  $= eg$ , and likewise in the particular cases wherein the fluent of the last written fluxion \* can be so assigned.

$$\text{Case 1. } dh = eg. \quad y = \frac{cb - ef}{f + 2gv + bv^2}$$

$$F = \frac{1}{2b^{\frac{1}{2}}} \times \text{fl. } \frac{gy^{-\frac{1}{2}}j}{\sqrt{e+y} \times \sqrt{cb^2 - efb + g^2 - fby}} - \text{fl. } \frac{y^{-1}j}{\sqrt{e+y}}$$

$$\text{Case 2. } ce - d^2 \times h^2 = fh - g^2 \times e^2. \quad y = \frac{cb - ef + 2\overline{db-eg}.v}{f + 2gv + bv^2}$$

$$F = \frac{1}{2b^{\frac{1}{2}}} \times \left\{ \begin{array}{l} \text{fl. } \frac{gy}{\sqrt{e+y} \times \sqrt{P^2 + \frac{P^2}{e}y + ry^2}} - \text{fl. } \frac{y^{-1}j}{\sqrt{e+y}} \\ + \frac{db-eg}{2e^{\frac{1}{2}}} \times \text{fl. } \frac{z^{-\frac{1}{2}}z}{\sqrt{P^2+z} \times \sqrt{4e^2.g^2 - fb+z}} - \text{fl. } \frac{z^{-1}z}{\sqrt{P^2+z}} \end{array} \right.$$

$$z = \frac{e.g^2 - fb.y^2}{y+e}$$

T H E O.

T H E O R E M LIX.

$$F = \frac{v^2 \dot{v}}{\sqrt{c+2dv+ev^2} \times \sqrt{f+2gv+bv^2}} = \frac{-w^{-2} \dot{w}}{\sqrt{cw^2+2dw+e} \times \sqrt{fw^2+2gw+h}}$$

$$F = \frac{1}{2b^{\frac{1}{2}}} \times \left\{ \begin{array}{l} \frac{2}{ey} \times \sqrt{e+y} \times \sqrt{P^2+Qy+ry^2} \\ + 2g \times \text{fl.} \frac{y^{-1} \dot{y}}{\sqrt{e+y}} - 2.dh - eg \times \text{fl.} \frac{y^{-2} \dot{y}}{\sqrt{e+y}} \\ + \frac{fh - 2g^2}{e} \times \text{fl.} \frac{y}{\sqrt{e+y} \times \sqrt{P^2+Qy+ry^2}} \\ + \frac{fb - g^2}{e} \times \text{fl.} \frac{y \dot{y}}{\sqrt{e+y} \times \sqrt{P^2+Qy+ry^2}} \\ + \frac{db - eg \times db + eg}{e} \times \text{fl.} \frac{y^{-1} \dot{y}}{\sqrt{e+y} \times \sqrt{P^2+Qy+ry^2}} * \end{array} \right.$$

P, Q, r, v, w, and y as in the two preceding theorems.

NOTE. The value of F will be assigned by the arcs of the conic sections when  $dh$  is  $= eg$ , or  $dh = -eg$ ; and likewise in the particular cases wherein the fluent of the last written fluxion \* can be so assigned.

# S C H E M E

FOR

## T A B L E XII.

acd (Fig. 1.) is a quadrantal arc of an ellipsis = E'.

Semi-transverse axis  $cd = \sqrt{a^2 + b^2}$ .

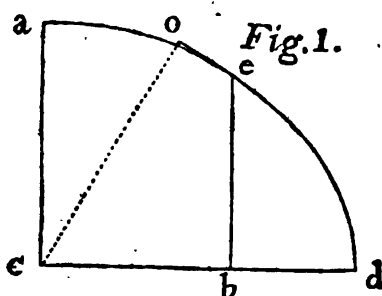
Semi-conjugate axis  $ac = b$ .

Abscissa  $cb = \left( \frac{a^2 + b^2}{a} \right)^{\frac{1}{2}} \times \left( a - x \right)^{\frac{1}{2}}$ .

Ordinate  $be = b \times \left( \frac{x}{a} \right)^{\frac{1}{2}}$ .

Tangent  $co = a - x \times \frac{a^{\frac{1}{2}} x^{\frac{1}{2}}}{\sqrt{b^2 + 2fx - x^2}} = a \times \frac{ax - x^2}{b^2 + ax}$ .

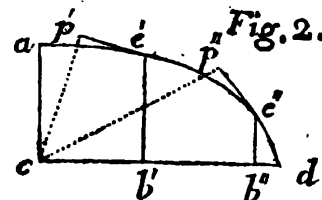
co being perpendicular to eo.



ae'e'd (Fig. 2.) is a quadrantal arc of another ellipsis = E''.

Semi-transverse axis  $cd = \frac{1}{2} \sqrt{a^2 + b^2} + \frac{1}{2} a$ .

Semi-conjugate axis  $ae = \frac{1}{2} \sqrt{a^2 + b^2} - \frac{1}{2} a$ .



e'p' and its equal e''p'' (each =  $\sqrt{a^2 - ax}$ ) are tangents, to which cp', cp'' are perpendiculars.

The abscissa cb', or cb'', corresponding to the ordinate b'e',

or b''e'', is =  $\frac{\sqrt{a^2 + b^2} + a - x \mp \sqrt{\frac{b^2}{a}x + x^2}}{2\sqrt{a^2 + b^2}} \times cd$ .

$ae - de' = e'p' = e''p''$ .

SCHEME

AD (Fig. 3.) is a hyperbola,  
whose vertex is A and cen-  
ter C.

Semi-transverse axis =  $a$ .

Semi-conjugate axis =  $b$ .

DP is a tangent, to which CP  
is perpendicular.

$$CP = \sqrt{ax}.$$

$$DP = \frac{a}{x} \times \sqrt{b^2 + 2fx - x^2}.$$

$$f = \frac{a^2 - b^2}{2a}.$$

The abscissa CAB, corresponding to the ordinate BD,

$$\text{is} = a \times \frac{a}{x} \times \sqrt{\frac{ax + b^2}{a^2 + b^2}}.$$

L, the *limit* of DP - AD, is =  $2E'' - E'$ ;

$$AD \text{ being always} = DP + ac + 2e'e'' - 2E''.$$

VW (Fig. 4.) is an  
arc of a conic sec-  
tion, whose vertex  
is V.

Abscissa VQ, corre-  
sponding to the or-  
dinate QW, =  $y$ ;

RVQ, or VQR being an axis of the section:

$$\text{which axis is} = \frac{2p}{\mp q}.$$

NOTE. All the Theorems in TABLE XII. (where any  
reference is necessary) refer to this Scheme.

TABLE

Fig. 3.

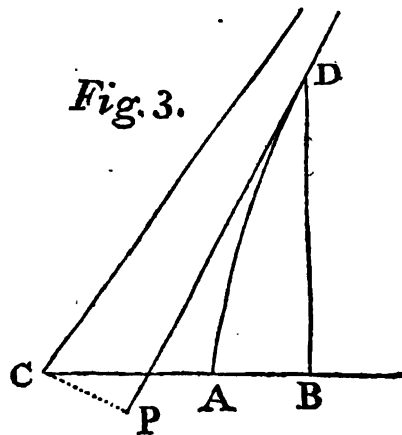
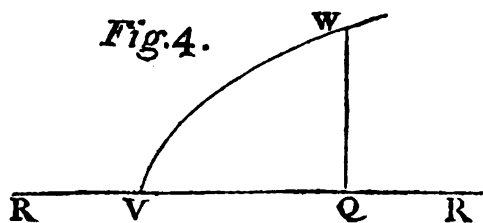


Fig. 4.



---

T A B L E XIII.  
CONTAINING  
T H E O R E M S  
FOR THE  
CALCULATION of FLUENTS.

---

T H E O R E M I.

$$\begin{aligned} \dot{A} &= V^p z^m \dot{z}, & \dot{B} &= V^p z^{m+n} \dot{z}, & \dot{C} &= V^p z^{m+2n} \dot{z}, & \&c. \\ \dot{P}_1 &= V^{p+1} z^m \dot{z}, & \dot{P}_2 &= V^{p+1} z^{m+n} \dot{z}, & \dot{P}_3 &= V^{p+1} z^{m+2n} \dot{z}, & \&c. \\ \dot{Q}_1 &= V^p \dot{V} z^{m+1}, & \dot{Q}_2 &= V^p \dot{V} z^{m+n+1}, & \dot{Q}_3 &= V^p \dot{V} z^{m+2n+1}, & \&c. \end{aligned}$$

$$V = a + bz^n + cz^{2n} + dz^{3n} (t).$$

$$\begin{aligned} P_1 &= aA + bB + cC (t). & Q_1 &= n \times \frac{bB + 2cC + 3dD}{t-1} \\ P_2 &= aB + bC + cD (t). & Q_2 &= n \times \frac{bC + 2cD + 3dE}{t-1} \\ P_3 &= aC + bD + cE (t). & Q_3 &= n \times \frac{bD + 2cE + 3dF}{t-1} \\ \&c. & \&c. & \&c. & \&c. \end{aligned}$$

$$\begin{aligned} V^{p+1} z^{m+1} &= \frac{m+1}{m+1} \times P_1 + \frac{p+1}{p+1} \times Q_1 \\ V^{p+1} z^{m+n+1} &= \frac{m+n+1}{m+n+1} \times P_2 + \frac{p+1}{p+1} \times Q_2 \\ V^{p+1} z^{m+2n+1} &= \frac{m+2n+1}{m+2n+1} \times P_3 + \frac{p+1}{p+1} \times Q_3 \\ &\&c. & \&c. \end{aligned}$$

Hence, if  $t-1$  of the fluents A, B, C, &c.  $P_1, P_2, P_3, \&c.$   $Q_1, Q_2, Q_3, \&c.$  be given, the rest will be determined.

T H E O.

T H E O R E M II.

$$\begin{aligned} \dot{A} &= V^p W^q z^m \dot{z}, & \dot{B} &= V^p W^q z^{m+n} \dot{z}, & \&c. \\ \dot{P}_1 &= V^{p+1} W^q z^m \dot{z}, & \dot{P}_2 &= V^{p+1} W^q z^{m+n} \dot{z}, & \&c. \\ \dot{Q}_1 &= V^p W^{q+1} z^m \dot{z}, & \dot{Q}_2 &= V^p W^{q+1} z^{m+n} \dot{z}, & \&c. \\ \dot{R}_1 &= V^{p+1} W^{q+1} z^m \dot{z}, & \dot{R}_2 &= V^{p+1} W^{q+1} z^{m+n} \dot{z}, & \&c. \\ \dot{S}_1 &= V^p \dot{V} W^{q+1} z^{m+1}, & \dot{S}_2 &= V^p \dot{V} W^{q+1} z^{m+n+1}, & \&c. \\ \dot{T}_1 &= V^{p+1} W^q \dot{W} z^{m+1}, & \dot{T}_2 &= V^{p+1} W^q \dot{W} z^{m+n+1}, & \&c. \end{aligned}$$

$$V = a + bz^n + cz^{2n} + dz^{3n} (t).$$

$$W = a'' + b''z^n + c''z^{2n} + d''z^{3n} (t'').$$

$$P_1 = aA + bB + cC (t). \quad Q_1 = a''A + b''B + c''C (t'').$$

$$P_2 = aB + bC + cD (t). \quad Q_2 = a''B + b''C + c''D (t'').$$

$$\&c. \quad \&c. \quad \&c. \quad \&c. \quad \cdot$$

$$R_1 = aQ_1 + bQ_2 + cQ_3 (t) = a''P_1 + b''P_2 + c''P_3 (t'').$$

$$R_2 = aQ_2 + bQ_3 + cQ_4 (t) = a''P_2 + b''P_3 + c''P_4 (t'').$$

$$\&c. \quad \&c. \quad \&c. \quad \&c.$$

$$S_1 = n \times \overline{bQ_2 + 2cQ_3 + 3dQ_4} (t-1). \quad T_1 = n \times \overline{b''P_2 + 2c''P_3 + 3d''P_4} (t''-1).$$

$$S_2 = n \times \overline{bQ_3 + 2cQ_4 + 3dQ_5} (t-1). \quad T_2 = n \times \overline{b''P_3 + 2c''P_4 + 3d''P_5} (t''-1).$$

$$\&c. \quad \&c. \quad \&c. \quad \&c.$$

$$V^{p+1} W^{q+1} z^{m+1} = \overline{m+1} \times R_1 + \overline{p+1} \times S_1 + \overline{q+1} \times T_1.$$

$$V^{p+1} W^{q+1} z^{m+n+1} = \overline{m+n+1} \times R_2 + \overline{p+1} \times S_2 + \overline{q+1} \times T_2.$$

$$\&c. \quad \&c.$$

Hence, if  $t + t'' - 2$  of the fluents A, B, C, &c. P<sub>1</sub>, P<sub>2</sub>, P<sub>3</sub>, &c. Q<sub>1</sub>, Q<sub>2</sub>, Q<sub>3</sub>, &c. R<sub>1</sub>, R<sub>2</sub>, R<sub>3</sub>, &c. S<sub>1</sub>, S<sub>2</sub>, S<sub>3</sub>, &c. T<sub>1</sub>, T<sub>2</sub>, T<sub>3</sub>, &c. be given, the rest will be determined.

T A B L E

---

T A B L E XIV.  
CONTAINING  
T H E O R E M S  
FOR THE  
CALCULATION of FLUENTS.

---

T H E O R E M I.

The fluent of  $\frac{x}{x} \times \text{fl. } \frac{x}{1+x}$ , generated whilst  $x$  from 0 becomes equal to  $b$ , is equal to  $\frac{a^2}{3}$ , or  $\frac{4a^2}{15} - \frac{1}{3} \text{sq. Log. } \frac{5^{\frac{1}{2}}-1}{2}$ , according as  $b$  is equal to 1, or  $\frac{5^{\frac{1}{2}}-1}{2}$ .

$a$  = the quadrantal arc of the circle whose radius is 1.

T H E O R E M II.

The fluent of  $\frac{x}{x} \times \text{fl. } \frac{x}{1-x}$ , generated whilst  $x$  from 0 becomes equal to  $b$ , is equal to  $\frac{2a^2}{3}$ ,  $\frac{a^2}{3} - \frac{1}{3} \text{sq. Log. } 2$ ,  $\frac{2a^2}{5} - \text{sq. Log. } \frac{5^{\frac{1}{2}}-1}{2}$ , or  $\frac{4a^2}{15} - \text{sq. Log. } \frac{5^{\frac{1}{2}}-1}{2}$ ; according as  $b$  is equal to 1,  $\frac{1}{2}$ ,  $\frac{5^{\frac{1}{2}}-1}{2}$ , or  $\frac{3-\sqrt{5}}{2}$ .

$a$  being as in the preceding theorem.

---

T H E O R E M III.

The *whole* fluent of  $\frac{x}{x} \times \text{fl. } \frac{x^{\frac{1}{2}pn-1}}{1-x^2}^p$ , generated whilst  $x$  from 0 becomes equal to 1, is  $= \frac{aF}{2^{p-1}ns}$ .

$F$  = the contemporary fluent of  $\frac{x^{\frac{1}{2}pn-1}}{\sqrt{1-x^2}}$ .

$a$  as in the preceding theorems.

$s$  = the sine of  $pa$ .  $p$  and  $2-p$  positive.

THEO-

T A B L E    X I V .  
T H E O R E M    I V .

137

The fluent of  $\frac{x}{\sqrt{b^2-x^2}} \times \text{fl. } \frac{x}{\sqrt{1-x^2}}$ , generated whilst  $x$  from 0 becomes equal to  $b$ , (= the contemporary fluent of  $\frac{x}{x} \times \text{fl. } \frac{x}{1-x^2}$ ), is equal to  $\frac{a^2}{2}, \frac{a^2}{3} - \frac{1}{2} \text{sq. Log. } \frac{5^{\frac{1}{2}-1}}{2}, \frac{a^2}{4} - \frac{1}{2} \text{sq. Log. } 2^{\frac{1}{2}-1}$ , or  $\frac{a^2}{6} + \frac{1}{2} \text{sq. Log. } \frac{5^{\frac{1}{2}-1}}{2} - \frac{1}{2} \text{Log. } 5^{\frac{1}{2}-2} \times \text{Log. } \frac{5^{\frac{1}{2}-1}}{2}$ ; according as  $b$  is equal to 1,  $\frac{5^{\frac{1}{2}-1}}{2}$ ,  $2^{\frac{1}{2}-1}$ , or  $5^{\frac{1}{2}-2}$ .  
 $a$  being as in the preceding theorems.

T H E O R E M    V .

The *whole* fluent of  $\frac{x}{\sqrt{1-x^2}} \times \text{fl. } \frac{x}{\sqrt{x+x^2}}$ , generated whilst  $x$  from 0 becomes equal to 1, (=  $2 \times$  the contemporary fluent of  $\frac{x}{x} \times \text{fl. } \frac{x}{\sqrt{1-x^2}}$ ), is =  $2a \text{ Log. } 2$ .  
 $a$  being as in the preceding theorems.

T H E O R E M    V I .

The *whole* fluent of  $\frac{x^{rn-1} \dot{x}}{b^n - x^n)^p} \times \text{fl. } \frac{x^{sn-1} \dot{x}}{c + dx^n)^q}$ , generated whilst  $x$  from 0 becomes equal to  $b$ , is =  $\frac{F'F''}{b^{rn} c^{1-p}}$ .

$$\left. \begin{array}{l} F' = \\ F'' = \end{array} \right\} \text{the contemporary fluent of } \left\{ \begin{array}{l} \frac{x^{rn+sn-1} \dot{x}}{b^n - x^n)^p} \\ \frac{x^{sn-1} \dot{x}}{c + dx^n)^{p+q-1}} \end{array} \right.$$

$\frac{d}{c}$  greater than  $-\frac{1}{b^n}$ .  $p+q=r+s+1$ .  $1-p$  and  $r+s$  positive.

NOTE. This theorem may be of use in computing the fluent of  $\frac{x^{rn-1} \dot{x}}{b^n - x^n)^p} \times \text{fl. } \frac{x^{sn+wn-1} \dot{x}}{c + dx^n)^{q+w}}$ , as the value of  $\text{fl. } \frac{x^{sn+wn-1} \dot{x}}{c + dx^n)^{q+w}}$

may be assigned in terms of  $\text{fl. } \frac{x^{sn-1} \dot{x}}{c + dx^n)^q}$  and algebraic quantities by Theorem XIV, XVI, XVIII, or XX, TAB. VII.  
 $v$  and  $w$  being positive or negative integers.

s

T H E O-



## T H E O R E M VII.

$$\dot{F} = x^{n-1} z' x. \quad z = \text{the hyp. log. of } \frac{x}{c} = \text{fl. } \frac{x}{c}.$$

$$F = K + \frac{x^n}{n} \times z^n - \frac{n}{m} z^{n-1} + \frac{n \cdot n - 1}{m^2} z^{n-2} - \frac{n \cdot n - 1 \cdot n - 2}{m^3} z^{n-3} + \&c.$$

NOTE. The fluent of  $\overline{mz + n} \times x^{n-1} z^{n-1} x$  is  $K + x^n z^n$ .

## T H E O R E M VIII.

$$\dot{F} = z' x. \quad z = \text{fl. } \frac{ax}{\sqrt{b + 2cx + dx^2}}$$

$$F = K + vz' - \frac{r}{d} ay z^{r-1} + \frac{r \cdot r - 1}{d^2} a^2 v z^{r-2} \\ - \frac{r \cdot r - 1 \cdot r - 2}{d^3} a^3 y z^{r-3} + \frac{r \cdot r - 1 \cdot r - 2 \cdot r - 3}{d^4} a^4 v z^{r-4} \\ \&c.$$

$$v = x + \frac{c}{d} \quad y = \sqrt{b + 2cx + dx^2}.$$

## T H E O R E M IX.

$$\dot{F} = x^p y' z' x. \quad y \text{ and } z \text{ as in the preceding theorem.}$$

$$F = K + Px - naQz^{n-1} + n \cdot n - 1 \cdot a^2 Rz^{n-2} - n \cdot n - 1 \cdot n - 2 \cdot a^3 Sz^{n-3} + \&c.$$

$$P = \text{fl. } x^p y' x, \quad Q = \text{fl. } \frac{P x}{y}, \quad R = \text{fl. } \frac{Q x}{y}, \quad S = \text{fl. } \frac{R x}{y}, \quad \&c.$$

T H E O-

T H E O R E M X.

$$Pp + Qqx + Rrx^2 + Ssx^3 + Ttx^4 \&c.$$

$$is = PG + D'F'x + D''F''\frac{x^2}{2} + D'''F'''\frac{x^3}{2.3} \&c.$$

$$G \text{ being } = p + qx + rx^2 + sx^3 \&c.$$

P, Q, R, &c. p, q, r, &c. any invariable quantities.

$$F' = \frac{\dot{G}}{x}, F'' = \frac{\ddot{G}}{x^2}, F''' = \frac{\dot{\dot{G}}}{x^3}, \&c. x \text{ being considered as invariable.}$$

$$D' = Q - P,$$

$$D'' = R - 2Q + P,$$

$$D''' = S - 3R + 3Q - P,$$

$$D^{iv} = T - 4S + 6R - 4Q + P,$$

$$\&c. \qquad \qquad \&c.$$

NOTE. If P, Q, R, &c. be equal to 1,  $\frac{m+n}{n}$ ,  $\frac{m+n.m+n+1}{n.n+1}$ ,  $\frac{m+n.m+n+1.m+n+2}{n.n+1.n+2}$ , &c. respectively; D', D'', D''', &c.

will be respectively equal to  $\frac{m}{n}$ ,  $\frac{m.m-1}{n.n+1}$ ,  $\frac{m.m-1.m-2}{n.n+1.n+2}$ , &c.

And if p, q, r, &c. be equal to 1,  $\frac{e}{f}b$ ,  $\frac{e.e+1}{f.f+1}b^2$ ,  $\frac{e.e+1.e+2}{f.f+1.f+2}b^3$ , &c. respectively; G will be equal to

$$\frac{f-1}{x^{f-1} \times 1-bx} \times \text{fl. } 1-bx \times x^{f-2}.$$

---

T A B L E XV.  
CONTAINING  
T H E O R E M S  
FOR THE  
CALCULATION of FLUENTS.

---

T H E O R E M I.

$$\overline{ax+by+c} \times \dot{x} + \overline{fx+gy+h} \times \dot{y} = 0.$$

$$\overline{2a.x+r+b+f-q.y+s}^{b-f+q} \times \overline{2a.x+r+b+f+q.y+s}^{f-b+q} = K.$$

$$q = \sqrt{b+f}^2 - 4ag. \quad r = \frac{bb-cg}{bf-ag} \quad s = \frac{cf-ab}{bf-ag}$$

NOTE. If  $bf = ag$ ,  $x$  will be = fl.  $\frac{z}{a} \times \frac{fx+ab}{f-bx+ab-bc}$   
 $z$  being =  $ax + by$ .

---

T H E O R E M II.

$$\overline{ax^p y^q + bx^{p-r} y^{q+r} \&c.} \times \dot{x} + \overline{fx^{p-i} y^{q+i} + gx^{p-i} y^{q+i} \&c.} \times \dot{y} = 0.$$

$$\text{Fl. } \frac{z}{x} = - \text{fl. } \frac{fv^{q+i} + gv^{q+i} \&c. \times \dot{v}}{av^q + bv^{q+r} \&c. + fp^{q+i+i} + gp^{q+i+i} \&c.}$$

$$v = \frac{y}{x}$$

THEO-

T H E O R E M III.

If  $aP^p + \overline{b'Q + b''y} \times P^{p-1} + \overline{c'Q^2 + c''Qy + c'''y^2} \times P^{p-2} \&c. = 0$   
 $p a P^{p-1} + \overline{p-1. b' P^{p-2} Q} + \overline{p-2. c' P^{p-3} Q^2} \&c.$   
 $+ \overline{p-1. b'' P^{p-2} y} + \overline{p-2. c'' P^{p-3} Q y} \&c.$   
 $+ \overline{p-2. c''' P^{p-3} y^2} \&c. \quad \&c.$   
 $+ b' P^{p-1} y^{n-1} + 2 c' P^{p-2} Q y^{n-1} + 3 d' P^{p-3} Q^2 y^{n-1} \&c.$   
 $+ c'' P^{p-2} y^{n-2} y' + 2 d'' P^{p-3} Q y^{n-2} y' \&c.$   
 $+ d''' P^{p-3} y^{n-2} y'^2 \&c. \quad \&c.$

will be = 0.

P being  $= y^{n+1} \dot{x} + nxy^n \dot{y}$ , Q  $= y^{n+1} \dot{x} + mxy^n \dot{y}$ .

By means of which equations  $\frac{\dot{x}}{y}$  may be exterminated, and the relation of  $x$  and  $y$  determined by a particular equation of the fluents respecting the first equation: and the same will be the general equation of the fluents respecting the second equation.

Moreover  $\overline{n - m. xy^{n+1} + hy^n = ky^n}$  (derived from the equation  $mnxy^2 + m + n + 1. yy\dot{x} + y^2 \dot{x} = 0$ ) will be an equation of the fluents respecting the first equation; the relation of the invariable quantities  $h$  and  $k$  (by which the equation of the fluents may be adjusted) being expressed by the equation

$ak^p + \overline{b'h + b''} \times k^{p-1} + \overline{c'h^2 + c''h + c'''} \times k^{p-2} \&c. = 0.$

NOTE. The second equation is deduced from the first by taking the fluxions of the several terms, considering  $y$  as invariable, and dividing by  $y^{n-1} \times \overline{mnxy^2 + m + n + 1. yy\dot{x} + y^2 \dot{x}}$  ( $= \dot{P} = y^{n-1} \dot{Q}$ ).

T H E Q.

## T H E O R E M IV.

$$\left. \begin{aligned} & \text{If } \overline{ay^{2m} + b'y^{m+n} + c'y^{2n} \times yx} \\ & + n\overline{ay^{2m} + \frac{m+n}{2}b'y^{m+n} + mc'y^{2n} \times xy} \\ & + \frac{b'y^m + c'y^n}{2} \times y \end{aligned} \right\} \text{be} = 0,$$

$$\begin{aligned} \text{or } & n\overline{ay^{2m} + \frac{m+n}{2}b'y^{m+n} + mc'y^{2n} \times yx} + \frac{b'y^m + c'y^n}{2} \times yx \\ & + n\overline{ay^{2m} + mn\overline{b'y^{m+n} + m^2c'y^{2n} \times x^2y} + nb'y^m + mc'y^n \times xy} \\ & + c''y \text{ be} = 0, \end{aligned}$$

$$\left. \begin{aligned} & \overline{m-n}^2 \times \overline{b^2 - 4ac} \times x^2y^{2m+2n} \\ & + 2.m-n \times \overline{bb'' - 2ac''} \cdot y^m - b'c'' - 2b'd' \cdot y^n \times xy^{m+n} \\ & + \overline{b'y^m + c'y^n}^2 - 4c'' \times \overline{ay^{2m} + b'y^{m+n} + c'y^{2n}} \end{aligned} \right\} \text{will be} = 0.$$

NOTE. This theorem is derived from the preceding one,  $p$  being therein taken equal to 2.

And it is observable that, though, with respect to the first equation here, the third is the *general* equation of the fluents; yet, with respect to the second equation here, and the equation

$$\begin{aligned} & \overline{a \times y^{n+1}x + nxy^ny}^2 + \overline{b' \cdot y^{n+1}x + mxy^ny + b'y \times y^{n+1}x + nxy^ny} \\ & + \overline{c' \cdot y^{n+1}x + mxy^ny}^2 + c''y \times y^{n+1}x + mxy^ny + c''y^2 = 0, \end{aligned}$$

(from which this theorem is derived,) the said third equation is only a *particular* equation of the fluents, unless  $m$  or  $n$  be = 0: and then it is *general*, as well with respect to the second equation as to the first.

T H E O-

T H E O R E M V.

If  $\frac{d+by+ay^2}{f+cx+ax^2} \times \dot{x} + \frac{1}{2} \times \frac{e-bx-cy-2axy}{e-bx-cy-2axy} \times \dot{y}$  be = 0,  
 or  $\frac{f+cx+ax^2}{d+by+ay^2} \times \dot{y} + \frac{1}{2} \times \frac{e-bx-cy-2axy}{e-bx-cy-2axy} \times \dot{x}$  = 0,

$$\text{or } \frac{\dot{x}}{\sqrt{f+cx+ax^2}} = \frac{\dot{y}}{\sqrt{d+by+ay^2}}$$

$$e^2 - 4df - \frac{2be + 4cd \cdot x - 2ce + 4bf \cdot y - 2bc + 4ac \cdot xy}{b^2 - 4ad \cdot x^2 + c^2 - 4af \cdot y^2} \text{ will be } = 0.$$

NOTE. This theorem is derived from the third,  $p$  being therein taken equal to 2,  $m = 0$ ,  $n = -1$ , and  $b', b'', c', c'', c'''$  equal to  $b, -c, d, e, f$  respectively.

And it is observable that, though, with respect to the three fluxional equations here, the fourth equation is the *general* equation of the fluents; yet, with respect to the equation  $a \times \overline{yx - xy} + \overline{bx - cy} \times \overline{yx - xy} + \overline{dx^2 + exy + fy^2} = 0$ , (from which the theorem is derived), the said fourth equation is only a *particular* equation of the fluents.

T H E O R E M VI.

If  $a \times \overline{y^{m+1}x} + \overline{nxy^m y} + b \times \overline{y^{n+1}x} + \overline{mxy^n y} + cy$  be = 0,  
 $\overline{n - m \cdot xy^{m+1}} + \overline{hy^m}$  will be =  $\overline{ky^n}$ ;

the relation of the invariable quantities  $h$  and  $k$  being expressed by the equation  $ak + bh + c = 0$ .

NOTE. This theorem is a particular case of the third,  $p$  being therein taken equal to 1, and  $b', b''$  equal to  $b, c$  respectively.

T H E O.

## T H E O R E M VII.

$$\begin{aligned} \text{If } p a P^{p-1} + \overline{p-1.} b' P^{p-2} Q + \overline{p-2.} c' P^{p-3} Q^2 \&c. \\ &+ \overline{p-1.} b'' P^{p-2} \dot{y} + \overline{p-2.} c'' P^{p-3} Q \dot{y} \&c. \\ &+ \overline{p-2.} c''' P^{p-3} y^2 \&c. \end{aligned}$$

&amp;c.

$$\begin{aligned} + b' P^{p-1} y^{n-m} + 2c' P^{p-2} Q y^{n-m} + 3d' P^{p-3} Q^2 y^{n-m} \&c. \\ + c'' P^{p-2} y^{n-m} \dot{y} + 2d'' P^{p-3} Q y^{n-m} \dot{y} \&c. \\ + d''' P^{p-3} y^{n-m} y^2 \&c. \end{aligned}$$

&amp;c.

$$be = F,$$

$$aP^p + \overline{b'Q + b''\dot{y}} \times P^{p-1} + \overline{c'Q^2 + c''Q\dot{y} + c'''\dot{y}^2} \times P^{p-2} \&c.$$

will be = fl.  $\dot{F}P$ .

$P$ ,  $Q$ , and  $\dot{P}$  ( $= y^{n-m} \dot{Q}$ ) being as in theorem III. and  $y$ , in computing the value of fl.  $\dot{F}P$ , being considered as invariable: which value will be assignable when  $F$  is a proper function of  $P$  and  $\dot{y}$ ; and then, by means of these equations,  $\frac{x}{y}$  may be exterminated, and the *general* relation of  $x$  and  $y$  determined; as well as when  $F$  is = 0.

*Example.* If  $\frac{x^2}{y}$  be =  $hy \times yx - xy$ ; to apply the theorem,  $m$  may be taken =  $-1$ ,  $n = 0$ ,  $p = 2$ ,  $c' = \frac{1}{2}h$ , and  $a$ ,  $b'$ ,  $b''$ ,  $c''$ ,  $c'''$ , &c. each = 0: then  $P$  being =  $x$ ,  $\dot{P} = \dot{x}$ , and  $Q = yx - xy$ ;

$$\dot{F}P = \frac{\dot{x}^2 \dot{x}}{y} \text{ will be } = \overline{hy.y\dot{x} - xy.\dot{x}},$$

$$\text{fl. } \dot{F}P = Ky^2 + \frac{1}{3} \frac{\dot{x}^3}{y} = \frac{1}{3} h \overline{y\dot{x} - xy}^2;$$

and, consequently,  $\overline{3x^3 - k^2} = h \times \overline{2x^2y^3 - 6kxy^3 + hky^6}$ ,  $k$  being the invariable quantity whereby the equation (of the fluents) may be *adjusted*.

T H E O-

T H E O R E M VIII.

If  $\frac{x^{-\frac{1}{2}} \dot{x}}{\sqrt{b + mx + anx^2}}$  be  $= \frac{y^{-\frac{1}{2}} \dot{y}}{\sqrt{bn + my + ay^2}}$ ,

the general equation of the fluents will be

$$axy - cnx - cy \pm 2\sqrt{cm + c^2n + ab} \times x^{\frac{1}{2}}y^{\frac{1}{2}} + b = 0.$$

T H E O R E M IX.

If  $a \times y^{m+1} \dot{x} + n + q \cdot y^m \dot{xy} + n - 1 \cdot qxy^{m-1} \dot{y}^2$   
 $+ b \times y^{m+1} \dot{x} + m + q \cdot y^m \dot{xy} + m - 1 \cdot qxy^{m-1} \dot{y}^2$   
 $+ cy^2$  be  $= 0$ .

the general equation of the fluents will be

$$n - m \cdot xy^{m+1} + gy^{m+1} + hy^{m+1} = ky^{m+1},$$

$y$  being invariable, and the relation of the invariable quantities  $h$  and  $k$  (by which that equation may be *adjusted*) being expressed by the equation

$$a \cdot q - m + 1 \cdot k + b \cdot q - n + 1 \cdot h + c = 0.$$

T H E O R E M X.

If  $\dot{v} + v \dot{y} + \frac{a \cdot n + q \cdot y^m + b \cdot m + q \cdot y^m}{ay^{m+1} + by^{m+1}} v \dot{y} + \frac{a \cdot n - 1 \cdot qy^{m-1} + b \cdot m - 1 \cdot qy^{m-1}}{ay^{m+1} + by^{m+1}}$   
 be  $= 0$ .

the general equation of the fluents will be

$$v = \frac{qy^{m+1} - n - 1 \cdot by^{m+1} - m - 1 \cdot ky^{m+1}}{by^{m+1} + ky^{m+1} - y^{m+1}}$$

$a \cdot q - m + 1 \cdot k$  being  $= b \cdot q - n + 1 \cdot h$ .

NOTE. This theorem is derived from the preceding, by substituting  $v$  for  $\frac{\dot{x}}{xy}$



## T H E O R E M XI.

$Qx^n = Ayx^n + Bxyx^{n-1} + Cx^2y^2x^{n-2} + Dx^3y^3x^{n-3} (n+1)$ .  
the coefficient of the last term being = 1, Q any function  
of  $x$  and  $y$ , and  $x$  invariable.

$$y = K'x^{r'} + K''x^{r''} + K'''x^{r'''} (n) + Fx^{r'} \quad (n) \quad (x)$$

$r', r'', r''', \&c.$  are the roots ( $r$ ) of the equation

$$A + B.r + C.r.r - 1 + D.r.r.r - 1.r - 2 \dots + r.r - 1.r - 2(n) = 0.$$

$$F = \text{fl. } x^{-r'-1} Qx, F'' = \text{fl. } x^{r'-r''-1} F'x, F''' = \text{fl. } x^{r''-r'''-1} F''x, \&c.$$

NOTE I. In deducing the theorem by repeatedly taking the  
fluents, after multiplying by  $x^{-r'-1}, x^{r'-r''-1}, x^{r''-r'''-1}, \&c.$

successively, the terms  $G'x^{r'-1}, \frac{G'x^{r'-r''}x^{n-2}}{r'-r''} + G''x^{r''-2},$

$\frac{G'x^{r'-r''}x^{n-3}}{r'-r', r'-r''} + \frac{G'x^{r''-r'''}x^{n-3}}{r''-r'''} + G'''x^{r'''-3}, \&c.$  successively arise;

the roots  $r', r'', r''', \&c.$  being supposed unequal. But when

$r''$  is =  $r'$ ,  $L'z + K''x^{r'-r''}$  arises (by such operation) instead

of  $K'x^{r'-r'} + K''x^{r''-r''} (= \frac{G'x^{r'-r'}}{r'-r'} + \frac{G''x^{r''-r''}}{r''-r''} (n-1) + \frac{G''x^{r''-r''}}{r''-r''} (n-2))$ ;

$G' \times \text{fl. } x^{r'-r''-1}x$  being then =  $G'z$  instead of  $\frac{G'x^{r'-r''}}{r'-r''}$ ,

$z$  being =  $\text{Log. } x$ : therefore, when  $r'$  is =  $r'' = r''' \dots = r$ ,

$L'z^{n-1} + L''z^{n-2} + L'''z^{n-3} (n-1) + K \times x^{r'}$  must be taken  
(in the value of  $y$ ) instead of  $K'x^{r'} + K''x^{r''} + K'''x^{r'''} (m)$ .

NOTE II. If  $r$  be =  $a + b\sqrt{-1}$ ,  $x^r$  will be =  $x^a N^{b\sqrt{-1}}$ ,  
the value whereof is shewn in the Scholium at the end  
of the Tables.

T H E O-

T H E O R E M XII.

If  $A + Bv + Cv^2 + Dv^3 + \overline{C + D} \cdot \frac{x\dot{v}}{x} + 3D \frac{xv\dot{v}}{x} + D \frac{x^2\ddot{v}}{x^2}$  be = 0,

$v$  will be =  $\frac{xy}{y\dot{x}}$ ;

$\dot{x}$  being invariable, and  $y = K'x^{r'} + K''x^{r''} + K'''x^{r'''};$

where  $r', r'', r'''$  are the roots ( $r$ ) of the equation

$$A + Br + Cr^2 + Dr^3 = 0.$$

T H E O R E M XIII.

If  $A + Bv + Cv^2 + Dv^3 \times \dot{v} + Cw\dot{v} + 3Dwv\dot{v} + Dw\dot{w}$  be = 0,

$v$  will be =  $\frac{xy}{y\dot{x}}$ , and  $w = \frac{x\dot{v}}{x}$ ;

$y$  being as in the preceding theorem.

T H E O R E M XIV.

If  $A\dot{v} + Bv\dot{v} + Cz\dot{v} + Dzv\dot{v} - Dv^2\dot{z} + Dz\dot{z}$  be = 0,

$v$  will be =  $\frac{xy}{y\dot{x}}$ , and  $z = \frac{x\dot{v}}{y\dot{x}} + \frac{x^2\ddot{v}}{y\dot{x}^2}$ ;

$\dot{x}$  and  $y$  being as in theorem XII.

T H E O R E M XV.

If  $A\dot{v} + B + C + D \cdot v\dot{v} + \overline{C + D} \cdot z\dot{v} + Dvz - Dv^2\dot{z} + Dzz\dot{z}$  be = 0,

$v$  will be =  $\frac{xy}{y\dot{x}}$ , and  $z = \frac{x^2\ddot{v}}{y\dot{x}^2}$ ;

$\dot{x}$  and  $y$  being as in theorem XII.

## T H E O R E M XVI.

If,  $z$  being  $= \sqrt{x^2 + y^2}$  and  $x$  invariable,  $z^2 + yy'$  be  $= axz$ ,

$$v \text{ will be } = \frac{k}{1-aw}, \quad v - ay = k,$$

$$yz - ayx = kx, \text{ and } x = \text{fl. } \frac{yy'}{\sqrt{k + 2aky + a^2 - 1.y^2}} :$$

where  $x$  is considered as the abscissa of a curve ( $z$ );  $y$  as the correspondent ordinate, at right angles to the base upon which  $x$  is measured;  $v$  as the normal to the curve, terminated by that base; and  $w$  as the sine (to the radius 1) of the angle made by the said normal and base;  $k$  being an invariable quantity serving to *adjust* the equation of the fluents, with another such quantity that must be added

upon taking the value of fl.  $\frac{yy'}{\sqrt{k + 2aky + a^2 - 1.y^2}}$

NOTE.  $y \text{ is } = vw = \frac{vx}{z}, \quad \dot{y} = \dot{v}w + v\dot{w}, \quad \ddot{y} = -\frac{w.\dot{v}w + v\dot{w}}{w.1 - w^2},$

$$\dot{x} = w\dot{z} = \frac{w.\dot{v}w + v\dot{w}}{\sqrt{1 - w^2}}, \text{ and } \dot{z} = \frac{\dot{v}w + v\dot{w}}{\sqrt{1 - w^2}} :$$

and, by properly substituting as many of these values as may be requisite, the equations of the fluents may be readily deduced from some other fluxional equations.

T A B L E

---

T A B L E XVI.

CONTAINING

T H E O R E M S

FOR THE

CALCULATION of FLUENTS.

---

T H E O R E M I.

$$\dot{F} = N^x x^p \dot{x}.$$

$$F = K + \frac{N^x}{r} \times x^p - \frac{p}{r} x^{p-1} + \frac{p \cdot p - 1}{r^2} x^{p-2} - \frac{p \cdot p - 1 \cdot p - 2}{r^3} x^{p-3} + \&c.$$

N being (everywhere in this Table) the number whose hyp. log. is 1.

NOTE.  $p \times \text{fl. } N^x x^{p-1} \dot{x} + r \times \text{fl. } N^x x^p \dot{x}$  is  $= K + N^x x^p$ .

---

T H E O R E M II.

$$\ddot{F} = r y N^x \dot{x}^2 - \frac{y N^x}{r} \dot{x} \text{ invariable.}$$

$$\dot{F} = y N^x \dot{x} - \frac{y N^x}{r} + K \dot{x}.$$

The fluent of  $\dot{F} N^{-2rx} = -\frac{y N^{-rx}}{r} - \frac{KN^{-2rx}}{2r}$

THEO-

## T H E O R E M III.

$$sy^{r-1}\dot{y} + qy^r\dot{z} + py^r x^{-1}\dot{x} = \dot{P}.$$

P and z any functions of x or y.

$$y^r = x^{-r} N^{-rz} \times K + \text{fl. } x^p N^{rz} \dot{P}.$$

*Example 1.* If z be  $= x^b$ , and P  $= bx^m + cx^n$  &c.  $= bz^{\frac{m}{b}} + cz^{\frac{n}{b}}$  &c.

y<sup>r</sup> will be  $= x^{-r} N^{-rz} \times K + \frac{bm}{b} \times \text{fl. } N^{rz} z^e z + \frac{bn}{b} \times \text{fl. } N^{rz} z^f z$  &c.

$$e \text{ being } = \frac{m+r-b}{b}, f = \frac{n+r-b}{b}.$$

*Example 2.* If  $ax + by$  be  $= Py + Q$ , P and Q being functions of  $\frac{x}{y}$  without x or y being concerned therein;

$$a\dot{x} + b\dot{y} - P\dot{y} = \overline{ar + b - P} \times \dot{y} \text{ will be } = \dot{P}y + \dot{Q},$$

$$\dot{y} - y\dot{w} = \frac{\dot{Q}}{v}, \text{ and } y = N^w \times K + \text{fl. } \frac{\dot{Q}}{vN^w}.$$

$$r \text{ being } = \frac{x}{y}, v = ar + b - P, \text{ and } w = \text{fl. } \frac{\dot{P}}{v}.$$

By which means r may be exterminated, and the equation of the fluents obtained.

*Example 3.* If  $\dot{w} + p\dot{w}v^{-1}\dot{v}$  be  $= b\dot{H}x$ , x being invariable;

$\dot{w}$  (=y) will be  $= v^{-p} \times Kx + \text{fl. } bv^p Hx$ , and  $w = \text{fl. } \frac{Kx + \text{fl. } bv^p Hx}{v^p}$ .

$$\text{And if } pv^{-1}\dot{v} \text{ be } = -\frac{2a^2x^{-1}\dot{x}}{a^2+x^2}, \text{ and } b\dot{H} = -\frac{a\dot{x}}{a^2+x^2};$$

$$v^p \text{ will be } = \frac{a^2+x^2}{x^2},$$

and  $w = k'' + k'x + \frac{1}{2}a \text{Log. } \frac{a^2+x^2}{a^2} - k \times \text{Circ. Arc, rad. } a, \text{ tang. } x.$

REMARK.

REMARK. Sometimes, when the value of fl.  $x^p N^{qz} \dot{P}$  cannot be immediately obtained, it may be of use to suppose a new variable quantity  $u = K + \text{fl. } x^p N^{qz} \dot{P} = y x^p N^{qz}$ ; and, by such means, to exterminate  $x$  (or  $y$ ): the relation of  $u$  and  $y$  (or  $x$ ) being afterwards assignable, as in the following examples.

*Example 4.*  $\dot{y} + qy\dot{x}$  being  $= r N^{mz} y^n \dot{y}$ ,

$y$  (by the theorem) will be  $= N^{-qz} \times K + \text{fl. } r N^{mz+qz} y^n \dot{y}$ .  
Now, supposing  $u = N^{qz} y = K + \text{fl. } r N^{mz+qz} y^n \dot{y}$ ,  $N^{mz+qz}$

will be  $= \left(\frac{u}{y}\right)^{\frac{m+q}{q}}$ ,  $r N^{mz+qz} y^n \dot{y} = r u^{\frac{m+q}{q}} y^{-\frac{m+q}{q}} \dot{y} = \dot{u}$ ,

and  $r \times \text{fl. } y^{-\frac{m+q}{q}} \dot{y} = \text{fl. } u^{-\frac{m+q}{q}} \dot{u}$ .

*Example 5.* If  $\dot{y} - y \frac{\dot{x}}{x}$  be  $= bx\dot{x} + \frac{cy^2 \dot{y}}{x}$ ,

$y$  (by the theorem) will be  $= x \times K + \text{fl. } bx\dot{x} + \text{fl. } c \frac{y^2 \dot{y}}{x}$ .

Suppose  $u = \frac{y}{x} = K + \text{fl. } bx\dot{x} + \text{fl. } c \frac{y^2 \dot{y}}{x}$ :

then,  $x$  being  $= \frac{y}{u}$ ,  $\dot{x} = \frac{\dot{y}}{u} - \frac{y\dot{u}}{u^2}$ ,

and  $bx\dot{x} + c \frac{y^2 \dot{y}}{x} = \frac{by}{u} - \frac{by\dot{u}}{u^2} + cu^2 \dot{y} = \dot{u}$ ,

$\dot{y} - y \frac{\dot{x}}{x}$  will be  $= \frac{\frac{1}{b} u \dot{u}}{1 + \frac{c}{b} u^2}$ :

whence, by applying the theorem a second time,

$y$  is found  $= \frac{u}{1 + \frac{c}{b} u^2} \times K + \frac{1}{b} \text{fl. } \frac{\dot{u}}{1 + \frac{c}{b} u^2}$ ;

$\frac{u^{-1} \dot{u}}{1 + \frac{c}{b} u^2}$  being  $= \frac{\dot{w}}{w}$ , when  $w$  is  $= \frac{u}{1 + \frac{c}{b} u^2}$ , by Theor. I. and III. Tab. I.

## T H E O R E M IV.

$Q\dot{x}^n = Ay\dot{x}^n + B\dot{y}\dot{x}^{n-1} + C\ddot{y}\dot{x}^{n-2} + D\dot{y}\dot{x}^{n-3} (n+1)$ ;  
the coefficient of the last term being = 1, Q any function  
of  $x$  and  $y$ , and  $\dot{x}$  invariable.

$$y = K'N^{r'x} + K''N^{r''x} + K'''N^{r'''x} (n) + F^{(n)}N^{r'x}.$$

$r', r'', r''', \&c.$  the roots ( $r$ ) of the equation  $A + Br + Cr^2 + Dr^3 \dots r^n = 0$ .

$$F' = \text{fl. } QN^{-r'x}, F'' = \text{fl. } N^{r'x-r''x} F'x, F''' = \text{fl. } N^{r'x-r''x} F''x,$$

$$F^{iv} = \text{fl. } N^{r''x-r'''x} F'''x, \&c.$$

NOTE. In deducing the theorem by repeatedly taking the  
fluents, after multiplying by  $N^{-r'x}$ ,  $N^{r'x-r''x}$ ,  $N^{r'x-r''x}$ , &c.  
succesively, the terms  $G'\dot{x}^{n-1}$ ,  $\frac{G'N^{r'x-r''x}\dot{x}^{n-2}}{r'-r''} + G''\dot{x}^{n-2}$ ,  
 $\frac{G'N^{r'x-r''x}\dot{x}^{n-3}}{r'-r'', r''-r'''} + \frac{G''N^{r''x-r'''x}\dot{x}^{n-3}}{r''-r'''}$  +  $G'''\dot{x}^{n-3}$ , &c. succesively  
arise; the roots  $r', r'', r''', \&c.$  being supposed unequal.  
But when  $r''$  is =  $r'$ ,  $\overline{L'x + K''} \times N^{r'x-r''x}$  arises (by such ope-  
ration) instead of  $K'N^{r'x-r''x} + K''N^{r''x-r''x}$  ( $= \frac{G'N^{r'x-r''x}}{r'-r'', r''-r'''}(n-1)$   
 $+ \frac{G''N^{r''x-r'''x}}{r''-r''', r'''-r^{iv}}(n-2)$ );  $G' \times \text{fl. } N^{r'x-r''x} \dot{x}$  being then =  $G'x$  in-  
stead of  $\frac{G'N^{r'x-r''x}}{r'-r''}$ : therefore, when  $r'$  is =  $r'' = r''' \dots = r^{(n)}$ ,  
 $L'x^{n-1} + L''x^{n-2} + L'''x^{n-3} (m-1) + \overline{K} \times N^{r'x}$  must be taken  
(in the value of  $y$ ) instead of  $K'N^{r'x} + K''N^{r''x} + K'''N^{r'''x} (m)$ .

T H E O-

T A B L E XVI.

153

T H E O R E M V.

If  $A + Bv + Cv^2 + Dv^3 + C\frac{\dot{v}}{x} + 3D\frac{\dot{v}\dot{v}}{x^2} + D\frac{\ddot{v}}{x^3}$  be = 0.

$v$  will be =  $\frac{j}{y^2}$ ;

$x$  being invariable, and  $y = K'N^{r'} + K''N^{r''} + K'''N^{r'''};$   
 where  $r', r'', r'''$  are the roots ( $r$ ) of the equation  
 $A + Br + Cr^2 + Dr^3 = 0.$

T H E O R E M VI.

If  $Ax^3 + Bx^2 - Cx + 2D.w + Cx - 3D.w^2 + Dw^3$

$+ Cx - 2D.\frac{x\dot{w}}{x} + 3D\frac{xw\dot{w}}{x^2} + D\frac{x^2\ddot{w}}{x^3}$  be = 0,

$w$  will be =  $\frac{xy}{y^2}$ ;  $y$  being as in the preceding theorem.

T H E O R E M VII.

$\ddot{y} - c^2 x^{4m} y \dot{x}^2 = \dot{v} + v^2 \dot{x} - c^2 x^{4m} \dot{x} = 0.$   $v = \frac{j}{y^2}$

$$y = \left\{ \begin{array}{l} \frac{K'N^{r'} + K''N^{r''}}{2c.2m+1} \times x^{-m} + \frac{Cx^{-5m-2}}{2c.2m+1} + \frac{Ex^{-9m-4}}{2c.2m+1} \&c. \\ + K'N^{r'} + K''N^{r''} \times \frac{Bx^{-3m-1}}{2c.2m+1} + \frac{Dx^{-7m-3}}{2c.2m+1} + \frac{Fx^{-11m-5}}{2c.2m+1} \&c. \end{array} \right.$$

$$r = \frac{c}{2m+1} \quad z = x^{2m+1}.$$

$$B = -\frac{m}{m+1}, C = -\frac{1}{2} \cdot \frac{3m+1}{3m+2} \times B, D = -\frac{1}{2} \cdot \frac{5m+2}{5m+3} \times C,$$

$$E = -\frac{1}{2} \cdot \frac{7m+3}{7m+4} \times D, F = -\frac{1}{2} \cdot \frac{9m+4}{9m+5} \times E, \&c.$$

NOTE. Both the series will terminate, if,  $n$  being any integer whatever,  $m$  be =  $-\frac{n}{2n \pm 1}$ .

And then, from the above value of  $y$ , that of  $v (= \frac{j}{y^2})$  may be easily obtained.

u

T H E O-



## T H E O R E M VIII.

$F' = \text{fl. } N^{\prime\prime} x^{\rho} \dot{z}, \quad F'' = \text{fl. } N^{\prime\prime} x^{\rho-1} y \dot{z}, \quad F''' = \text{fl. } N^{\prime\prime} x^{\rho-2} y^2 \dot{z}, \quad \&c.$   
 $G' = \text{fl. } N^{\prime\prime} x^{\rho-1} \dot{z}, \quad G'' = \text{fl. } N^{\prime\prime} x^{\rho-2} y \dot{z}, \quad G''' = \text{fl. } N^{\prime\prime} x^{\rho-3} y^2 \dot{z}, \quad \&c.$   
 $H' = \text{fl. } N^{\prime\prime} x^{\rho-2} \dot{z}, \quad H'' = \text{fl. } N^{\prime\prime} x^{\rho-3} y \dot{z}, \quad H''' = \text{fl. } N^{\prime\prime} x^{\rho-4} y^2 \dot{z}, \quad \&c.$   
 $\&c.$

$$y = \sqrt{b + 2cx + dx^2}. \quad z = \text{fl. } \frac{ax}{\sqrt{b + 2cx + dx^2}}$$

$$F' = \frac{x^{\rho} N^{\prime\prime}}{r} - \frac{\rho F''}{ar}$$

$$F'' = \frac{x^{\rho-1} y N^{\prime\prime}}{r} - \frac{\rho-1 \cdot F'''}{ar} - \frac{dF'}{ar} - \frac{cG'}{ar}$$

$$F''' = \frac{x^{\rho-2} y^2 N^{\prime\prime}}{r} - \frac{\rho-2 \cdot F^{iv}}{ar} - \frac{2dF''}{ar} - \frac{2cG''}{ar}$$

$$F^{iv} = \frac{x^{\rho-3} y^3 N^{\prime\prime}}{r} - \frac{\rho-3 \cdot F^v}{ar} - \frac{3dF'''}{ar} - \frac{3cG'''}{ar}$$

( $\rho + 1$ ).

$$G' = \frac{x^{\rho-1} N^{\prime\prime}}{r} - \frac{\rho-1 \cdot G''}{ar}$$

$$G'' = \frac{x^{\rho-2} y N^{\prime\prime}}{r} - \frac{\rho-2 \cdot G'''}{ar} - \frac{dG'}{ar} - \frac{cH'}{ar}$$

( $\rho$ ).

$$H' = \frac{x^{\rho-2} N^{\prime\prime}}{r} - \frac{\rho-2 \cdot H''}{ar}$$

$$H'' = \frac{x^{\rho-3} y N^{\prime\prime}}{r} - \frac{\rho-3 \cdot H'''}{ar} - \frac{dH'}{ar} - \frac{cI'}{ar}$$

( $\rho - 1$ )

&c.

Hence, when  $\rho$  is a positive integer, all the fluents  $F', F'', (\rho + 1); G', G'', (\rho); H', H'', (\rho - 1); \&c.$  may be found.

T H E O -

T H E O R E M IX.

The fluent of  $\frac{x^n \dot{x}}{\sqrt{1-x^2}}$  is  $= \frac{2^{-n}}{-1} \times \text{fl. } N^{x\sqrt{-1}} - N^{-x\sqrt{-1}} \times \dot{x}$ ,

which is  $= K + \frac{2^{1-n}}{-1} \times \frac{\overset{(n)}{s}}{n} - n \cdot \frac{\overset{(n-2)}{s}}{n-2} + \frac{n \cdot n-1}{2} \cdot \frac{\overset{(n-4)}{s}}{n-4} \left(\frac{1}{2}n\right) \pm * \frac{1}{2} M \dot{x}$ ,

or  $= K - \frac{2^{1-n}}{-1} \times \frac{\overset{(n)}{c}}{n} - n \cdot \frac{\overset{(n-2)}{c}}{n-2} + \frac{n \cdot n-1}{2} \cdot \frac{\overset{(n-4)}{c}}{n-4} \left(\frac{n+1}{2}\right)$ ,

according as  $n$  is an even or an odd positive number.

$M = \frac{n \cdot n-1 \cdot n-2 \left(\frac{1}{2}n\right)}{1 \cdot 2 \cdot 3 \left(\frac{1}{2}n\right)}$ .  $z = \text{fl. } \frac{\dot{x}}{\sqrt{1-x^2}} = \text{Circ. Arc, rad. 1, sine } x$ .

$\overset{(n)}{s}$ ,  $\overset{(n-2)}{s}$ ,  $\overset{(n-4)}{s}$ , &c.  $\overset{(n)}{c}$ ,  $\overset{(n-2)}{c}$ ,  $\overset{(n-4)}{c}$ , &c. the sines and cosines of  $n \cdot z$ ,  $n-2 \cdot z$ ,  $n-4 \cdot z$ , &c. respectively.

\* + or - according as  $\frac{1}{2}n$  is even or odd.

S C H O L I U M.

In computations wherein exponentials are concerned, it may sometimes be necessary to observe, that,

$a \times \frac{N^{\frac{rx\sqrt{-1}}{a}} - N^{-\frac{rx\sqrt{-1}}{a}}}{2\sqrt{-1}}$  and  $a \times \frac{N^{\frac{rx\sqrt{-1}}{a}} + N^{-\frac{rx\sqrt{-1}}{a}}}{2}$

denoting the sine and cosine of  $rx$  respectively,

$x$  being an arc of the circle whose radius is  $= a$ ;

if that sine and cosine be also respectively denoted by

$s$  and  $c$ ,  $N^{\frac{rx\sqrt{-1}}{a}}$  will be  $= \frac{c+s\sqrt{-1}}{a}$ , and  $N^{-\frac{rx\sqrt{-1}}{a}} = \frac{c-s\sqrt{-1}}{a}$

☞ All the *Logarithms* mentioned in these Tables, and likewise in the Memoirs, are of the *hyperbolic* kind.

